

Electron-ion-ion triple-coincidence spectroscopic study of site-specific fragmentation caused by Si:2*p* core-level photoionization of F₃SiCH₂CH₂Si(CH₃)₃ vapor

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Site-specific fragmentation caused by Si:2*p* core-level photoionization of F₃SiCH₂CH₂Si(CH₃)₃ vapor was studied by means of high-resolution energy-selected-electron photoion-photoion triple-coincidence spectroscopy. The *ab initio* molecular orbital method was used for the theoretical description. F₃SiCH₂CH₂⁺-Si(CH₃)₃⁺ ion pairs were produced by the 2*p* photoionization of the Si atoms bonded to the three methyl groups, and SiF⁺-containing ion pairs were produced by the 2*p* photoionization of the Si atoms bonded to the three F atoms.

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The core-level chemical shift shown by atoms in a molecule depends on their chemical environments, so atoms that have the same atomic number but are in different chemical environments will show different chemical shifts. Furthermore, the core hole created by the photoionization in an atom is localized very close to the nucleus of that atom. These properties have been used to study site-specific fragmentation [1–4], in which bonds around the site of core-ionized atoms are dissociated selectively. Site-specific fragmentation is potentially useful for controlling chemical reactions and also offers possibilities for analyzing the structures and properties of molecules, molecular assemblies, and nanoscale devices by controlling matter at the level of individual atoms. To realize these exciting possibilities, we need to understand what controls fragmentation at the atomic level.

A molecule *M* including several atoms with the same atomic number in different chemical environments is expected to show core-photoelectron peaks reflecting chemical shifts differing site by site. In a normal Auger transition, the core-electron emission makes a valence electron fall into the core orbital and creates a valence hole that spatially overlaps the core-ionized atomic site in the molecule (*M*⁺). The electron falling into the core orbital usually gives its energy to another valence electron, which is emitted as a normal Auger electron, creating a second valence hole in the molecule (*M*²⁺). Since these valence holes weaken chemical bonds around the initially core-ionized atom, site-specific fragmentation (*M*²⁺ → F₁⁺+F₂⁺) often occurs around it. To observe such a fragmentation process selectively, we have to detect ion pairs (F₁⁺-F₂⁺) in coincidence with energy-selected photoelectrons that originate from each of several nonequivalent atomic sites in the different chemical environments. That is, we have to use energy-selected-photoelectron/Auger-electron photoion-photoion triple-coincidence (PEPIPICO/AEPIPICO) spectroscopy.

In this work we used the PEPIPICO and AEPIPICO methods and the *ab initio* molecular orbital (MO) method to study site-specific fragmentation caused by Si:2*p* core-level photoionizations of 1-trifluorosilyl-2-trimethylsilylethane [F₃SiCH₂CH₂Si(CH₃)₃, FSMSE] in the vapor phase. FSMSE is useful for such a study because the chemical environment of a Si atom bonded to three F atoms (here denoted Si[F]) is very different from that of one bonded to three methyl groups (Si[Me]). The dimethylene group (-CH₂CH₂-) between the two Si atoms maintains the site-specificity of the fragmentation by inhibiting intersite electron migration [5]. This study has provided clear evidence for almost 100% site-selectivity of the ion/ion-pair production in Si:2*p* photoionization.

The electron-ion coincidence measurements were performed using a hemispherical electron energy analyzer (Gammadata-Scienta SES-2002) and a time-of-flight (TOF) ion spectrometer, both of which were equipped with position-sensitive delay-line detectors [6]. FSMSE vapor was introduced to the photoionization region through a needle as an effusive beam. The coincidence apparatus was mounted behind the high-resolution plane-grating monochromator installed on the *c* branch of the soft x-ray figure-8 undulator beamline 27SU at the SPring-8 facility [7]. The data analysis method was modified from the original one [6] to extract PEPIPICO and AEPIPICO counts [8]. The experimental and computational methods and procedures used in the present work have been described in detail in previous papers [6,9] and an EPAPS document [10].

