

**Estimation of Effective Thermal Conductivity of  
Graphite flakes/Al Composites by Using Two-  
Dimensional Microstructure Images**

(二次元微細組織画像を用いたグラファイトフレーク/Al  
複合材料の有効熱伝導率の推定)

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# *Chapter 1*

## **Background and Objective**

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## 1.1 Introduction

Today is an era of the rapid development of electronic technology. As shown in Figure 1-1, various electronic products are gradually integrated into people's daily lives, making people's lives convenient, comfortable, and safe. Therefore, ensuring the electronic products' efficiency, stability, and safety is essential to people's lives.

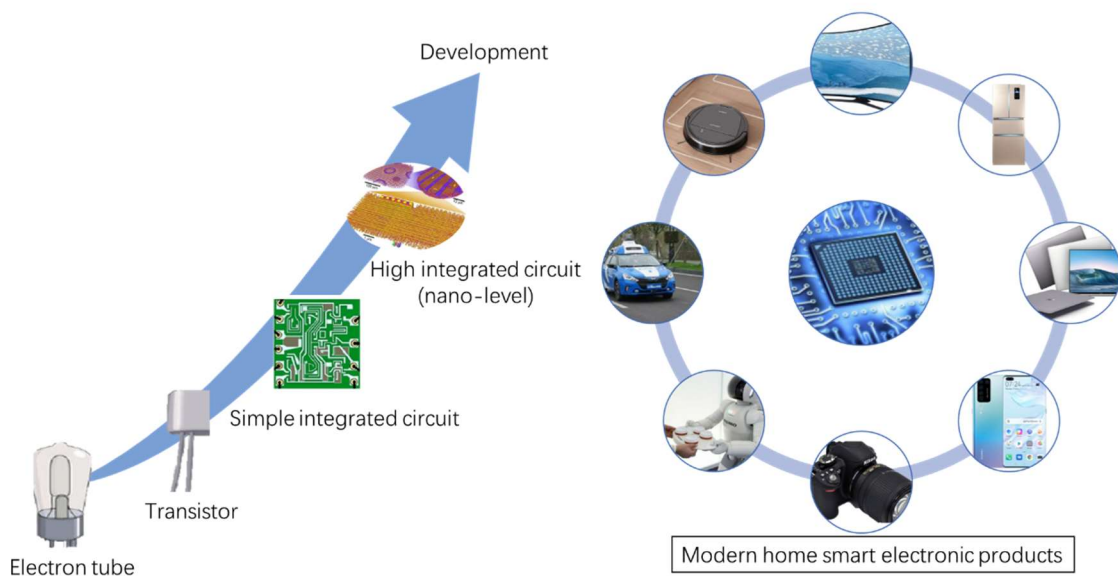


Figure 1-1 The development of microelectronics and modern smart products.

However, with the rapid development of electronic technology in the direction of high power, high frequency, miniaturization, and integration, the heat generated by the increase in the power density of electronic components is increased dramatically. As a result, the generated heat would heat electronic components, affecting their working effectiveness and life span, and even resulting in a malfunction of the device if not dissipation in time<sup>[1]</sup>. In other words, the mismatch in the heat dissipation performance and the thermal expansion coefficient between the materials used to make the components cause the electronic components' damage and failure. The key to solving this problem is to develop new electronic materials with high thermal conductivity (TC) and lower coefficient of thermal expansion (CTE) <sup>[2]</sup>. Semiconductors are the core component of electronic products; hence, electronic packaging materials must match

thermal properties to semiconductors. Table 1-1 lists the thermal properties of semiconductors materials.

Table 1-1 thermal properties of semiconductor materials <sup>[1, 3-5]</sup>.

Materials	Density / $\text{g}\cdot\text{m}^{-3}$	TC / $\text{W}\cdot\text{m}^{-1}\cdot\text{K}^{-1}$	CTE / $10^{-6} / \text{K}^{-1}$
Silicon (Si)	2.3-2.34	139-150	2.8-4.2
Silicon carbide (SiC)	3.1-3.21	80-490	2.8-5.0
Gallium arsenide (GaAs)	5.3	46-58	5.8-5.9
Gallium nitride (GaN)	6.1	150	3.2

Traditional electronic packaging materials contain metals, ceramics, and composite materials. Table 1-2 shows the TC and CTE of traditional electronic packaging materials. However, the traditional electronic packaging materials, such as Copper (Cu), Argentum (Ag), Aurum (Au), Alumina ( $\text{Al}_2\text{O}_3$ ), W-Cu, etc., all have high density or high TC or high CTE. As a result, they can no longer meet electronic products' current comprehensive development requirements, such as high integration, efficient heat dissipation, and green environmental protection<sup>[6, 7]</sup>.

Table 1-2 thermal properties of traditional electronic packaging materials.

Materials	Density / $\text{g}\cdot\text{m}^{-3}$	TC / $\text{W}\cdot\text{m}^{-1}\cdot\text{K}^{-1}$	CTE / $10^{-6} / \text{K}^{-1}$
Al	2.7	236	23.0
Cu	8.9	400	17.0
Ag	10.49	429	18.9
Au	19.3	318	14.2
$\text{Al}_2\text{O}_3$	3.9	20	6.5-6.7
BeO	2.9	250-275	6.7-8.0
SiC	3.2	270	4.5
AlN	3.3-4.5	220	3.3-4.5
W-Cu	8.4	167	6.5
Mo-Cu		220-270	7.9-9.3

To cater for the future development characteristics of electronic products: e.g., high integration of integrated circuits and the miniaturization, portability, efficiency, high power, high stability, and low cost, the new generation of electronic packaging materials is required to have the following mechanical and thermal properties [8]:

1, High TC: Timely homogenize and dissipate the heat generated by the semiconductor.

2, Low CTE: Reduce the thermal stress between the semiconductor and electronic packaging materials.

3, Sufficient strength, stiffness, and toughness: Play a good role in supporting and protecting semiconductors and devices.

4, Lightweight: The density is as low as possible to facilitate the lightweight design of the device structure.

5, Cost control and competitive requirements: high yield, suitable for mass production, low price.

Metal-matrix composites have a high TC, low expansion of non-metallic reinforcement, easy processing of metal-matrix, and good plastic toughness. Therefore, metal matrix composites have attracted significant attention from researchers. Al is of low density, low cost, and high TC, and it is the commonly used metal matrix of the composites. The reinforcements include mainly carbon materials (carbon fiber, graphite, diamond), silicon carbide particles (SiC), silicon particles (Si), etc. Table 1-3 shows the TC, CTE, and density of those reinforcements.

SiC/Al, Si/Al, Diamond/Al are the most mature metal matrix composites. These composites have been widely used in the world and have gradually replaced the traditional electronic packaging materials. However, the metal matrix composites currently used in electronic packaging, e.g., SiC/Al, Si/Al, Diamond/Al, show high prices and TC lower than expected in large-scale industrial production. Figure 1-2 shows the density, TC, and CTE of the composites used in electronic packaging.

Table 1-3 Thermal properties of reinforcements<sup>[1, 7, 9-11]</sup>.

Materials	Density/ $\text{g}\cdot\text{cm}^{-3}$	TC / $\text{W}\cdot\text{m}^{-1}\cdot\text{K}^{-1}$	CTE / $10^{-6}/\text{K}$
Carbon fiber	1.75 ~ 2.19	530-1200	-1.45 ~ 7
Graphite	2.23	~3000 (xy) 6 ~ 38 (z)	-1.5 ~ -1.0 (xy) 29 (z)
Diamond	3.51	~2500	1.3
SiC	3.2	270	4.5
Si	2.3	150	4.1
BeO	2.9	250	6.7

Where a and b: the TC and ETC could be less than those of crystalline graphite due to the defects and impurities.

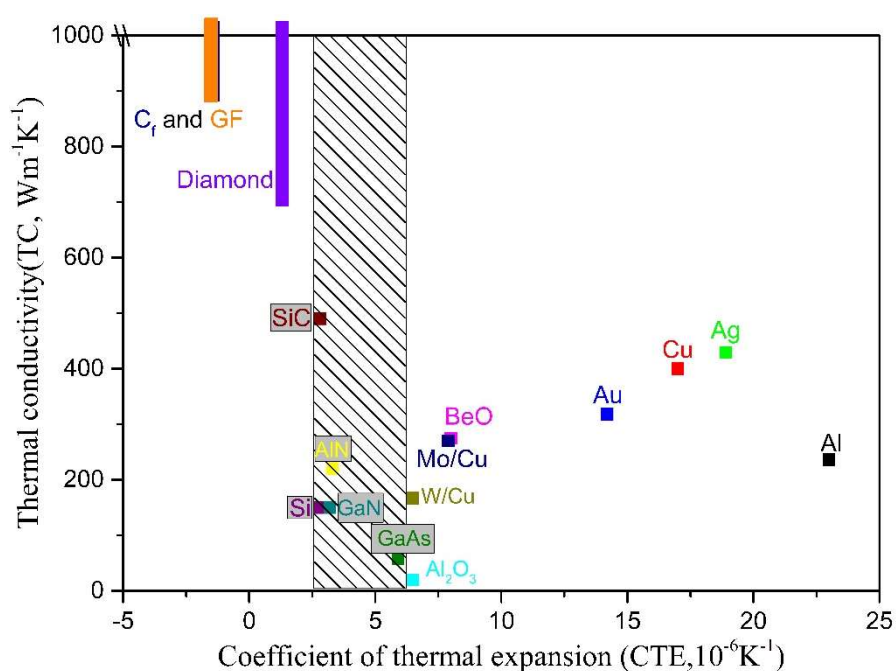


Figure 1-2 TC as a function of CTE for composites.

It is fortunate that carbon fibers and graphite flakes also have excellent thermal properties and are low-cost. Particularly for graphite flake, their TCs are values ( $\sim 2200$ ,  $\sim 38$ )  $\text{W m}^{-1} \text{K}^{-1}$  [9, 12, 13] in the basal-plane and the out-plane, respectively, and their CTE are  $(-1.0, 8) \times 10^{-6} \text{K}^{-1}$  [10, 14, 15] in the basal-plane and the out-plane, respectively. In

view of GFs' excellent thermal properties, GFs/Al composites are considered to be one of the most promising candidates for electronic packaging. Many experimental studies of the effective thermal conductivity (ETC) of GFs/Al composites have been conducted extensively, e.g., Chen et al. [14] reported that the ETC of 80 vol.% GFs/Al composites ( $783 \text{ W m}^{-1} \text{ K}^{-1}$ ) was higher than that of 80 vol.% diamond/copper composites [16] ( $724 \text{ W m}^{-1} \text{ K}^{-1}$ ). Li et al. [17] reported that the ETC of 70 vol.% GFs (graphite flakes)/Al composite was  $714 \text{ W m}^{-1} \text{ K}^{-1}$  in the plane parallel to the GFs layers. Comparison with SiC/Al, Si/Al, Diamond/Al composites, GFs/Al composites show higher TC, lower CTE, lower cost, and better machining performance. However, the ETC of GFs/Al composites is not high as the expected theoretical value due to the following deficiencies<sup>[10, 18-21]</sup>:

(1) GFs show strong anisotropic TC, resulting in the orientation of GFs has a great effect on the effective thermal conductivity (ETC) of GFs/Al composites.

(2) Interface reaction can occur readily during the fabricating process of GFs/Al composites, and the harmful formation of  $\text{Al}_4\text{C}_3$  is often observed at the GFs-Al interface. The  $\text{Al}_4\text{C}_3$  is a brittle compound, and it can reduce the plasticity of the composites and promote accelerated fatigue crack growth rates.

(3) The interfacial thermal resistance at the interface between GFs and Al considerably affects composites' ETC.

Additionally, to predict the ETC of composites, many theoretical models are proposed in previous studies [10, 11, 13, 14, 17, 18, 22-32]. However, those theoretical models are not suitable for calculating the ETC of GFs/Al composites due to the diversity orientation of the GFs in composites. Moreover, to evaluate the interfacial thermal resistance at the GFs-Al interface, the commonly used methods include the phonon diffuse mismatch model (DMM) [33], acoustic mismatch model (AMM)<sup>[34]</sup>, etc. However, those models are based on an idea interface to calculate interfacial thermal resistance. In reality, the contact interface microstructure of two solids is extremely complex. Contact interfaces are affected by many factors, such as the thermophysical properties, the hardness, shape, roughness of the material, the gap at the interface, the



roughness of the contact interface, as well as the pressure and temperature during composite fabrication. Furthermore, it is almost impossible to obtain the direct relationship between those influence factors and the contact interface microstructure. Thus, those models have specific limitations for calculating interface thermal resistance, and they are not entirely reliable for calculating interfacial thermal resistance.

This study reviews many previous studies on the ETC of GFs/Al composites and develops a new method to calculate the ETC of GFs/Al composites. Research method: a newly developed two-dimensional image-based model. Research contents: 1, Comparing the thermal conductivity of GFs in the 3D model and the 2D model, and determining the relationship between the thermal conductivity of GFs in the 3D model and the 2D model; 2, Study the effect degrees of orientation of GFs and interfacial thermal resistance on the effective thermal conductivity of GFs/Al composites using two-dimensional image-based simulations.

## **1.2 Experimental research on ETC of GFs/Al composites**

### **1.2.1 Fabrication methods of GFs/Al composites**

Metal matrix composites can be fabricated via a variety of methods like stir casting, squeeze casting, ultrasonic infiltration, pressure infiltration, spark plasma sintering, and hot pressing, etc. However, not all fabrication techniques are suitable for fabricating GFs/Al composites since there are some challenges to overcome, such as the oriented GFs, avoiding the formation of  $Al_4C_3$ , and reducing the interfacial thermal resistance. In this part, some common methods for fabricating GFs/Al composites are reviewed, along with their advantages and disadvantages. The fabrication methods can be classified into two routes: powder metallurgy process and liquid metallurgy process.