The photoelectron spectrum (PES) of FSMSE vapor shown in Fig. 1(a) has two peaks in the region of Si:2*p* photoemission. The peaks at lower and higher binding energies (106.4 and 109.8 eV) are, respectively, assigned to Si[Me]:2*p* and Si[F]:2*p* photoelectron emissions [11]. Thus, the chemical shift difference originating from the two Si sites of FSMSE vapor is clearly evident in the photoelectron spectrum. Computational results [10] show that the molecular motion immediately after emission of the Si:2*p* photoelectrons of FSMSE is negligible.

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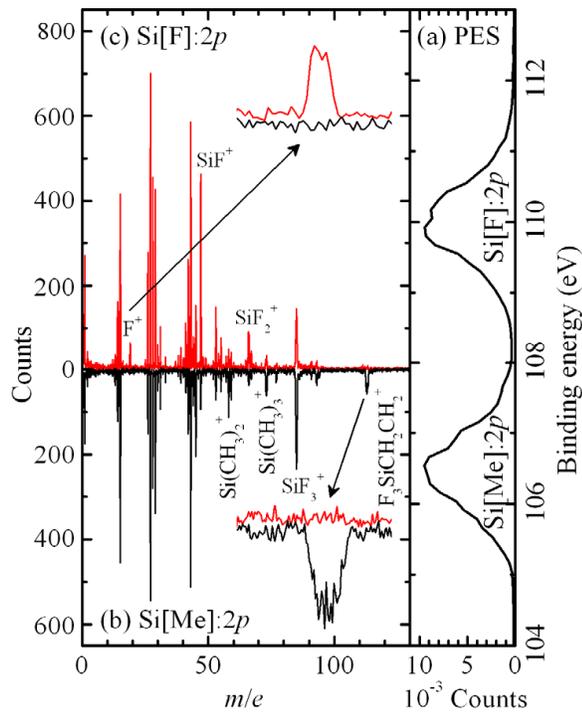
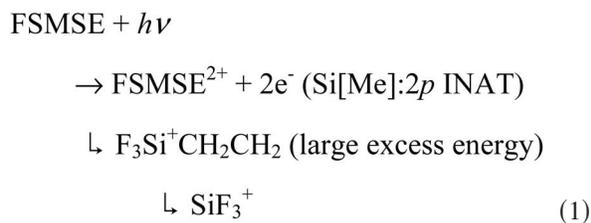


FIG. 1. (Color online) (a) Si:2*p* PES of FSMSE vapor obtained during the PEPICO measurement. (b) and (c) Ion mass spectra measured in coincidence with the Si[Me]:2*p* and Si[F]:2*p* photoelectrons. The insets show the regions of F⁺ and F₃SiCH₂CH₂⁺ with enlarged scales.

The ion mass spectra measured in coincidence with the Si[Me]:2*p* and Si[F]:2*p* photoelectrons (PEPICO spectra) are shown in Figs. 1(b) and 1(c). The peaks corresponding to C_{*n*}H_{*m*}⁺ groups (*m/e* < 40) are larger than those corresponding to ions containing F atoms because C_{*n*}H_{*m*} is easily ionized than F [12]. As seen in Fig. 1 and described below, site-specific fragmentation is clearly evident in the ion mass spectra. The Si[Me]:2*p* photoionization increased the production of Si(CH₃)₂⁺, Si(CH₃)₃⁺, SiF₃⁺, and F₃SiCH₂CH₂⁺, whereas the Si[F]:2*p* photoionization increased the production of F⁺, SiF⁺, and SiF₂⁺ [10]. The fragments most sensitive to the site of the initial energy deposition were F₃SiCH₂CH₂⁺ and F⁺, and the site-selectivity for the production of these ions was almost 100%. Except in the case of SiF₃⁺, the above site-specific bond dissociation occurs at the Si site where the photoionization has taken place. The Si[Me]:2*p* photoionization may have produced SiF₃⁺ in the following two-step dissociation mechanism:



where INAT denotes ionization and normal Auger transition. Similar two-step dissociation mechanisms could also account