#### **1.2.1.1 Powder Metallurgy Processing**

The powder metallurgy process (PMP) is a low-temperature sintering technology that can effectively inhibit the interfacial reaction and the formation of  $Al_4C_3$ . The PMP methods can be divided into spark plasma sintering (SPS) [35-37] and vacuum hot pressing [14, 23, 26, 27]. Figure 1-3 shows the schematic diagram for SPS and vacuum hot pressing.

The process of using the PMP methods to prepare GFs/Al composites is to fill the mixtures of raw materials Al and GFs into a graphite mold for sintering. Some techniques in past studies were proposed to ensure GFs have a consistent orientation in composites. Such as that, to stack the GFs by layers, Chen et al. [14] tapped the graphite mold while filling naturally; Xue et al. [26] used a shaking device to vibrate the graphite mold filled with mixtures of Al and GFs, hoping to obtain the oriented GFs; Huang et al. [23] fill the graphite mold with layers of artificial graphite films, and Al foils, the orientation of graphite film and Al foil can be well controlled.

The formation of  $Al_4C_3$  at the GFs-Al interface is thermodynamically favorable, and it can damage the mechanical and thermal properties. To avoid the formation of  $Al_4C_3$ , Chen et al. [14] used some approaches, including a lower temperature than

melting point of Al during the hot-pressing process, and GFs with nearly 100% graphitization degree and a larger flat surface, to prevent the formation of  $\text{Al}_4\text{C}_3$ . Salvo [38] and Candan et al. [39] filled a certain minimum critical Si content to prevent the formation of  $\text{Al}_4\text{C}_3$ . Similarly, Xue et al. [26] coated GFs surface with Si to prevent the formation of  $\text{Al}_4\text{C}_3$ , and Kurita et al. [27] added Al-Si<sub>11.3at%</sub> alloy into the raw materials Al powders to avoid the formation of  $\text{Al}_4\text{C}_3$  during fabricated GFs/Al composites. The introduction of silicon not only prevents the formation of  $\text{Al}_4\text{C}_3$ , but silicon can effectively fill the pores at the interface between aluminum and graphite. Meanwhile, the disadvantages of the introduction of silicon were also revealed, which caused the microstructure of GFs to disarray in proximity to the GFs-Al interface, which can weaken the interlaminar strength between GF in proximity to the Al/GF interface and Al/GF interface [27].

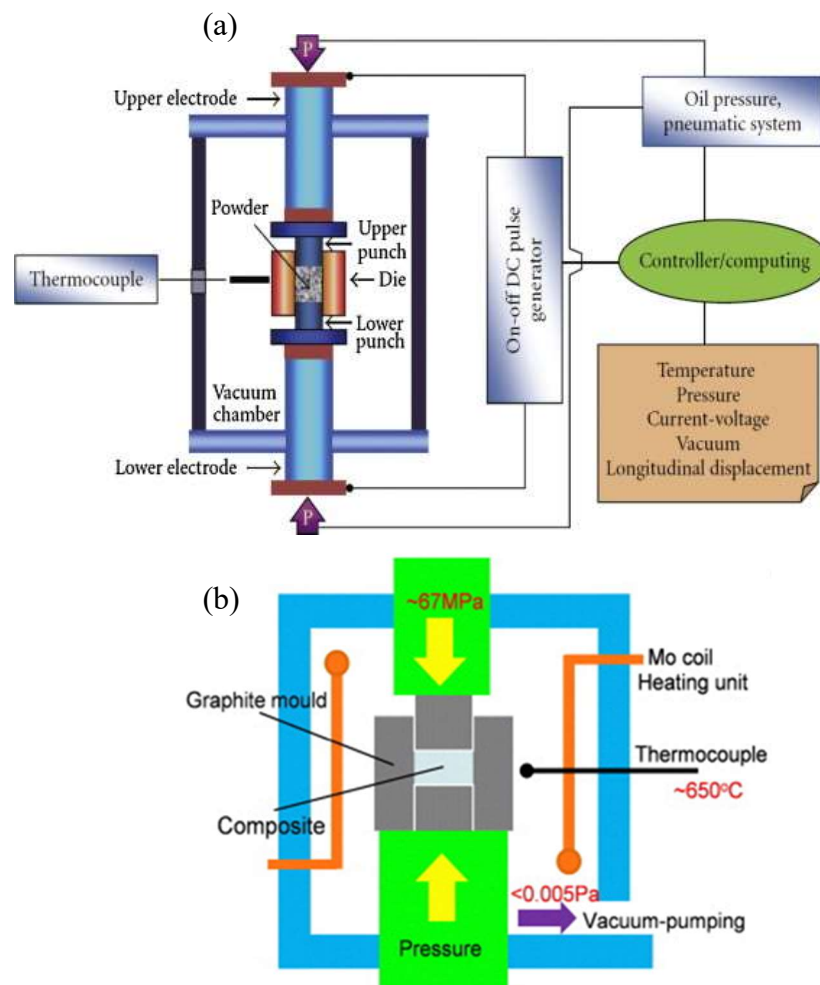


Figure 1-3 Schematic diagram for (a) spark plasma sintering [40], (b) vacuum hot pressing [41].

### 1.2.1.2 Liquid metallurgy processing

The liquid metallurgy process is also widely used to fabricate GFs/Al composites. The commonly used methods include gas pressure infiltration <sup>[11, 13, 31]</sup> and squeeze casting <sup>[17, 28, 29]</sup>. Figure 1-4 shows the schematic diagram of gas pressure infiltration and squeeze casting.

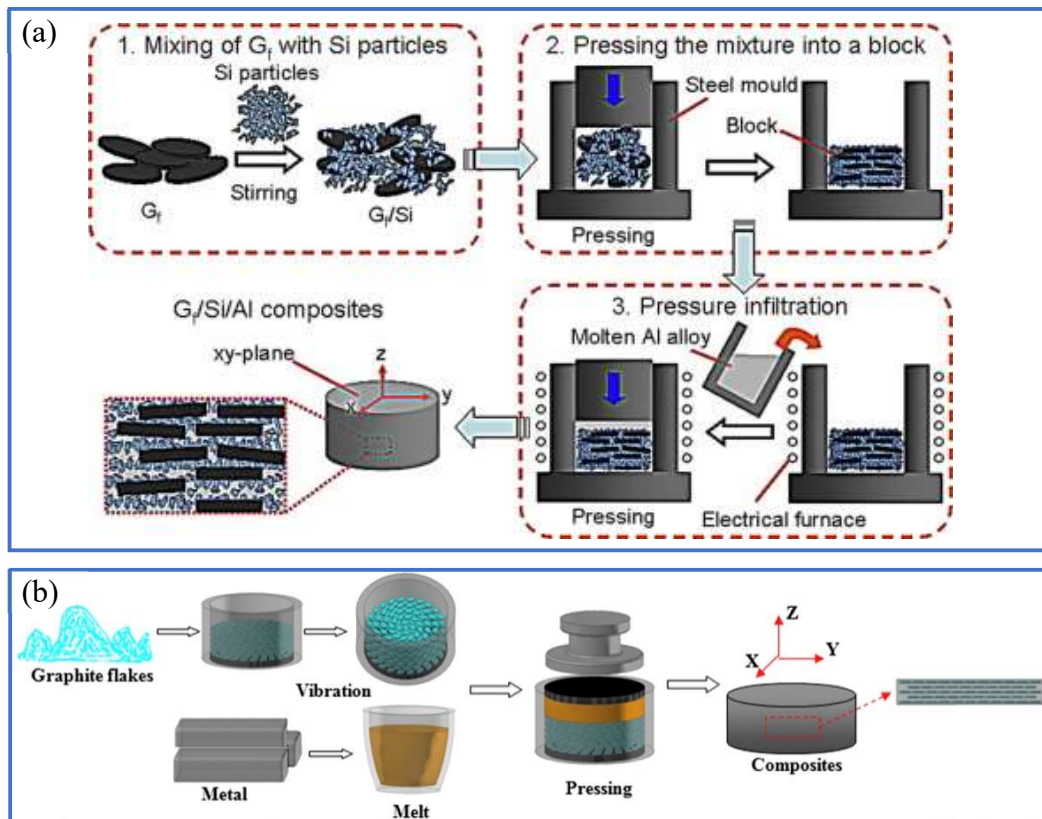


Figure 1-4 Schematic diagram of the fabrication process of GFs/Al composites, (a) gas pressure infiltration method <sup>[11]</sup>, (b) squeeze casting method <sup>[17]</sup>.

The common features of gas pressure infiltration and squeeze casting in GF/Al composite fabrication are: First, it is to put the raw material GFs into a mold and press GFs into preforms; Then, the molten Al is poured into a preheated mold containing the GFs preforms, and apply pressure to infiltrate the molten Al into the preforms. As a result, GFs are able to achieve uniform orientations through the fabrication of GFs preforms, and as reported in the papers <sup>[11, 13, 17, 28, 29]</sup>, GFs were stacked by layers in GFs/Al composites.

As known that high temperature is an essential condition for liquid metallurgy to improve the fluidity of molten Al, but conversely, it promotes the interface reaction as well. According to Okura's studies <sup>[42]</sup> on the rate of formation of intermetallic compounds in Al matrix-carbon fiber composites, the interface reaction rate can be evaluated by:  $X = [2kt \cdot \exp(-Q/RT)]^{1/2}$ .  $X$  is the thickness of  $Al_4C_3$  reaction layer,  $k$  is a constant,  $t$  is heating time,  $Q$  is the activation energy, and  $T$  is temperature. The formula shows that the formation of  $Al_4C_3$  is related to the temperature  $T$  and the heating time  $t$ . Consequently, the rapid cooling method<sup>[17, 29, 31]</sup> in the pressure infiltration process and squeeze casting process is one of the most used methods to suppress the formation of  $Al_4C_3$ . In addition, coating Si on the surface of GFs, or mixing a small amount of Si with raw GFs, or replacing pure aluminum with Al-Si alloys are also often used to prevent the formation of  $Al_4C_3$  <sup>[11, 13, 20, 28, 29]</sup>.

### 1.2.2 Problems of previous studies

Lots of experimental studies, as mentioned above 1.2.1 and 1.2.2, have tried to conquer the problems of the orientation of GFs and the formation of  $Al_4C_3$ . But the measured ETCs of GFs/Al composites were not as high as the expected theoretical ETC. It is evident that the high TC of GFs cannot be fully utilized in GFs/Al composites. The reasons can be analyzed from the following aspects.

Firstly, it is the method to control the orientation of GFs. The vibration method is often used to obtain oriented GFs in powder metallurgy processing. However, it is almost impossible to ensure that all the graphite flakes are entirely parallel to each other by vibrating the mold containing the mixtures of Al and GFs. Moreover, the pressure applied during the thermoforming process is usually given along one direction, the orientations of GFs cannot be well controlled, and angles between GFs basal-plane and the surface of GFs/Al composites are easy to form. Therefore, for the GFs/Al composites with a high-volume fraction of Gfs, it is more difficult to control GFs orientation by the vibration method. In the liquid metallurgy processing, obtaining the oriented GFs is to pre-press the GFs into a preform. However, the infiltration process

of molten Al is challenging because the distance between GFs is minimal. Especially the prefabricated GFs preforms with a preferred orientation, and there is almost no space between GFs<sup>[10, 11, 13, 14, 17, 29]</sup>. In order to enhance the infiltration process, the third phases (SiC particles, Si particles, and short carbon fiber) were added into GFs preforms<sup>[10, 11, 13, 29, 43-46]</sup>. But the third phase would reduce the ETC of GFs/Al composites due to the low TC of the third phase itself, and it may destroy the parallel stacking of GFs in the preform.

Secondly, it is the approach to avoid the formation of  $Al_4C_3$ . There are two approaches used usually in experiments: 1 using highly graphitized GF and low sintering temperature.; 2 introducing the third phase (Si or SiC). However, according to Huang<sup>[23]</sup> and Kurita et al.<sup>[27]</sup>, low sintering temperature can cause low density and more pores embedded in GFs/Al composites. The pores can reduce the ETC of composites<sup>[47]</sup>. In addition, the low TC of the third phase (Si or SiC) also lowers the ETC of GFs/Al composites.

Thirdly, it is the effect of interfacial thermal resistance on the ETC of composites. Many studies have proposed that the interfacial thermal resistance between carbon and Al dramatically affects the ETC of composites<sup>[13, 23, 29, 31, 48-54]</sup>. Besides, previous studies have investigated the effects of the shape<sup>[55-57]</sup> and size<sup>[14, 17, 18, 52, 53, 58, 59]</sup> of enhanced phase on the ETC of composites. For example, according to the studies of the effect of particle size on the ETC, the smaller filler particles can result in a lower ETC of composites because they caused a more extensive interface area<sup>[14, 17, 18]</sup>.

### 1.3 Theoretical research on interfacial thermal resistance

Interfacial thermal resistance is derived from the Kapitza<sup>[60]</sup> thermal resistance discovered in 1941. Interfacial thermal resistance is defined as the ratio of the temperature discontinuity at the interface to the power per unit area flowing across that interface<sup>[33]</sup>. Up to now, lots of scholars have studied the interface thermal resistance [33, 34, 61-71].