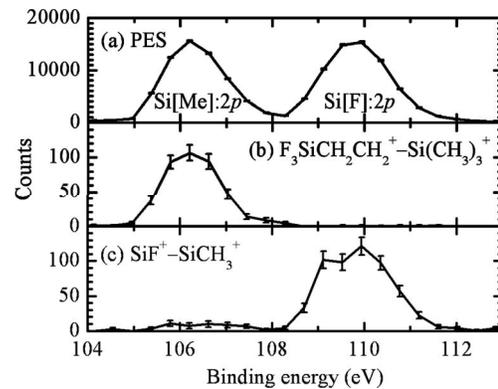


FIG. 2. (a) Si:2*p* PES of FSMSE vapor obtained during the PEPICO measurement. (b) and (c) Plots of the PEPICO counts of F₃SiCH₂CH₂⁺-Si(CH₃)₃⁺ and SiF⁺-SiCH₃⁺ versus the photoelectron binding energy.

for the production of Si(CH₃)₂⁺ from Si(CH₃)₃⁺ in the Si[Me]:2*p* photoionization and the production of F⁺ and/or SiF⁺ from SiF₂⁺ in the Si[F]:2*p* photoionization.

Figure 2 shows plots of the PEPICO counts of F₃SiCH₂CH₂⁺-Si(CH₃)₃⁺ and SiF⁺-SiCH₃⁺ versus the photoelectron binding energy, together with the PES of FSMSE. The plots for F₃SiCH₂CH₂⁺-Si(CH₃)₂⁺ and SiF₃⁺-Si(CH₃)₂⁺ are similar to that for F₃SiCH₂CH₂⁺-Si(CH₃)₃⁺, and the plots for SiF⁺-SiCH₂⁺, SiF⁺-SiC₂H₅⁺, and SiF⁺-SiH⁺ (and/or -C₂H₅⁺) are similar to that for SiF⁺-SiCH₃⁺ [10]. The ion pair most sensitive to the site of the initial energy deposition is F₃SiCH₂CH₂⁺-Si(CH₃)₃⁺, and the site-selectivity for its production is almost 100%. In similar plots, the PEPICO counts of F⁻ and SiF₂⁺-containing ion pairs are concealed by the background. Possible ion pairs with the same *m/e*, like F⁻-F⁺, were not detected because of the dead time of the ion detector.

Since the plot of the PEPICO count of F₃SiCH₂CH₂⁺-Si(CH₃)₃⁺ shows a peak at the binding energy of the Si[Me]:2*p* photoelectron, the Si[Me]-C₂H₄ bond dissociation and subsequent formation of F₃SiCH₂CH₂⁺-Si(CH₃)₃⁺ are thought to result from the Si[Me]:2*p* photoionization. The production of F₃SiCH₂CH₂⁺-Si(CH₃)₂⁺ and SiF₃⁺-Si(CH₃)₂⁺ by the Si[Me]:2*p* photoionization seems to be possible in two-step dissociation (and rearrangement) mechanisms similar to (1). Since the plots of the PEPICO counts of the SiF⁺-containing ion pairs show a peak at the binding energy of the Si[F]:2*p* photoelectron, we know that the Si[F]-F and Si[F]-C₂H₄ bond dissociations and subsequent formation of the ion pairs result from the Si[F]:2*p* photoionization.

In Fig. 2 and EPAPS document [10], the Si[F]:2*p* photoionization always breaks both of the Si[F]-C₂H₄-Si[Me] bonds and the Si(CH₃)₃ group falls apart, while the Si[Me]:2*p* photoionization leaves the Si[F]-C₂H₄ bond intact except in mechanism (1). So in general the Si[F]:2*p* photoionization seems to lead to more violent fragmentation than the Si[Me]:2*p* photoionization. This may be due to the different degrees of freedom at the Si[Me] and Si[F] sites. Although the dissociation in the SiF₃ moiety takes place only