To determine interfacial thermal resistance, the commonly used experimental methods include steady-state measurement<sup>[72-74]</sup>, transient measurement (such as laser photothermal measurement<sup>[75, 76]</sup> and photoacoustic method<sup>[77-79]</sup>, thermal imaging method<sup>[72]</sup>, “flash” flash method<sup>[80, 81]</sup>, etc.), and micro/nano-scale measurement methods (such as transient thermoreflectance (TTR) technology and  $3\omega$  method<sup>[33, 82, 83]</sup>). However, those measurements are very strict, and the measurement results are not accurate enough because those measurements are susceptible to many external factors.

Therefore, many theoretical models are developed to predict interfacial thermal resistance, and interface thermal resistance is evaluated mainly from macro and micro perspectives.

#### 1.3.1 Macro theory

At the macroscopic level, the interfacial thermal resistance comes from the incomplete contact of the material interface<sup>[84, 85]</sup>. The interfacial thermal resistance,  $R_c$ , can be expressed by  $R_c = R_d + R_L$ .  $R_d$  is the thermal resistance caused by heat diffusion and heat shrinkage resistance at the contact interface,  $R_L$  is the thermal resistance of the contact area at the contact interface<sup>[86]</sup>. Microstructure at the contact interface dominates the interfacial thermal resistance. However, the microstructure at the contact interface is affected by many factors, such as surface roughness, microscopic morphology, elastoplastic deformation of surface peaks, and stages in the surfaces of contact materials. Based on those influencing factors, some theoretical

methods have been proposed to calculate interface thermal resistance, including CMY plastic contact model [61, 86], Multi-scale contact model [87, 88], Monte-Carlo simulation<sup>[89]</sup>, Micro-scale contact model [90], etc. But those macroscopic models and simulations all have specific usage restrictions due to the complexity of the interface microstructure.

### 1.3.2 Micro theory

At the micro-level, heat is transferred by phonons, electrons, and photons. Therefore, the calculation mechanism of interfacial thermal resistance is to calculate phonon and electron couplings. In the past decades, different microscale models were proposed to predict interfacial thermal resistance. The common models include interfacial layer model (ILM) [91, 92], AMM model<sup>[34]</sup>, DMM model and Boltzmann mesoscopic simulation<sup>[68]</sup>.

ILM model considers that the contact interface consists of subsurface layers on the two sides of the contact interface and the middle contact interface layer. The thickness of the subsurface layer is 1 micron. At the contact interface, hot carriers transfer energy in the form of radiation. The subsurface layer is mainly composed of local materials, but it also contains the atoms diffused from the other side of the interface and some other microstructures such as microprotrusions, dislocations, and oxidation from the other side of the interface. These microstructures cause the hot carriers scattering, the lattice defect scattering, dislocation scattering, and others scattering. The hot carriers scattering at the interface caused the interfacial thermal resistance. The interfacial thermal resistance,  $R_c$ , can be calculated by the following formulas:

$$R_c = \sum R_{CI} + R_{Sub} \quad \dots\dots (1.2.1.2-1)$$

$$R = \frac{\Delta T}{q} \quad \dots\dots (1.2.1.2-2)$$



$$q_{CI} = \frac{K_B^4}{8\pi^2\hbar^3} (\sum_{CI} V_{1,j}^{-2}) \alpha_{1 \rightarrow 2} \left( T_1^4 \int_{z=0}^{-\hbar\omega/K_B T_1} \frac{z^3 dz}{\exp(z)-1} - T_2^4 \int_{z=0}^{-\hbar\omega/K_B T_2} \frac{z^3 dz}{\exp(z)-1} \right) \dots\dots (1.2.1.2-3)$$

$$q_{Sub} = \frac{-1}{3d} CVl\Delta T \dots\dots (1.2.1.2-4)$$

where subscripts *CI* and *Sub* are the contact interface layer and subsurface layer, respectively;  $R$  is thermal resistance;  $\Delta T$  and  $q$  are the temperature difference of the interface layer and the heat flux density through the interface, respectively;  $K_B$  is Boltzmann constant;  $V$  is mean sounding velocity;  $\alpha$  is the transmittance of phonon;  $T$  is temperature;  $\hbar$  is Planck's constant;  $\omega$  is angular frequency;  $z$  is the acoustic resistance of solid materials;  $d$  is the thickness of the subsurface layer;  $C$  is volume-specific heat;  $l$  is the mean free path of phonon scattering process;  $j$  is the phonon mode for 1 and 2. However, the interface thickness  $d$  is closely related to the interface heat flux density, contact area, load pressure on the contact surface, and other factors. Therefore, the correlation between them cannot be obtained directly. Thus, it is not easy to obtain the interface thickness  $d$ .

AMM model assumes that the contact interface of two solid materials is a plane, and phonons are treated as continuous plane waves. The propagation of phonons on the interface is complete specular reflection and refraction, and the propagation is no possibility of diffuse scattering. The interfacial thermal resistance can be calculated as follows:

$$R_c = \left[ \frac{\pi^2 K_B^4 \Gamma_{1 \rightarrow 2}}{15 \hbar^3} \left( \sum_j V_{1,j}^{-2} \Gamma_{1,j} \right) \right]^{-1} T^{-3} \dots\dots (1.2.1.2-5)$$

$$\Gamma_{1,j} = \int_0^{\pi/2} \alpha_{1 \rightarrow 2}(\theta, j) \sin \theta \cos \theta d\theta \dots\dots (1.2.1.2-6)$$

where  $\theta$  is phonon incident angle. Under low-temperature conditions, phonon

wavelength is much larger than interface roughness size and interface defect size, and the contact interface can be approximately regarded as a plane. The AMM model is reliable to calculate the interfacial thermal resistance under low-temperature conditions. However, under high-temperature conditions, the wavelength of phonons is reduced to close to the size of interface roughness, so the contact interface of two solids cannot be approximated as a plane, and the possibility of phonons scattering at the interface is very high. The interfacial thermal resistance calculated by the AMM model may have a more significant error.

DMM model assumes that phonons are completely elastically scattered at the contact interface of two solid materials, and the transmittance of phonons ( $\alpha$ ) is related to the phonon density of states of the material. Therefore, the calculation expressions of the interface thermal resistance are as follows:

$$R_c = \left( \frac{\pi^2 K_B^4 V_1^{-2} \alpha_{1 \rightarrow 2}}{30 \hbar^3} \right)^{-1} T_2^{-3} \dots\dots (1.2.1.2-7)$$

$$\alpha_{1 \rightarrow 2} = \frac{\sum_j V_{2,j}^{-2}}{\sum_j V_{1,j}^{-2} + \sum_j V_{2,j}^{-2}} \dots\dots (1.2.1.2-8)$$

Comparing with the AMM model, the DMM model can better predict the interfacial thermal resistance under high-temperature conditions. For example, Stevens<sup>[93]</sup> and Stoner<sup>[94]</sup> reported that the value of interfacial thermal resistances calculated with the DMM model under room temperature was one order of magnitude less than the experimental value of interfacial thermal resistances. AMM and DMM models both have a specific temperature range to calculate interfacial thermal resistance. The reasons are that AMM ignores the phonons scattering at the interface, while DMM assumes that phonons are completely scattered at the interface.

The heat transfer mechanism of hot carriers (phonons and electrons) at the contact interface is very complicated. Therefore, it is difficult to obtain ideal results if only the

macro or micro model calculates the interface heat conduction. Fortunately, Lattice Boltzmann proposed a mesoscopic simulation at multi-scale levels. The model combines the discrete macroscopic simulation and the continuous microscopic method, and the calculation of interface thermal resistance considers the scattering and radiation of phonons and electrons. Han et al. [68, 70] calculated the interfacial thermal resistance using the thermal lattice Boltzmann method framework.

### 1.3.3 Theoretical calculation for the interfacial thermal resistance at GFs-Al interface

For GFs/Al composites, the interfacial thermal resistance is considered to have a considerable effect on the ETC, and the effect cannot be eliminated. The interfacial thermal resistance at the GFs-Al interface is commonly calculated using AMM and DMM models. Interfacial thermal resistance is equal to the reciprocal of interfacial thermal conductance ( $h$ ). Chen et al. [32] obtained the experimental value of  $h$  to be  $5.0 \times 10^7 \text{ W m}^{-2} \text{ K}^{-1}$ , and the calculated value by the DMM model showed the same order of magnitude as  $10^7$ . In addition, The values of  $h$  calculated by AMM were as  $4.5 \times 10^7$  [11, 29]  $\text{W m}^{-2} \text{ K}^{-1}$ ,  $4.8 \times 10^7$  [26]  $\text{W m}^{-2} \text{ K}^{-1}$ , and  $7.4 \times 10^7$  [13]  $\text{W m}^{-2} \text{ K}^{-1}$ .

However, as mentioned in 1.3.1 and 1.3.2 parts, the AMM and DMM models are suitable for calculating interfacial thermal resistance in different temperature ranges. Moreover, the calculation results are not completely reliable because the calculations are based on an ideal contact interface. Therefore, it is crucial to understand how heat conducts at the interfacial surface to obtain the accurate value of interfacial thermal resistance. Researchers considered that the hot carriers of metal heat transfer are mainly electrons at the micro-level, and the hot carriers of nonmetal heat transfer are mainly phonons. Hence, the heat transfer across the metal-nonmetal interface has two possible pathways. One is the coupling between electrons of metal and phonons of nonmetal through a harmonic interaction at the contact interfaces. Another is the coupling

between electrons and phonons in the metal matrix and then the couplings of the metal electrons and the nonmetal phonons<sup>[95]</sup>. As Battabyal et al.<sup>[96]</sup> reported that the heat conduction mechanisms at the diamond-Al interface were the coupling between Al electrons and Al phonons and the coupling between diamond and Al phonons that occurred at the interface. Similarly, the heat conduction mechanism at the diamond-Al should be suitable for the heat conduction at the GFs-Al interface.

## **1.4 Theoretical and simulation calculations for the ETC of composites**

### **1.4.1 Theoretical calculation for the ETC of composites**

The influential factors of the ETC of composite include the size, shape, volume fraction, dispersion and orientation of the filling phase, porosity, and interfacial thermal resistance. In past studies, many models were proposed to calculate the effects of those factors on the ETC of composites. The models<sup>[18, 48, 49, 57, 97-105]</sup> can be divided into theoretical and empirical models. Table 1-4 shows some summarized theoretical models.

Maxwell's model is the original theoretical model used to predict the ETC of composites<sup>[106, 107]</sup>. Maxwell's model assumes that the shape of the filled phase is spherical, the volume fraction of the filled phase is small, and the filled phases are uniform dispersion and independent of each other. However, Maxwell's model is a very simple theoretical model, and its application has significant limitations. For example, the ETC value calculated by Maxwell's model would significantly differ from the experimental ETC value if the filled phase has a high-volume fraction, or the shape of the filled phase was ellipsoid, fiber, and sheet, or considering the effects of interfacial thermal resistance and porosity. Therefore, many new theoretical models extended based on Maxwell's model were proposed, such as Bruggeman's model<sup>[58, 108]</sup>, Fricke's model<sup>[109]</sup>, Hamilton-Crosser's model<sup>[110]</sup>, Every's model<sup>[58]</sup>, Nan's model<sup>[111]</sup>, etc.

For example, Bruggeman's model was developed by integrating the differential equation of Maxwell's model, and it is suitable for the ETC calculations of composites with a high-volume fraction of the filled phase. Fricke's model is suitable for the ETC calculations of composites with filled elliptical phases and a high volume fraction of filled phase. Hamilton-Crosser's model is suitable for calculating the effect of the filled phase shape on the ETC. Furthermore, Every's model was extended based on Bruggeman's model, and it further considers the effect of the interfacial thermal resistance on ETC. Nan's model can consider the effects of size, shape, orientation,

dispersion, the volume fraction of filled phase, and interfacial thermal resistance on ETC.

Besides, many scholars also developed other theoretical models, such as the Cheng-Vachon model <sup>[112, 113]</sup>, Nielsen model <sup>[114, 115]</sup>, etc. Meanwhile, the empirical models have the Springer-Tsai model <sup>[116]</sup>, Y. Agari's models<sup>[107, 117, 118]</sup>, etc.

Table 1-4 Theoretical models using to calculate the ETC of composites.

Models	Equation
Maxwell	$\lambda_c = \lambda_m \frac{2\lambda_m + \lambda_p - 2V_p(\lambda_m - \lambda_p)}{2\lambda_m + \lambda_p + V_p(\lambda_m - \lambda_p)}$
Bruggeman	$(1 - V_p)^3 = \frac{\lambda_m}{\lambda_c} \left( \frac{\lambda_c - \lambda_p}{\lambda_m - \lambda_p} \right)^3$
Fricke	$\lambda_c = \lambda_m \left\{ \frac{1 + V_p [S(\lambda_p / \lambda_m - 1)]}{1 + V_p (S - 1)} \right\};$ $S = \frac{1}{3} \sum_{i=1}^3 [1 + (\lambda_p / \lambda_m - 1) f_i]^{-1};$ $\sum_{i=1}^3 f_i = 1; f_i \text{ is the half-axis length of elliptical particles.}$
Hamilton-Crosser	$\lambda_c = \lambda_m \left[ \frac{\lambda_p + (n-1)\lambda_m + (n-1)V_p(\lambda_p - \lambda_m)}{\lambda_p + (n-1)\lambda_m - V_p(\lambda_p - \lambda_m)} \right];$ $n = 3/\psi; \psi \text{ is Particle sphericity.}$
Every	$(1 - V_p)^3 = \left( \frac{\lambda_m}{\lambda_c} \right)^{(1+2a)/(1-a)} \left[ \frac{\lambda_c - \lambda_p(1-a)}{\lambda_m - \lambda_p(1-a)} \right]^{3/(1-a)}$ $a \text{ is the radius of the reinforcement particles.}$
Cheng-Vachon	$\frac{1}{\lambda_c} = \frac{1-B}{\lambda_m} +$ $\frac{1}{\{C(\lambda_p - \lambda_m)[\lambda_m + B(\lambda_p - \lambda_m)]\}^{1/2}} \ln \frac{[\lambda_m + B(\lambda_p - \lambda_m)]^{1/2} + B/2[C(\lambda_p - \lambda_m)]^{1/2}}{[\lambda_m + B(\lambda_p - \lambda_m)]^{1/2} - B/2[C(\lambda_p - \lambda_m)]^{1/2}};$ $B = \left( \frac{3V_p}{2} \right)^{1/2}, C = B^{-1};$
Nielsen	$\lambda_c = \lambda_m \frac{1 + (K_E - 1) \left[ \left( \frac{\lambda_p}{\lambda_m} - 1 \right) / \left( \frac{\lambda_p}{\lambda_m} + K_E - 1 \right) \right] V_p}{1 - \left[ \left( \frac{\lambda_p}{\lambda_m} - 1 \right) / \left( \frac{\lambda_p}{\lambda_m} + K_E - 1 \right) \right] \left[ 1 + \frac{(1 - V_m)V_p}{V_m^2} \right] V_p};$ $K_E \text{ is the Einstein coefficient.}$

where  $\lambda$  and  $V$  represent the TC and volume fraction; Subscript  $c$  represents composite, and  $m$  and  $p$  represent the matrix and filled phase of composites, respectively.

### 1.4.2 Theoretical calculation for the ETC of GFs/Al composites

As above summarized, many theoretical models have been developed to predict the ETC of composites. However, according to the structural characteristics of GFs/Al composites, the mixing rule and Nan's model are the favorite of researchers [14, 18, 26, 28, 29].

The mixing rules can be expressed:

$$\lambda_c^{x-y} = V_{GFs} \times \lambda_{GFs}^{x-y} + (1 - V_{GFs}) \times \lambda_{Al} \quad \dots\dots (1.3.2-1)$$

$$\frac{1}{\lambda_c^z} = \frac{V_{GFs}}{\lambda_{GFs}^z} + \frac{1 - V_{GFs}}{\lambda_{Al}} \quad \dots\dots (1.3.2-2)$$

where  $\lambda_c$  is the ETC of GFs/Al composites,  $\lambda_{GFs}^{x-y}$  is the TC of GFs along the basal-plane,  $\lambda_{GFs}^z$  is the TC of GFs along the out-plane,  $\lambda_{Al}$  is the TC of Al matrix, and the  $V_{GFs}$  is the volume fraction of GFs. Equation 1.3.2-1 uses to calculate the ETC of GFs/Al composites when the basal-plane of GFs is parallel to the heat flow direction; Equation 1.3.2-2 uses to calculate the ETC of GFs/Al composites when the out-plane of GFs is parallel to the heat flow direction. However, it is worth noting that the mixing rule is a simple theoretical model to calculate the ETC of composites, and the mixing rules only consider the effect of volume fraction of the filled phase on the ETC. Therefore, the ETC of GFs/Al composites calculated using the mixing rules would be higher than the ETC obtained by experiment methods, mainly attributed to the effects of the orientation and the dispersion of GFs and interfacial thermal resistance.

To consider the effect of interfacial thermal resistance, TC of GFs,  $\lambda_{GFs-eff}^i$ , can be calculated by the following equation 1.3.2-3, as follows:

$$\lambda_{GFs-eff}^i = \frac{\lambda_{GFs}^i}{1 + \frac{2\lambda_{GFs}^i}{hD}} \quad \dots\dots (1.3.2-3)$$

where  $i$  is the  $x$ - $y$  plane or  $z$ -axis direction,  $h$  is the interfacial thermal conductance, and

$D$  is the thickness and diameter of GFs.  $h$  is the reciprocal of interfacial thermal resistance ( $R$ ),  $h = 1/ R$ . To calculate ETC of composites considering the interfacial thermal resistance, the equations can be obtained using  $\lambda_{GFs-eff}$  to replace the  $\lambda_{GFs}^{x-y}$  and the  $\lambda_{GFs}^z$  in the mixing rules.

Nan's model<sup>[111]</sup> is an effective medium approximate model (EMA) to predict the ETC of GFs/Al composites. It includes the effects of the size, shape, and orientation<sup>[28, 29]</sup>. The equations can be expressed as:

$$\lambda_c^{x-y} = \lambda_{Al} \frac{2 + V_{GFs} [\lambda_{GFs}^{x-y} (1 - S_{11}) (1 + \langle \cos^2 \theta \rangle) + \lambda_{GFs}^z (1 - S_{33}) (1 - \langle \cos^2 \theta \rangle)]}{2 - V_{GFs} [\lambda_{GFs}^{x-y} S_{11} (1 + \langle \cos^2 \theta \rangle) + \lambda_{GFs}^z S_{33} (1 - \langle \cos^2 \theta \rangle)]} \dots\dots (1.3.2-4)$$

$$\lambda_c^z = \lambda_{Al} \frac{1 + V_{GFs} [\lambda_{GFs}^{x-y} (1 - S_{11}) (1 - \langle \cos^2 \theta \rangle) + \lambda_{GFs}^z (1 - S_{33}) \langle \cos^2 \theta \rangle]}{1 - V_{GFs} [\lambda_{GFs}^{x-y} S_{11} (1 - \langle \cos^2 \theta \rangle) + \lambda_{GFs}^z S_{33} \langle \cos^2 \theta \rangle]} \dots\dots (1.3.2-5)$$

$$\langle \cos^2 \theta \rangle = \frac{\int \rho(\theta) \cos^2 \theta \sin \theta d\theta}{\int \rho(\theta) \sin \theta d\theta} \dots\dots (1.3.2-6)$$

where  $\theta$  is the angle between heat flow direction and the basal plane of GFs;  $S_{ij}$  is the geometrical factors (such as size and shape) of GFs. To consider the effects of the geometrical factors  $S_{ij}$  of GFs on the ETC of GFs/Al composites, the thermal conductivities of GFs along the basal-plane ( $\lambda_{GFs}^{x-y}$ ) and the out-plane ( $\lambda_{GFs}^z$ ) are replaced by  $\lambda_{GFs-e}^{x-y}$  and  $\lambda_{GFs-e}^z$ .

$$\lambda_{GF}^{x-y} = \frac{\lambda_{GFs}^{x-y} - \lambda_{Al}}{\lambda_{Al} + S_{11} (\lambda_{GFs}^{x-y} - \lambda_{Al})} \dots\dots (1.3.2-7)$$



$$\lambda_{GFs-eff}^z = \frac{\lambda_{GFs}^z - \lambda_{Al}}{\lambda_{Al} + S_{33}(\lambda_{GFs}^z - \lambda_{Al})} \dots\dots (1.3.2-8)$$

Assuming that GFs is an ellipsoidal shape, the semi-axes  $a$  and  $b$  of GFs on the two-dimensional plane are the same ( $a = b$ ), and the thickness of GFs is  $c$ . According to the ratio ( $p$ ) between  $c$  and semi-axes  $a$  or  $b$ ,  $S_{ii}$  can be obtained using the following equations<sup>[119]</sup>:

$$S_{11} = S_{22} = \begin{cases} \frac{p^2}{2(p^2 - 1)} - \frac{p}{2(p^2 - 1)^{2/3}} \cosh^{-1}p, & p > 1 \\ \frac{p^2}{2(p^2 - 1)} + \frac{p}{2(p^2 - 1)^{2/3}} \cos^{-1}p, & p < 1 \end{cases} \dots\dots (1.3.2-9)$$

If GFs are highly oriented in the Al matrix, the  $\langle \cos^2\theta \rangle$  is infinitely close to 1. Moreover, to regard the shape of GFs as a disk,  $S_{ij}$  were calculated<sup>[119]</sup> as,  $S_{11} = S_{22} = \pi p/4$  and  $S_{33} = 1 - \pi p/2$ .

### 1.4.3 Simulation calculation for the ETC of composites

With the rapid development of computer technology, simulation methods have become a new and effective research tool to calculate the ETC of composites. Many simulation methods have been developed to predict the ETC of composites, such as molecular dynamics (MD), lattice Boltzmann method (LBM), finite volume method (FVM), finite element method (FEM), etc. Those methods analyze the ETC of composites at different simulation scales from micro-scale, mesoscale, to macro-scale.

#### 1.4.3.1 Microscale model

The microscopic simulation is an atomic level calculation, and it is used to analyzes the interaction between atoms through first principles. The calculation method's essence is to use the Monte Carlo method (MC) or MD to solve the approximate solution of the Schrödinger equation. MD is the most used model to solve the Schrödinger equation. MD includes equilibrium and non-equilibrium molecular dynamics (EMD and NEMD).

EMD calculates the ETC of composites under the steady temperature field, and its theoretical function to calculate the ETC( $\lambda_c$ ) is based on the Green-Kubo equation, as follow:

$$\lambda_c = \frac{1}{3VK_B T^2} \int_0^{\infty} \langle J(t)J(0) \rangle dt \quad \dots\dots (1.3.3.1)$$

where  $V$  is the volume of the system,  $K_B$  is Boltzmann constant,  $T$  is temperature,  $J$  is microscopic heat flow, and  $t$  is time. However, using EMD to calculate ETC would generate amounts of calculation and spend lots of time because of the very slow convergence calculation of  $J$  and its related functions. As a comparison with EMD, NEMD shows a better convergence calculation. NEMD establishes a non-equilibrium heat conduction process by applying disturbances to the system, and then to calculate ETC of composites using the Fourier function of heat conduction,  $J = -\lambda \Delta T$ . NEMD method is generally used to calculate the ETC of composites<sup>[120-124]</sup>. Tian et al<sup>[120]</sup>. studied the effects of the filled phase's orientation, dispersion, and arrangement, the interface mismatch, the interface density, and the filled phase polydispersity on the ETC of nanocomposites. Yang et al. <sup>[124]</sup> employed the NEMD method to calculate the TC of defective graphene oxide. However, NEMD also has the disadvantages, such as poor model homogeneity, momentum, energy conservation, etc. Müller-plathe et al<sup>[125]</sup>. proposed a reverse perturbation non-equilibrium molecular dynamic method (RNEMD) for calculating the thermal conductivity. RNEMD can correct the disadvantages in the momentum and energy conservation of NEMD method, making it more suitable for calculating ETCs.

#### 1.4.3.2 Mesoscale model

Compared with the microscopic simulation, mesoscopic simulation models contain huge numbers of atoms (about  $10^{23}$  /cm<sup>3</sup>), and they are much larger than the microscopic simulation models. The theoretical equation is different from the

Schrödinger equation, and it belongs to the average basic characteristic structure relationship. Due to the complex diversity of mesoscopic scale mechanisms and constitutive laws, different mesoscopic simulation models were developed<sup>[126]</sup>: 1, Space and time discretization dislocation dynamics model; 2, Phase-field dynamics or generalized Ginzburg-Landau model; 3, Deterministic or probabilistic meta-packet automata model; 4, Polymorphic Dynamics Potts Model; 5, Geometry topology and grouping model; 6, Topological mesh and vertex model. Although the mesoscopic simulation models are diverse, they have some common characteristics. E.g., the mesoscopic model does not contain obvious atomic-scale dynamics but an idealized assumption that the material is a continuum. So, the lattice defects have a common uniform matrix, and the matrix naturally couples the interactions between lattice defects. Thus, intrinsic space and time scales are not explicitly included in the corresponding governing equations. Moreover, a continuum mesoscopic scale model containing a single lattice defect usually has a series of phenomenological partial differential rates and intrinsic structural equations. Those differential equations are usually solved by the finite difference method, finite element, or Monte Carlo method.

In the mesoscopic scale model, the determinations of space-time scale parameters and dispersion are determined by the differential equation and its coefficients and the characteristics and properties of the variables. The control of differential equations is used for local or global situations according to the interaction between lattice defects. State variables (atomic concentration, dislocation density, structural parameters, displacement, lattice orientation) are usually combined with spatial grid coordinates. Because the phenomenological equation and the law of structural evolution have been well studied on the mesoscopic scale, it is of great significance to use the continuum approximation method to predict the structural evolution on the mesoscopic scale. Additionally, the experimental data obtained from the mesoscale model is easier than the data obtained from the micro-scale, and it is more detailed than the information obtained from the macro-scale model. The methods to solve the mesoscopic scale model mainly include DM, MC, and finite element method (FEM).

Mesosopic scale models used to calculate the ETC of composites have been

developed, and the most commonly used calculation methods include dissipative particle dynamics (DPD), LBM, off-lattice MC, etc. Such as DPD was proposed by Hoogerbrugge and Koelman in the 1992 year, and it was used to solve the difference between the lattice automata method and the reality and the fluid problem on the mesoscopic time and space scale that MD cannot solve. DPD is a coarse-grained technology based on MD, which means that each particle represents one or more molecules or a fluid region of a fragment of a polymer rather than a single atom. Moreover, the particle interaction is expressed by a pair of conservative, dissipation, and random forces. R. Qiao et al. [127] calculated the TC of nanocomposites by energy-conserving DPD, and the obtained results were well agreed with the results calculated by the Maxwell-Garnett model. Bing Zhou et al. [128] studied the effect of volume fraction, alignment, dispersion and size of CNT, and polymer size on the ETC of CNT/polymer composites with random and aligned CNT by using coupling smoothed particle hydrodynamics and DPD methods.