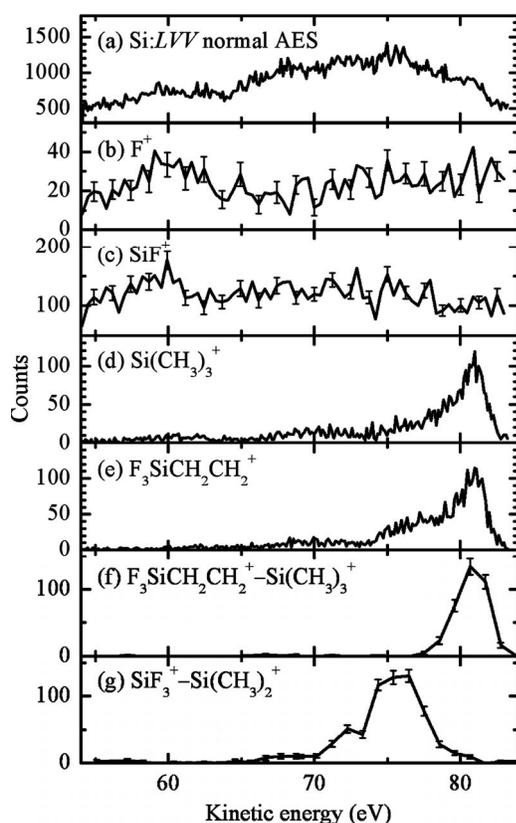


FIG. 3. (a) Si:LVV normal AES of FSMSE vapor obtained during the AEPICO measurement. (b)–(e) Plots of the AEPICO counts of F^+ , SiF^+ , $Si(CH_3)_3^+$, and $F_3SiCH_2CH_2^+$ versus the Auger-electron kinetic energy. (f) and (g) Plots of the AEPICO counts of $F_3SiCH_2CH_2^+-Si(CH_3)_3^+$ and $SiF_3^+-Si(CH_3)_2^+$.

at three Si-F bonds, the $Si(CH_3)_3$ moiety has a lot of bonds that can be broken: three Si-C bonds and nine C-H bonds. In other words, the Si[Me] site has a larger energy reservoir of the vibrational mode. As a result, in the Si[Me]: $2p$ photoionization, dissociation takes place mainly around the Si[Me] atom. In the Si[F]: $2p$ photoionization, however, it takes place not only around the Si[F] atom but also at various other bonds because of the fast energy transfer to the whole molecule.

Figure 3(a) shows the Si:LVV normal Auger-electron spectrum (AES) of FSMSE vapor, and Figs. 3(b), 3(c), 3(d), and 3(e), respectively, show plots of the Auger-electron photoion coincidence (AEPICO) counts of F^+ , SiF^+ , $Si(CH_3)_3^+$, and $F_3SiCH_2CH_2^+$ versus the Auger-electron kinetic energy. The plots of the AEPICO counts of $Si(CH_3)_2^+$ and SiF_3^+ show peaks at lower kinetic energies than those of $F_3SiCH_2CH_2^+$ and $Si(CH_3)_3^+$ [10]. Figures 3(f) and 3(g) show plots of the AEPICO counts of $F_3SiCH_2CH_2^+-Si(CH_3)_3^+$ and $SiF_3^+-Si(CH_3)_2^+$. In similar plots, the AEPICO counts of F^+ , SiF^+ , and SiF_2^+ -containing ion pairs are concealed by the background.

The plot of the AEPICO count of $F_3SiCH_2CH_2^+-Si(CH_3)_3^+$ [Fig. 3(f)], which is a site-specific ion pair formed by the Si[Me]: $2p$ photoionization (Fig. 2), shows a peak at the high-kinetic-energy edge of the Auger band (82 eV), as

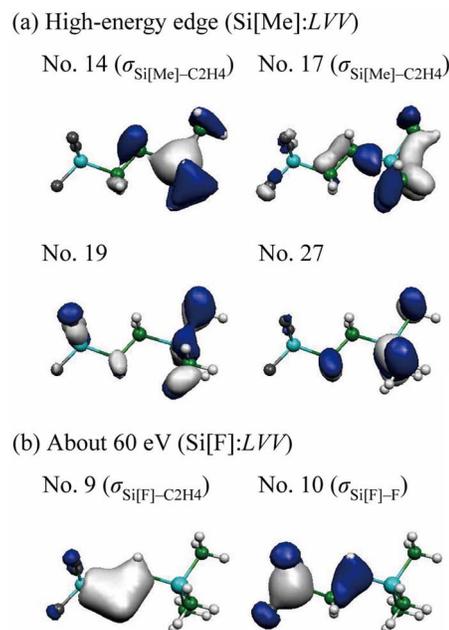
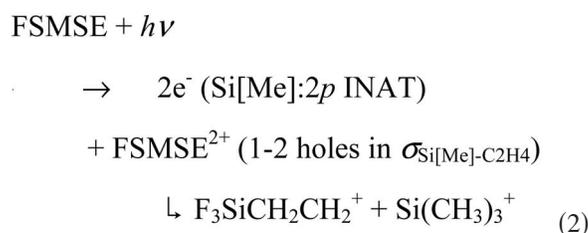