#### 1.4.3.3 Macroscale and multiscale model

The macroscopic simulation model treats the composite material as a single orthotropic material or completely anisotropic material. The macroscopic models of composites are commonly used to simulate the overall performances of the composites. Thus, the Non-linearity of materials and local failures are not considered. Furthermore, compared with microscopic and mesoscopic scale models, macroscopic models do not contain the calculations at the atomic and molecular level, which can save lots of calculations and calculation time. Using macroscopic models to calculate the ETC of composites, FEM<sup>[105, 129-133]</sup>, FVM<sup>[49]</sup>, and FDM<sup>[134]</sup> are the most used methods.

Multi-scale simulation is a combination of microsimulation and macro simulation. Macroscopic simulation has the advantage of being a small amount of calculation and saving time, but it is difficult to reflect the effect of the microscopic behavior of each material on the overall properties of the composite. On the other hand, although the

microscopic simulation can well understand the influence of the micro-behavior of each material on the overall behavior of the composite material, it has a considerable amount of calculation. As a result, because of the limitation of computer computing power and the structural complexity of composites, it is only possible to establish a micro-model based on the local characteristics of the composite. In contrast, the multi-scale model can well coordinate the collaborative calculation of the micro model and the macro model to study better the influence of the micro behavior of the material on the overall behavior and the overall properties of the material. Based on those advantages, multi-scale simulation has been widely used in the ETC calculation of composites<sup>[135-140]</sup>.

## **1.5 Objective of this thesis**

GFs are low-cost, lightweight, and possess excellent thermal properties, including high thermal conductivity and low thermal expansion coefficient. GFs/Al composites are thus chosen as the candidates of thermal management materials. Many experimental and theoretical studies have been conducted to investigate the ETC of GFs/Al composites. However, previous experimental studies found that the orientation of GFs, the interface reaction, and the interfacial thermal resistance greatly affect the composites' thermal properties.

Theoretical calculations and three-dimensional(3D) simulations are used to evaluate the ETC of GFs/Al composites considering the orientation of GFs and the interfacial thermal resistance. However, both theoretical calculations and 3D simulations exhibit the disadvantages of calculating the ETC of GFs/Al composites. For theoretical calculations, it is difficult for those theoretical models to calculate the ETC of GFs/Al composites due to the diversity of GFs' orientations in composites. Additionally, the 3D simulation calculations would take lots of time and be very expensive because these 3D models are constructed using images from ultra-high-resolution X-ray computed tomography. On the other hand, although two-dimensional (2D) images simulations are low-cost and high timeliness, few studies have used 2D images simulations to calculate the ETC of GFs/Al composites. The reason can be attributed to the orientation of GFs in 2D images may not be equivalent to the orientation of GFs in the corresponding experimental samples, which can result in the ETC of GFs/Al composites calculated by 2D images simulations are not reliable.

This study aims to study the effects of the orientation of GFs and the interfacial thermal resistance on the ETC of GFs/Al composites using 2D images simulations. The following steps did the works: In the first step, GFs/Al composites were fabricated using SPS. We measured the relative density of GFs/Al composites, analyzed the distribution and orientation of GF, investigated the formation of  $Al_4C_3$ , and evaluated the ETC and thermal expansion using the experimental and theoretical methods.

Secondly, we studied the relationship between the TCs of the GFs in 3D and 2D images. As a result, we obtained a fitted function used to convert TC of GFs from 2D to 3D. ETC of 10%vol GFs/Al composite was calculated using 2D image simulations, employing the fitted function. Thirdly, the effect degrees of the orientation of GFs and the interfacial thermal resistance on the ETC of GFs/Al composites were calculated by 2D image simulations.

## **1.6 Outline of this thesis**

### **Chapter 1 Background and Objective**

This chapter discusses the development of electronic packaging materials, the fabrication methods of the GFs/Al composites, some experimental studies of the thermal properties of GFs/Al composites, some theoretical and simulation methods for calculating the ETC of GFs/Al composites, and the objectives of this thesis.

### **Chapter 2 Experiment process and Fundamental properties of GFs/Al composites**

In this chapter, the process of preparing GFs/Al composites using SPS was showed. The orientation, distribution, and aspect ratio of GFs in the Al matrix were analyzed by OM images. The relative density of composite was determined by the Archimedes method. The interface reaction was checked by XRD method. The thermal expansion coefficient and the ETC of GFs/Al composite were obtained by theoretical and experimental methods.

### **Chapter 3 Investigation the relationships between the effective thermal conductivity of GFs in the three-dimensional and two-dimensional model**

In this chapter, we constructed 3D and 2D models of GFs/Al composites. The orientations and TCs of GFs in the 3D and 2D models were compared. Furthermore, the relationships between the orientations of GFs in 3D and 2D models were analyzed. As a result, a fitted equation was derived and used to convert the TC of GFs from 2D to 3D. Moreover, the effective thermal conductivity of the graphite flake/aluminum composite with 10 vol% graphite flakes was calculated using 2D image simulation and the fitted equation.

### **Chapter 4 Effect of the orientation of GFs and interfacial thermal resistance on the effective thermal conductivity of GFs/Al composites**

In this chapter, the ETC of GFs/Al composites were calculated by using 2D image



simulations, including the effect of GFs' orientation and interfacial thermal resistance. The effect degree of the orientation of GFs and the interfacial thermal resistance on the ETC of GFs/Al composites were analyzed. The interfacial thermal resistances at the GFs-Al interface were calculated.

## **Chapter 5 Conclusions**

The results from the studies mentioned above were summarized in this chapter.

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## *Chapter 2*

### **Fabrication and Thermal Properties of GFs/Al composites**

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## **2.1 Introduction**

Carbon materials reinforced Al-matrix composites attract great interest from researchers of electron packaging materials due to their lightweight, superior mechanical, and thermal properties [1-7]. Especially for graphite flakes (GFs) reinforced Al-matrix composites, their excellent thermal properties, such as high thermal conductivity (TC) and low coefficient of thermal expansion [5, 8-10], are able to meet the thermal performance requirements of electron packaging materials.

Lots of experiments have been performed to study the effective thermal conductivity (ETC) of GFs/Al composites and found that the high thermal conductivity of GFs cannot be fully displayed due to the following characteristics: 1, the thermal conductivity of GFs is significant anisotropic, its orientation in the composites hinder the effective thermal conductivity of GFs/Al composites; 2, GFs tends to react with Al during the fabricating the GFs/Al composites, the formation of  $Al_4C_3$  is bad for the mechanical properties and reduce the effective thermal conductivity of GFs/Al composites; 3, the interfacial thermal resistance is considered to have a great effect on the effective thermal conductivity of GFs/Al composites. Therefore, a suitable fabrication process of GFs/Al composites is essential to take advantage of the high thermal conductivity properties of GFs.

In previous studies on the effective thermal conductivity of GFs/Al composite, the commonly used fabrication methods of GFs/Al composite include the vacuum hot pressing<sup>[9-12]</sup>, squeeze casting<sup>[13-15]</sup>, and gas pressure infiltration<sup>[8, 16]</sup>. The interface reaction can be effectively avoided using those methods. For example, Chen et al. [12] fabricated the GFs/Al composites by hot pressure method, and a low sintering temperature, lower than the melting point of Al, was used to avoid the interface reaction. The ETC of GFs/Al composites with 80 vol. % GFs achieved up to  $783 \text{ W m}^{-1} \text{ K}^{-1}$ . Xue et al.<sup>[10]</sup> fabricated GFs/Al composites by vacuum hot pressing method, the surface of GFs were pre-processed with SiC particles coating to inhibit the interface reaction. The ETC of GFs/Al composites with 70 vol. % GFs was up to  $735 \text{ W m}^{-1} \text{ K}^{-1}$ . Li et al.<sup>[13]</sup>



prepared GFs/Al composites by squeeze casting method, and rapid cooling method was used to hinder the interface reaction. The effective thermal conductivity of the sample with 70 vol. % GFs was as high as  $714 \text{ W m}^{-1} \text{ K}^{-1}$ . Zhou et al.<sup>[8]</sup> used pressure infiltration to prepare GFs/Al composites. Si particles were introduced to suppress the interface reaction. The effective thermal conductivity of the sample with 71.1 vol. % GFs was  $526 \text{ W m}^{-1} \text{ K}^{-1}$ . The interface reaction can be hinted at by low sintering temperature, rapid cooling, or adding Si or SiC particles. However, the effective thermal conductivities of GFs/Al composites were lower than the expected theoretical effective thermal conductivities. The main reasons can be attributed to the effects of the orientation of GFs and the interfacial thermal resistance between Al and GFs.

This study employed the spark plasma sintering (SPS) method to prepare GFs/Al composites. SPS is a low-temperature sintering technology. The low sintering temperature can avoid the interface reaction, and SPS can generate a plasma atmosphere to effectively destroy the oxide film on the surface of Al powders, increase new metallic contacts and necks, and reduce the porosity<sup>[17-20]</sup>. The GFs/Al composites with 10-40 vol. % of GFs were prepared. Orientations of GFs in GFs/Al composites were analyzed using the OM images; The interface reaction was detected by the XRD method; The effective thermal conductivity and the coefficient of thermal expansion of GFs/Al composites were evaluated by experimental and theoretical methods.

## 2.2 Experimental procedure

### 2.2.1 Raw materials

The raw materials consisted of Al powders and GFs. The institute of high purity chemistry offered Al powders in Japan, and Ito Graphite Industry Co., Ltd. offered the highly crystalline GFs in Japan. Figure 2-1 shows the microstructure of Al powder and GFs. Al powders show a rugby shape, and their size is uneven, while GFs show a flat surface and uniform thickness. Table 2-1 shows the chemical components of the raw materials. The purity is as high as 99.9% for Al powder and 98.8% for GFs. GFs are highly crystalline, which can hinder the chemical reaction between Al and GFs. Table 2-2 shows some physical properties and thermal conductivity of Al powder and GFs.

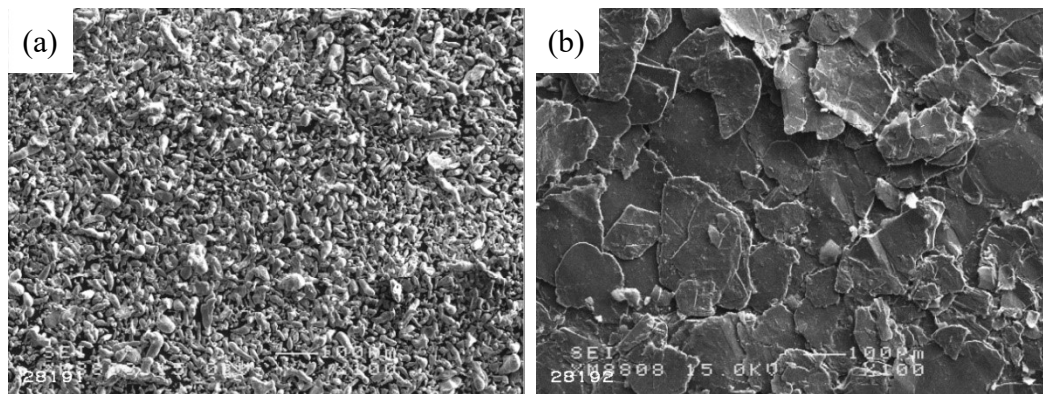


Figure 2-1 Microstructures of raw materials: (a) Al powder and (b) GFs.

Table 2 -1 Chemical components of Al and GFs.

Material	Carbon	Ash	Volatile	H <sub>2</sub> O	Al	Cu	Fe	Si
	%							
Al					99.9	0.001	0.06	0.03
GFs (XD150)	98.78	0.61	0.61	0.27				

Table 2-2 Shape and physical properties of Al and GFs.

Materials	Shape	$d$	$\rho$	$C_p$	TC
		$\mu\text{m}$	$\text{g}\cdot\text{cm}^{-3}$	$\text{J}\cdot\text{kg}^{-1}\cdot\text{K}^{-1}$	$\text{W}\cdot\text{m}^{-1}\cdot\text{K}^{-1}$
Al	powder	30	2.7	880	236
GFs	flake	137.02	2.23	710	880 (in-plane) 38 (out-plane)

where  $d$  is the average diameter,  $\rho$  is density, and  $C_p$  is the specific heat.

### 2.2.2 Preparation of graphite flakes and Al powders mixtures

The wet mixing process was used to mix Al powders and GFs. First, the Al powders, GFs, and alumina balls were filled into an aluminum jar. Subsequently, the mixtures were mixed for 2 h using a V-type mixer at 50rpm. Finally, the mixtures were dried in a constant temperature oven at 70 °C. The mixing procedure of raw material Al powder and GFs was shown in Figure 2-2.

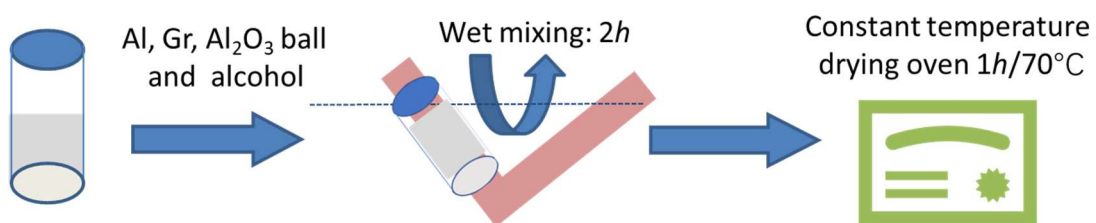


Figure 2-2 Flowchart of GFs and Al powder mixing.

According to Xu's studies <sup>[21]</sup>, the morphology of the Al particle would not be destroyed when the diameter of the Alumina ball is 10 mm. Thus, the Alumina ball with 10 mm diameter was used in the mixing process of Al powder and GFs. Besides, the choice of dispersion solution is crucial for the uniform mixing of GFs and Al powder. Because the density of GFs is lower than that of Al, the low-viscosity dispersion solution may cause Al particles and GFs to separate during the mixing process; Al particles settle at the bottom of the mixed solution, while GFs are in an aggregate state. Commonly used hybrid solutions are ethanol, butanol, and acetone, and their viscosity

and density are shown in Table 2-3. The sedimentation velocity ( $V$ ) of Al powders with  $30\mu\text{m}$  diameter in ethanol( $V_e$ ), butanol ( $V_b$ ), and acetone ( $V_a$ ) can be evaluated by the following equation [22]:

$$V = d^2 \cdot g \cdot (\rho_{Al} - \rho_s) / 18\eta \quad \dots\dots (2-1)$$

where the  $d$  is the diameter of Al particle,  $g$  is gravity,  $\rho_{Al}$  and  $\rho_s$  are the density of the Al powders and dispersion solution, and  $\eta$  is the viscosity of the dispersion solution.

Table 2-3 Viscosity and density of ethanol, butanol and acetone.

Dispersion solution	Butanol	ethanol	Acetone
Viscosity ( $\eta$ , $10^{-3}$ Pa·s)	2.95	1.07	0.32
Density ( $\rho$ , $10^6$ g/m <sup>3</sup> )	0.81	0.80	0.79

Figure 2-3 shows the calculated result. The sedimentation velocity is  $V_b < V_e < V_a$ . Therefore, butanol should be the most suitable dispersion solution in the wet mixing process. However, the higher viscosity of butanol can cause a lower drying speed of the mixed solution of GFs and Al powders, and there may be residual butanol in the mixture[21]. Moreover, the difference between  $V_b$  and  $V_e$  is not large (as shown in Figure 2-3). Thus, ethanol is used as the dispersion solution in the mixing process of GFs and Al powders.

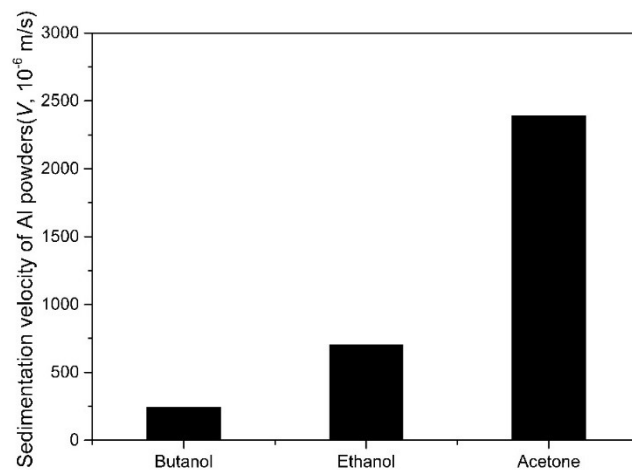


Figure 2-3 Sedimentation velocity of Al powders with  $30\mu\text{m}$  diameter in ethanol, butanol, and acetone [21].

### 2.2.3 Fabrication of graphite flakes/Al composites by sparking plasma sintering

Figure 2-4 shows the equipment of SPS. The mixed powders were filled into a graphite mold (graphite die) to fabricate GFs/Al composites. The mold was tapped during the filling of mixed powders in order to stack GFs in layers. Then, spark plasma sintering (SPS) was conducted at 873 K, and 60 MPa pressure for 0.5 h under a vacuum of  $1.3 \times 10^{-2}$  Pa. Sintering temperature was kept below the melting point of Al to prevent the reaction between Al and GF. The temperature was increased at a rate of 200 K/min. The details of the sintering procedure were shown in Figure 2-5. A pure Al sample was fabricated by SPS under the same condition as those composites. The dimensions of the sintered GFs/Al composites were  $65 \times 10 \times 15 \text{ mm}^3$ .

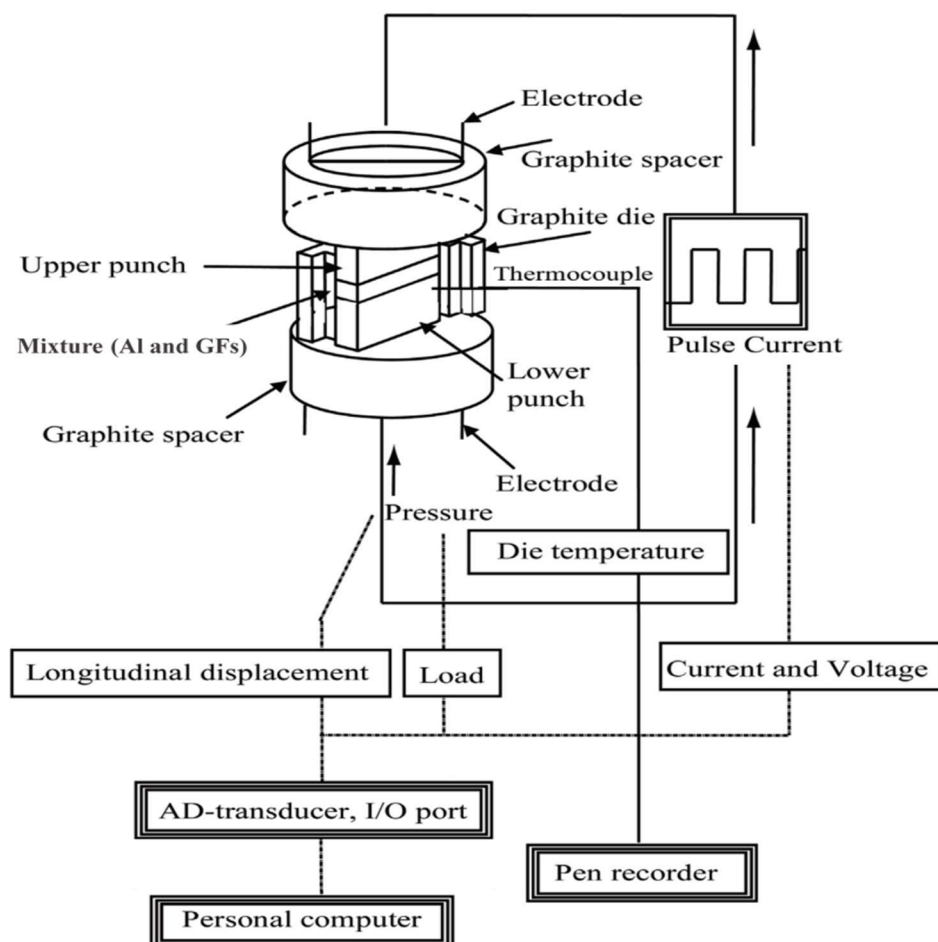


Figure 2-4 Schematic of the SPS equipment.

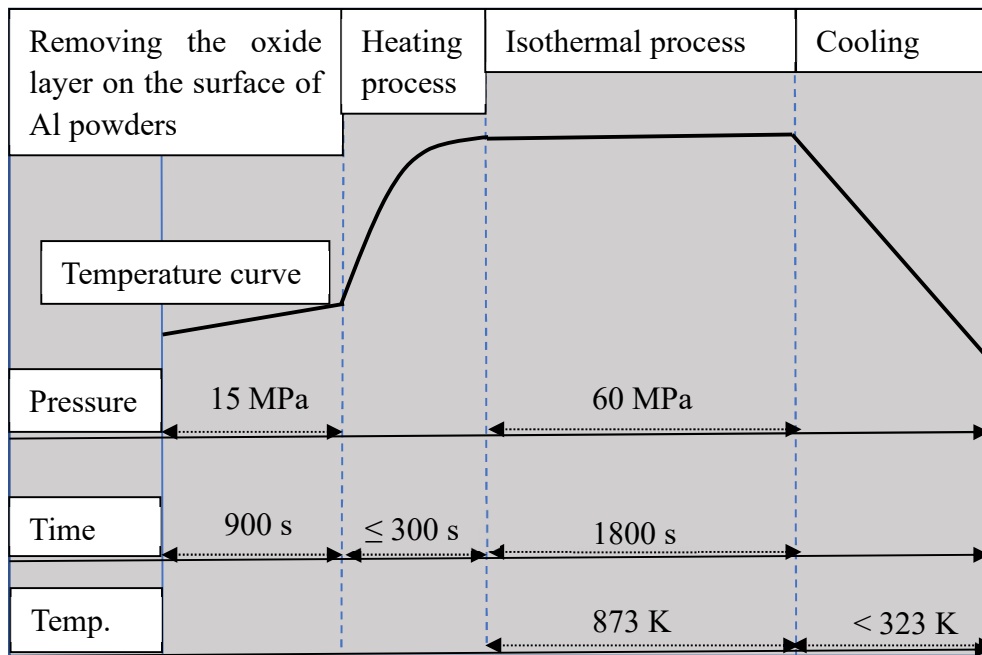


Figure 2-5 Flow chart of spark sintering process.

## 2.3 Evaluation methods

### 2.3.1 Microstructure and relative density

The electron probe micro-analyzer (EPMA) and an optical microscope (OM) were used to observe the microstructure of the composites and the mixtures of Al and GFs powders; the relative densities of composites were measured using the Archimedes method.

### 2.3.2 X-ray powder diffraction method

The formation of  $Al_4C_3$  was investigated by the X-ray powder diffraction method (XRD) (D/MAX-2500/PC). The measurement conditions are listed in Table 2-4.

Table 2-4 Measurement conditions of XRD.

Target	Cu
Wavelength ( $10^{-10}$ m)	1.54056 ( $K_{\alpha 1}$ )
Power (kW)	$\leq 18$
Voltage (kV)	40
Current (mA)	200
Sampling step (degree)	0.02
Scan speed (degree/min)	1

### 2.3.3 Effective thermal conductivity

The effective thermal conductivities of GFs/Al composites were measured at 25 °C using a steady-state thermal-conductivity measuring device<sup>[23]</sup>, as shown in Figure 2-5. The device consists of a heating end and a cooling end. The heating end includes a heating plate and a hot bar. The heating plate is connected to the cartridge heaters, and a heater controller controls the temperature ( $T_H$ ). The cooling end also includes two components: a cooling plate and a cold bar. The cooling plate is connected to the water-

cooled chiller through the cooling pipes. The water-cooled chiller controls the temperature of the cooling plate ( $T_c$ ). There are three thermocouples fixed on the hot bar and the cold bar, respectively. The hot bar and cool bar are cylindrical copper rods with a diameter of 10 mm. When measuring the effective thermal conductivity, the sample is sandwiched between the hot bar and cold bar and contacts three thermocouples. After the temperature of the nine thermocouples reaches a steady state, the temperatures for no less than 300 seconds at a 1-second interval are collected by the data acquisition system. Subsequently, the temperature gradients of the hot bar, cold bar, and sample ( $\Delta T_h$ ,  $\Delta T_s$ ,  $\Delta T_c$ ) are calculated by a least-square method as follows<sup>[23]</sup>:

$$\Delta T_h = \frac{\sum_{i=1}^3 (Z_i - \bar{Z}_h)(T_i - \bar{T}_h)}{\sum_{i=1}^3 (Z_i - \bar{Z}_h)^2}, \quad \bar{Z}_h = \frac{1}{3} \sum_{i=1}^3 Z_i, \quad \bar{T}_h = \frac{1}{3} \sum_{i=1}^3 T_i$$

..... (2-2)

$$\Delta T_s = \frac{\sum_{i=4}^6 (Z_i - \bar{Z}_s)(T_i - \bar{T}_s)}{\sum_{i=4}^6 (Z_i - \bar{Z}_s)^2}, \quad \bar{Z}_s = \frac{1}{3} \sum_{i=4}^6 Z_i, \quad \bar{T}_s = \frac{1}{3} \sum_{i=4}^6 T_i$$

..... (2-3)

$$\Delta T_c = \frac{\sum_{i=7}^9 (Z_i - \bar{Z}_c)(T_i - \bar{T}_c)}{\sum_{i=7}^9 (Z_i - \bar{Z}_c)^2}, \quad \bar{Z}_c = \frac{1}{3} \sum_{i=7}^9 Z_i, \quad \bar{T}_c = \frac{1}{3} \sum_{i=7}^9 T_i$$

..... (2-4)

where  $Z_i$  is the distance between the  $i$ -th thermocouple and the heating plate, and  $T_i$  is the temperature of the  $i$ -th thermocouple. Finally, the effective thermal conductivity of the sample ( $\lambda_{eff}$ ) can be calculated by following equation:

$$\lambda_{eff} = \lambda_{Cu} C (\Delta T_h + \Delta T_c) / 2 \Delta T_s$$

..... (2-5)

where  $\lambda_{Cu}$  is the thermal conductivity of copper bar ( $\lambda_{Cu} = 385 \text{ Wm}^{-1}\text{K}^{-1}$ );  $C$  is a correction factor, and it is a constant value (i.e.,  $C = 236/257$ ).