FIG. 4. (Color online) (a) Main MOs with one or two holes in the Si[Me]:LVV normal Auger final states corresponding to the high-kinetic-energy edge of the Auger band of FSMSE. The MOs are numbered in ascending order of orbital energy, and the highest occupied MO is No. 36. (b) Main MOs with one or two holes in the Si[F]:LVV normal Auger final states corresponding to an Auger-electron kinetic energy of about 60 eV.

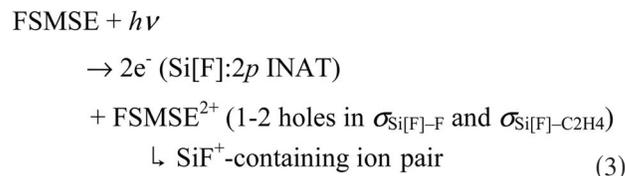
do the plots of the AEPICO counts of $Si(CH_3)_3^+$ and $F_3SiCH_2CH_2^+$ shown in Figs. 3(d) and 3(e). Therefore the Si[Me]- C_2H_4 bond dissociation and subsequent formation of $F_3SiCH_2CH_2^+-Si(CH_3)_3^+$ are results of an Auger decay with a kinetic energy of 82 eV. The emission of the Auger electron leads to various Si[Me]:LVV normal Auger final states, each of which has two holes in the valence MOs. We have used *ab initio* calculations to estimate the probabilities leading to the Auger final states, have selected some states with high probability, and in Fig. 4(a) have illustrated their valence MOs with one or two holes in the Si[Me]:LVV normal Auger final states at about 82 eV. The present finding that the Si[Me]- C_2H_4 bond dissociation and subsequent formation of $F_3SiCH_2CH_2^+-Si(CH_3)_3^+$ occur is consistent with the computational result that two MOs with a character of Si[Me]- C_2H_4 bonding ($\sigma_{Si[Me]-C_2H_4}$) have, with high probability, one or two holes in those Si[Me]:LVV normal Auger final states (MO Nos. 14 and 17). The hole creation in $\sigma_{Si[Me]-C_2H_4}$ is thought to result in the Si[Me]- C_2H_4 bond dissociation and the formation of $F_3SiCH_2CH_2^+-Si(CH_3)_3^+$:



The $\text{SiF}_3^+-\text{Si}(\text{CH}_3)_2^+$ ion pair, which like the $\text{F}_3\text{SiCH}_2\text{CH}_2^+-\text{Si}(\text{CH}_3)_3^+$ ion pair is site-specifically produced by the $\text{Si}[\text{Me}]:2p$ photoionization [10], seems to be formed from upper Auger final states with two holes in deeper-lying MOs because the AEPIPICO count of $\text{SiF}_3^+-\text{Si}(\text{CH}_3)_2^+$ shows a peak at a lower kinetic energy [Fig. 3(g)]. The plot for $\text{SiF}_3^+-\text{Si}(\text{CH}_3)_2\text{H}_2^+$ is similar to that for $\text{SiF}_3^+-\text{Si}(\text{CH}_3)_2^+$ [10]. The AEPIPICO counts of $\text{F}_3\text{SiCH}_2\text{CH}_2^+-\text{Si}(\text{CH}_3)_2\text{H}_2^+$ and some other ion pairs show peaks at other low kinetic energies [10].