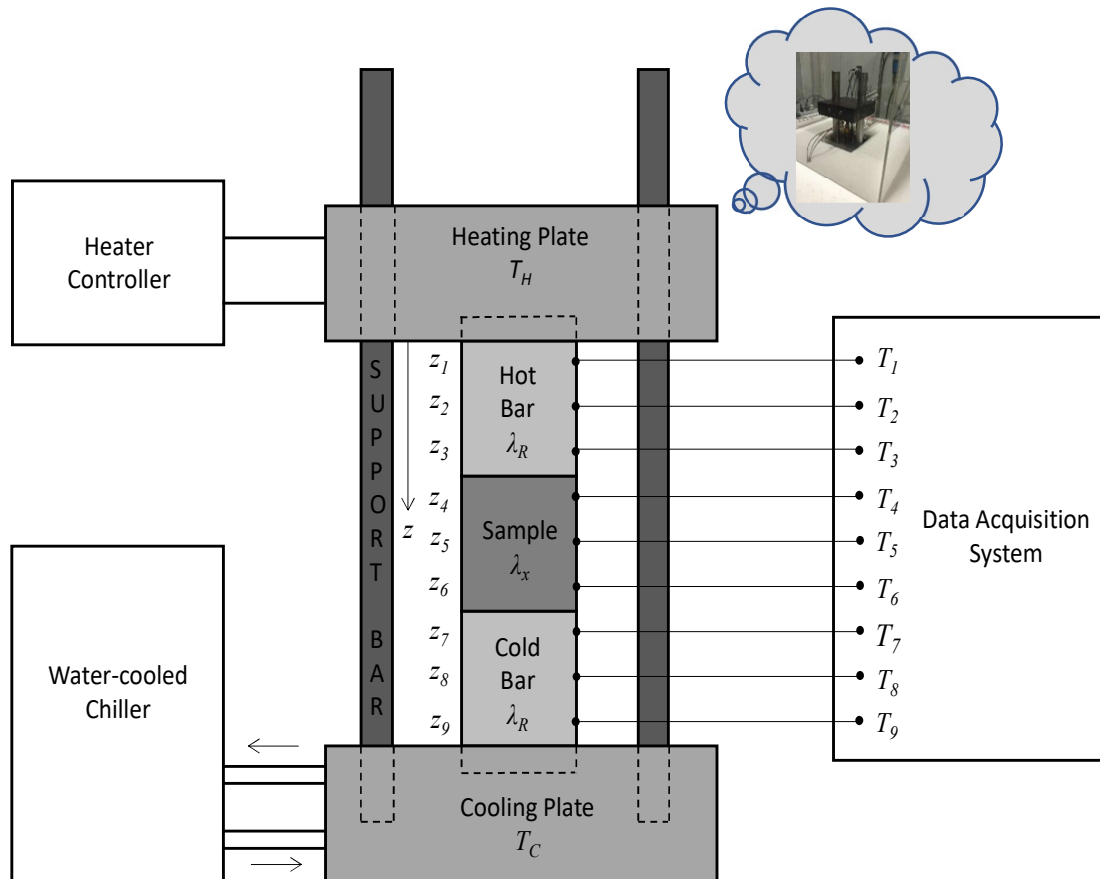


Figure 2-5 Schematic diagram of the steady-state thermal conductivity measurement device [23].

### 2.3.4 Coefficient of thermal expansion

The heat expansion of GFs/Al composites was measured from 25 to 200 °C by a thermal expansion meter (NETZSCH DIL 402 C).

## 2.4 Results and discussion

### 2.4.1 Microstructure characterization and orientation of GFs in composites

Figure 2-6 shows the microstructure images of the mixtures of Al and 10-40 vol.% GFs. GFs were evenly dispersed in the Al powders, and GFs in Figure 2-6 (a-d) have the same particle size. Moreover, compared with Figure 2-1, GFs did not show the obvious fracture, its average particle size was still about 130 $\mu$ m, and the morphology of Al powder was no obvious change.

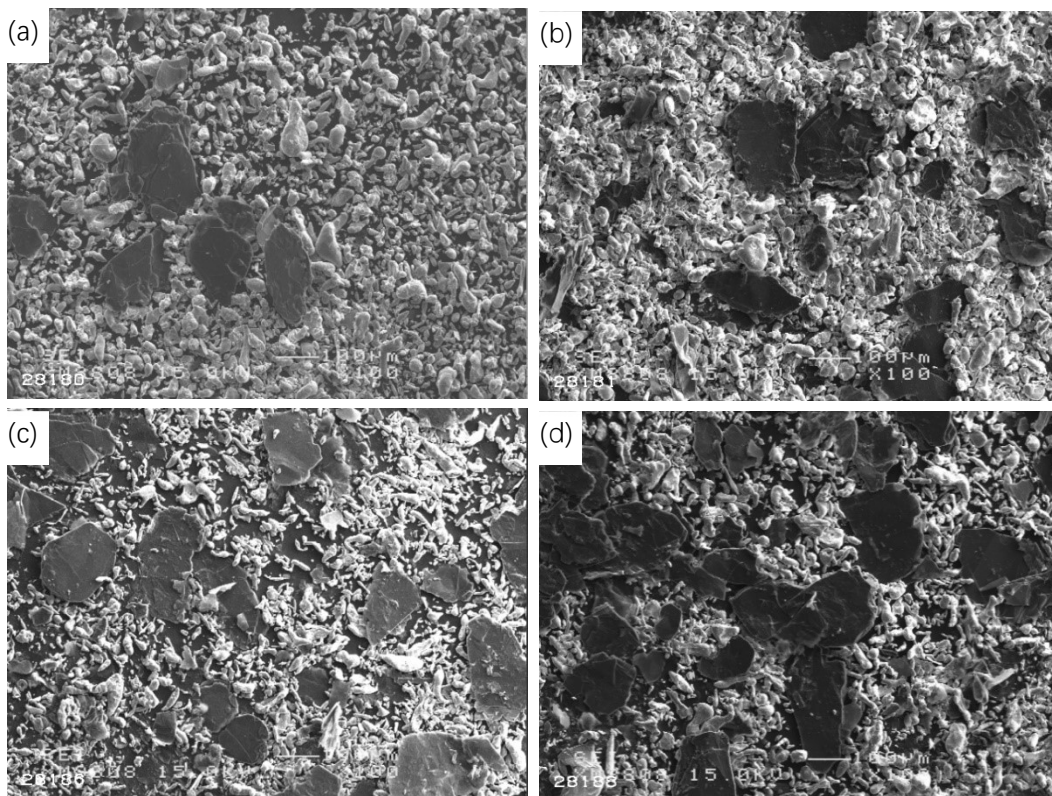


Figure 2-6 SEM images of the mixtures of Al and 10-40 vol. % GFs. 10 vol.% GFs (a), 20 vol.% GFs (b), 30 vol.% GFs (c), and 40 vol.% GFs (d).

10-40 vol. % GFs/Al composites and one pure Al sample were prepared by SPS. Their relative densities were over 98%, except for the pure Al sample being at 97%. The high relative density indicated that the fabricated GFs/Al composites were densified. The detailed values of the relative densities were listed in Table 2-5.

Table 2-5 Relative density and effective thermal conductivity (ETC) of GFs/Al composites.

Sample	Al		GFs/Al composites		
	0	10 vol.%	20 vol.%	30 vol.%	40 vol.%
Volume fraction of GFs	0	10 vol.%	20 vol.%	30 vol.%	40 vol.%
Relative density	97.0%	99.4%	99.3%	98.6%	97.1%
ETC ( $\lambda_{eff}$ , $\text{Wm}^{-1}\text{K}^{-1}$ )	217	248	280	313	346

Figure 2-7 depicts the shape of the samples sintered via SPS. The black arrows represent that C-plane was bearing the pressure during sintering. Therefore, GFs in the Al-matrix were stacked by layers, and the basal plane of GFs is parallel to the C surface. In order to better observe the orientation of GFs in the Al-matrix, the OM images of the GFs/Al composites were obtained from the A surface.

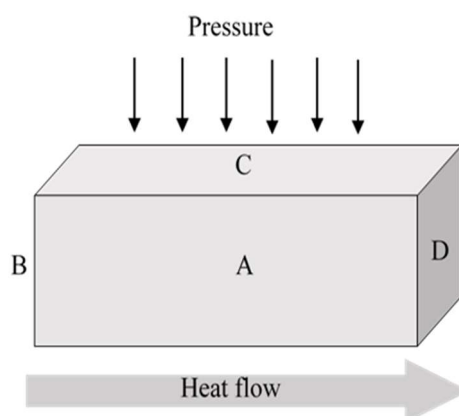


Figure 2-7 Schematic of the GFs/Al composites sintered by SPS. The grey arrow indicates that the heat flow was from B to D when measuring the ETC of the sample.

Figure 2-8 shows the OM images of 10-40 vol.% GFs/Al composites. The grey phases are of the Al matrix, and the elongated dark phases are GFs. The GFs were homogeneously distributed in the Al matrix and stacked by layers as expected. Most GFs were parallel to the heat flow direction, except for a few GFs whose basal plane direction formed an angle with the direction of the heat flow. In addition, as shown in Figure 2-8, with the increase of GFs content from 10 vol.% to 40 vol.%, GFs exhibits aggregation along the direction perpendicular to the heat flow direction, and GFs tend to form a network system by connecting GFs. As Chen et al. [12] reported, networks

formed by connecting GFs can effectively improve the ETC of composites.

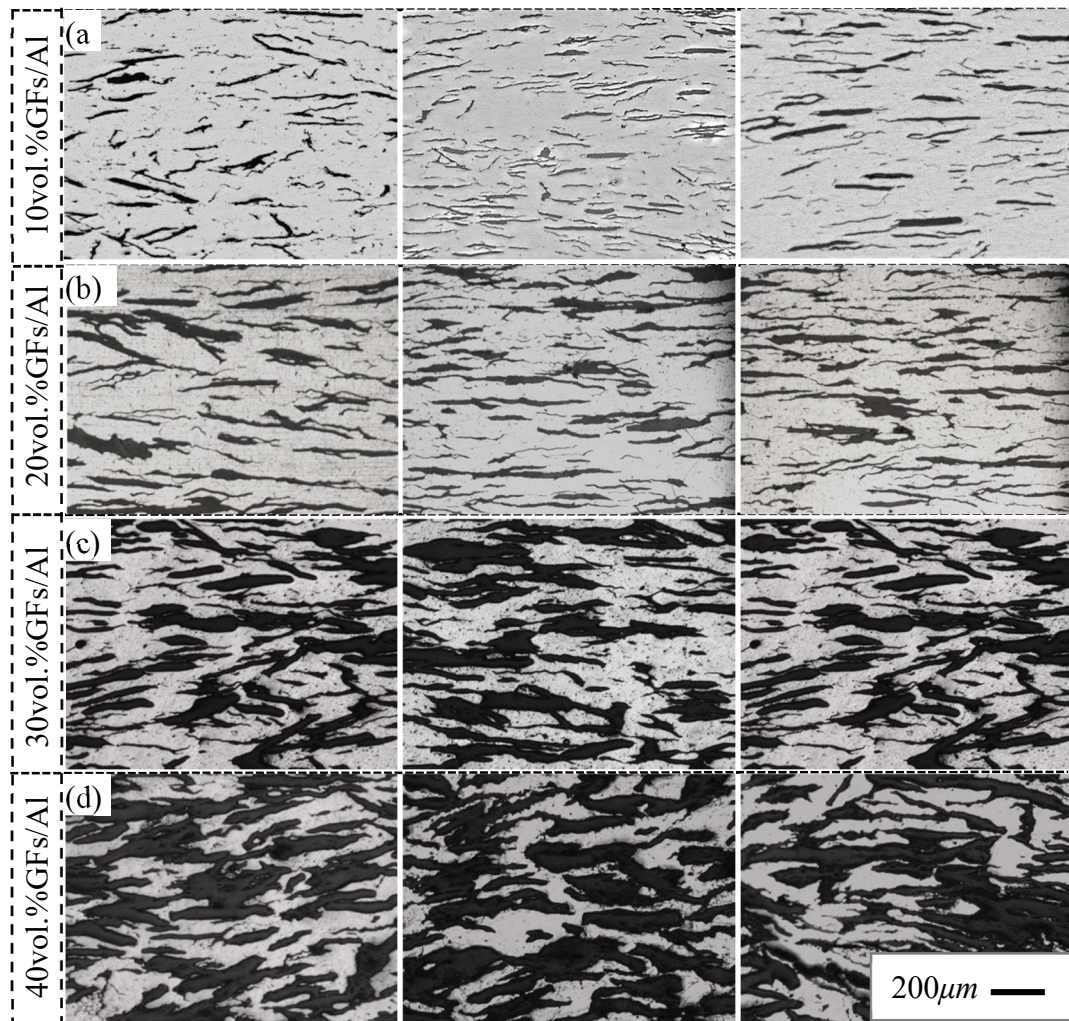


Figure 2-8 OM images of 10-40 vol.% GFs/Al composites.

An angle between GFs basal plane and heat flow direction was used to characterize the orientation of GFs in composites. The angles were counted statistically, and the results were shown in Figure 2-9. The  $x$ -axis showed that almost 70% of the angles were less than 10 degrees, which meant that nearly 70% of the GFs were oriented roughly parallel to the direction of heat flow. GFs parallel to each other is expected to enhance the overall thermal conduction.