The plots of the AEPIPICO counts of F^+ and SiF^+ , which are site-specific fragments formed by the $\text{Si}[\text{F}]:2p$ photoionization (Fig. 1 and [10]), show a hump at about 60 eV [Figs. 3(b) and 3(c)] as does the plot of the AEPIPICO count of F^+ in the condensed phase [13]. One can also see a hump at about 60 eV in the $\text{Si}:LVV$ normal AES [Fig. 3(a)]. By means of the PEPIPICO method, it is seen that some SiF^+ -containing ion pairs are efficiently produced by the $\text{Si}[\text{F}]:2p$ photoionization [Fig. 2(c) and [10]]. Therefore the $\text{Si}[\text{F}]-\text{F}$ and $\text{Si}[\text{F}]-\text{C}_2\text{H}_4$ bond dissociations and subsequent formation of the SiF^+ -containing ion pairs are thought to result from the $\text{Si}[\text{F}]:LVV$ normal Auger final states corresponding to an Auger-electron kinetic energy of about 60 eV. This interpretation is consistent with the computational result that a MO with a character of $\text{Si}[\text{F}]-\text{F}$ bonding ($\sigma_{\text{Si}[\text{F}]-\text{F}}$) and another MO with a character of $\text{Si}[\text{F}]-\text{C}_2\text{H}_4$ bonding ($\sigma_{\text{Si}[\text{F}]-\text{C}_2\text{H}_4}$) have, with high probability, one or two holes in the Auger

final states at about 60 eV [Fig. 4(b)]. The hole creation in $\sigma_{\text{Si}[\text{F}]-\text{F}}$ and $\sigma_{\text{Si}[\text{F}]-\text{C}_2\text{H}_4}$ is thought to result in the dissociation of the $\text{Si}[\text{F}]-\text{F}$ and $\text{Si}[\text{F}]-\text{C}_2\text{H}_4$ bonds and the formation of the SiF^+ -containing ion pairs:



In conclusion, site-specific fragmentation caused by the $\text{Si}:2p$ core-level photoionization of FSMSE vapor was measured by means of high-resolution PEPIPICO/AEPIPICO spectroscopy and was interpreted by means of the *ab initio* MO method. $\text{F}_3\text{SiCH}_2\text{CH}_2^+-\text{Si}(\text{CH}_3)_3^+$ and some SiF^+ -containing ion pairs were formed according to mechanisms (2) and (3). Site-specific fragmentation thus offers an approach to controlling chemical reactions [14] by controlling the sites at which holes are created (e.g., by a resonant excitation). This elucidation of the details of the fragment production mechanism brings the goal of chemical synthesis and some other applications one step closer to realization, but further investigations are needed.

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- [1] W. Eberhardt *et al.*, Phys. Rev. Lett. **50**, 1038 (1983).
 [2] K. Müller-Dethlefs, M. Sander, L. A. Chewter, and E. W. Schlag, J. Phys. Chem. **88**, 6098 (1984).
 [3] R. Romberg *et al.*, Phys. Rev. Lett. **84**, 374 (2000).
 [4] S. Nagaoka *et al.*, Chem. Phys. Lett. **412**, 459 (2005) and references cited therein.
 [5] S. Nagaoka *et al.*, Int. J. Mass Spectrom. Ion Process. **171**, 95 (1997).
 [6] G. Prümper *et al.*, J. Electron Spectrosc. Relat. Phenom. **144-147**, 227 (2005).
 [7] H. Ohashi *et al.*, Nucl. Instrum. Methods Phys. Res. A **467-468**, 529 (2001).
 [8] G. Prümper and K. Ueda (unpublished).
 [9] M. Mitani, O. Takahashi, K. Saito, and S. Iwata, J. Electron Spectrosc. Relat. Phenom. **128**, 103 (2003).
 [10] See EPAPS Document No. E-PLRAAN-75-R10702 for the experimental and computational details, the description of the comparison with the condensed phase, the computational animations of the molecular dynamics immediately after emission of the $\text{Si}[\text{Me}]:2p$ and $\text{Si}[\text{F}]:2p$ photoelectrons, and the detailed results of the coincidence measurements. This document can be reached via a direct link in the online article's HTML reference section or via the EPAPS homepage (<http://www.aip.org/pubservs/epaps.html>).
 [11] S. Nagaoka *et al.*, J. Chem. Phys. **107**, 10751 (1997).
 [12] J. D. Bozek, K. H. Tan, G. M. Bancroft, and K. J. Fu, Chem. Phys. **158**, 171 (1991).
 [13] K. Mase *et al.* (unpublished).
 [14] K. Tanaka *et al.*, J. Electron Spectrosc. Relat. Phenom. **119**, 255 (2001).