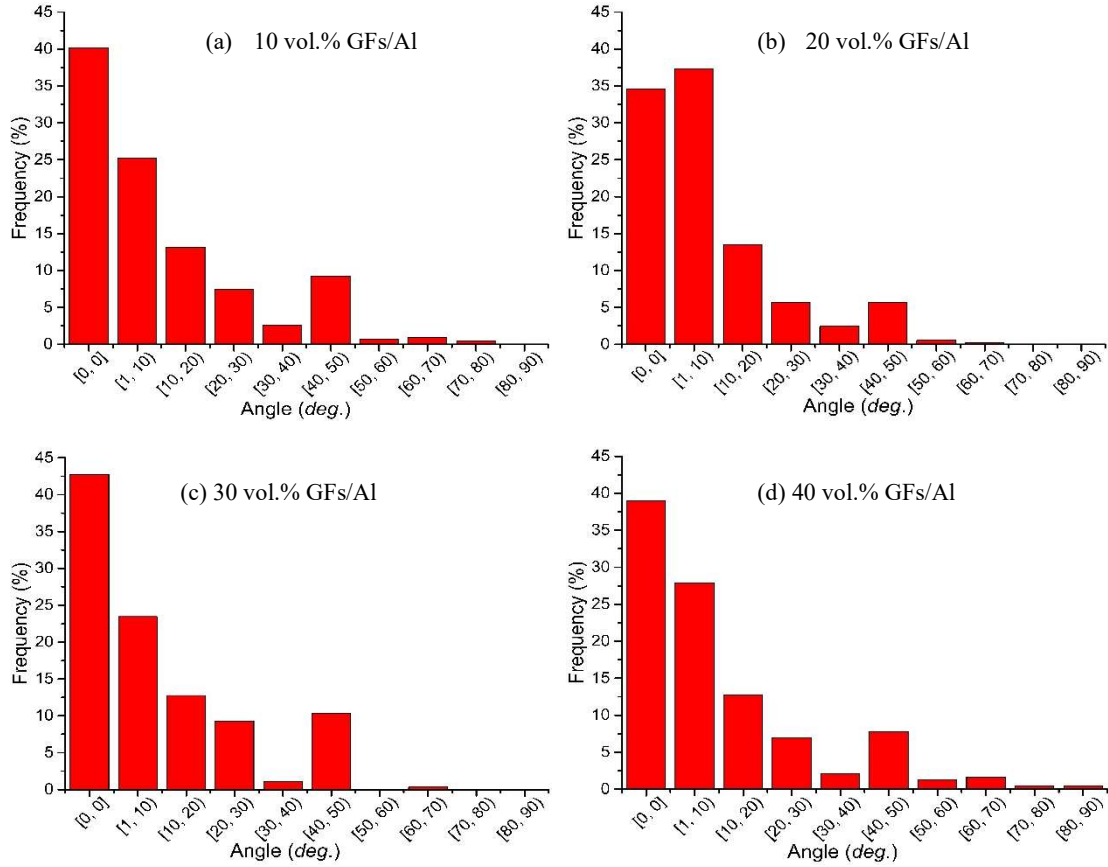


Figure 2-9 Distribution of GFs' orientation in the 10-40 vol.% GFs/Al composites, the orientation was described by the angle between GFs basal-plane direction and heat flow direction.

Figure 2-10 shows the statistical averages of the aspect ratio from the GFs in Figure 2-8. The aspect ratios values were mainly located at  $(1/10, 3/10]$ . Besides, the average diameters of GFs in composites ( $d$ ) were calculated by equivalent circle area (i.e.,  $d = 2\sqrt{S_{GFs}/\pi}$ ,  $S$  is the area), the values were 1.62, 2.54, 3.89, and 5.41 mm in 10~40 vol.% GFs/Al composites, respectively.

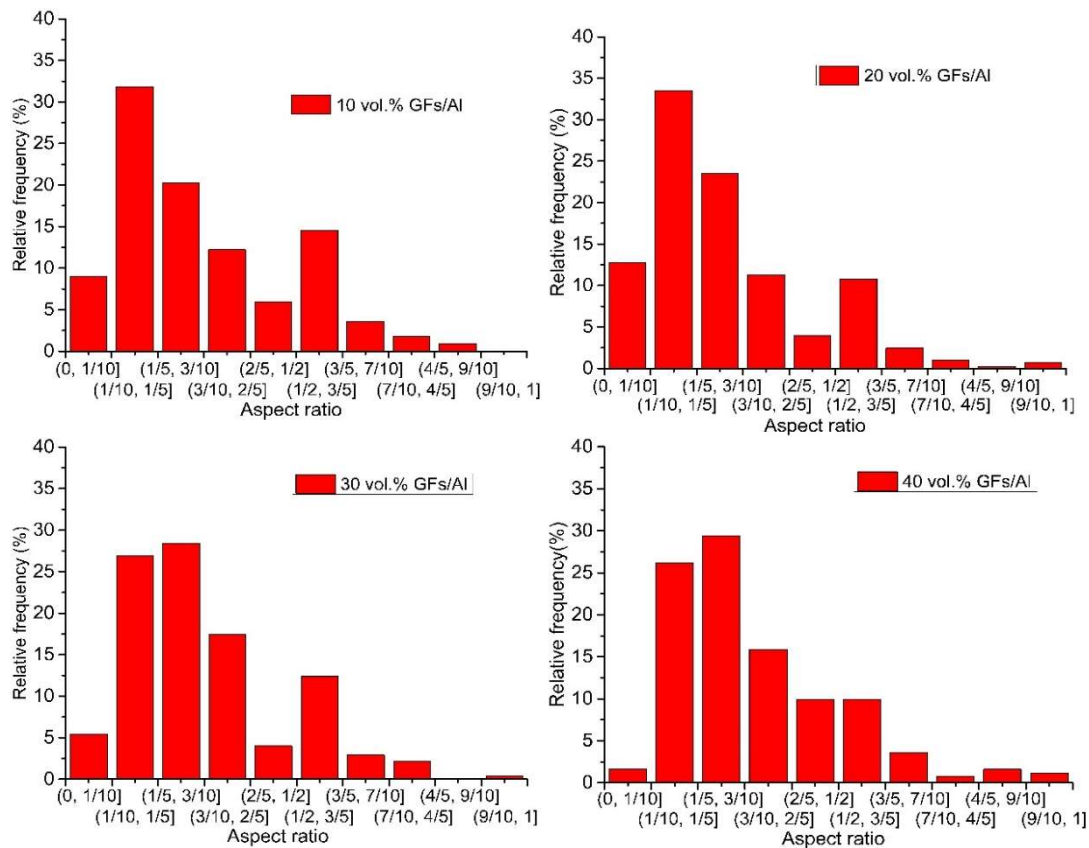


Figure 2-10 Statistical counting of the aspect ratio of GFs.

## 2.4.2 Interface reaction

As Etter et al. [24] reported if Al and GFs undergo a chemical reaction during the manufacturing process of GFs/Al composites, the formation of  $Al_4C_3$  exists at the Al-GFs interface. Therefore, XRD was used to determine the interface reaction at the Al-GFs interface, as shown in Figure 2-11. The result showed that the  $Al_4C_3$  phase was absent in 10-40 vol.% GFs/Al composites. The reason can be attributed to the following two points: first, the sintering temperature was less than the melting point of Al<sup>[25]</sup>; second, the raw material GFs was high graphitization<sup>[26]</sup>. On the other hand, high-temperature sintering methods were also commonly used to fabricate GFs/Al composites in the liquid metallurgy process. For example, Huang<sup>[9]</sup> and Chen<sup>[12]</sup> used vacuum hot pressing process to fabricate GFs/Al composites, Zhou<sup>[8]</sup> and Chang<sup>[16]</sup> used the pressure infiltration method to fabricate GFs/Al composites, and Li<sup>[13]</sup> used

squeeze casting technique to fabricate GFs/Al composites. The vacuum hot pressing process, pressure infiltration, and squeeze casting belong to the liquid metallurgy process, and a higher temperature than the melting point of Al is required to fabricate GFs/Al composites. Therefore, in order to hinder the formation of  $\text{Al}_4\text{C}_3$ , a high cooling rate was used in the liquid metallurgy process to reduce the reaction time between Al and GFs. Because the formation of  $\text{Al}_4\text{C}_3$  is associated with the reaction temperature and reaction time, according to the equation  $X = \sqrt{2Kt}$ , where  $X$  is the thickness of reaction layer,  $t$  is the reaction time, and  $K$  is the reaction rate related to the reaction temperature<sup>[27, 28]</sup>.

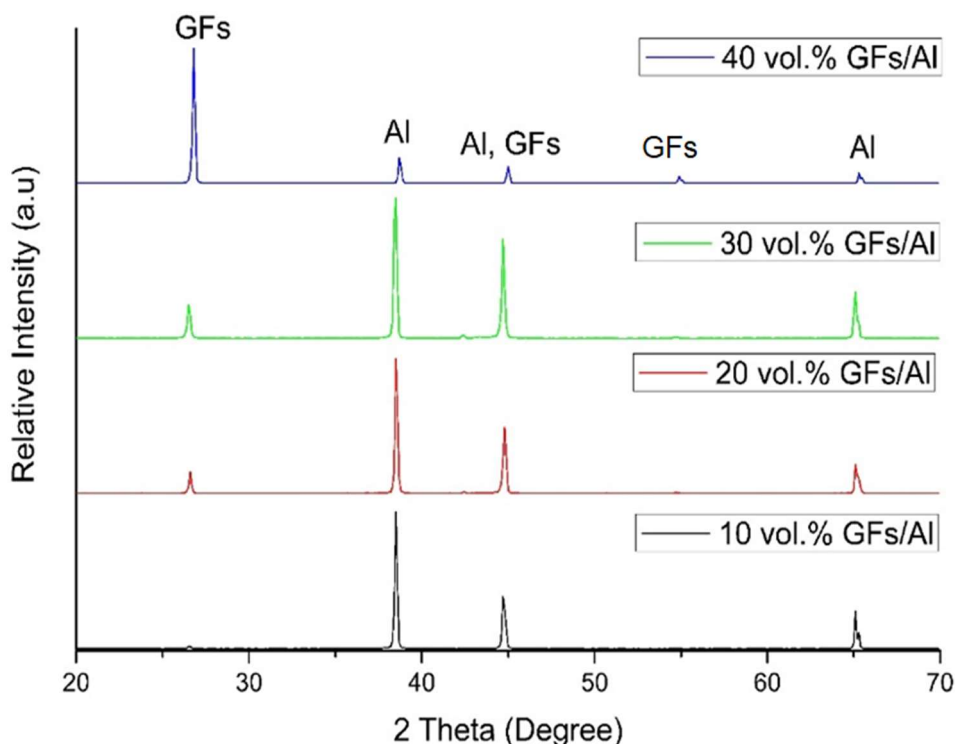


Figure 2-11 XRD analysis of GFs/Al composites.

Moreover, Huang et al.<sup>[9]</sup>, Chen et al.<sup>[12]</sup>, Zhou et al.<sup>[8]</sup>, Chang et al.<sup>[16]</sup>, and Li et al.<sup>[13]</sup> examined the microstructure and composition of the Al-GFs interface by HRTEM. Their studies revealed that an amorphous layer interface was formed at the Al-GFs interface, and the amorphous layer can be frequently observed, as shown in Figure 2-12<sup>[9]</sup>. Therefore, we assumed the amorphous layers were formed at the Al-Gfs

interfaces in the 10-40 vol.% GFs/Al composites.

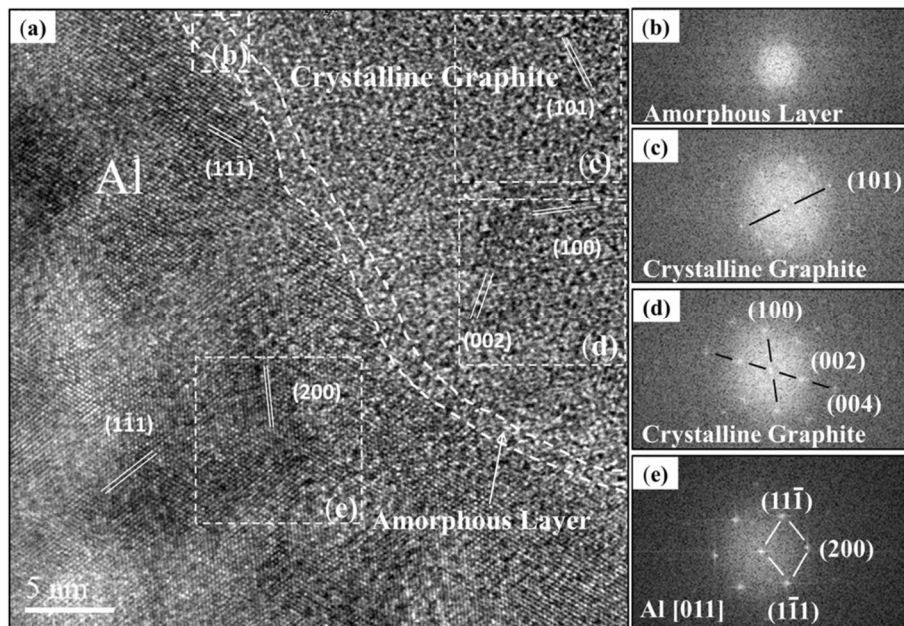


Figure 2-12 HRTEM image of the Al-GFs interface<sup>[9]</sup>: (a) HRTEM image; (b) FFT of amorphous layer; (c–d) FFT of crystalline graphite; (e) FFT aluminum.

### 2.4.3 Effective thermal conductivity and coefficient of thermal expansion

The effective thermal conductivities of 10-40 vol. % GFs/Al composites and the pure Al sample were measured by the steady-state thermal conductivity measuring device (see Figure 2-5). Figure 2-13 (a-d) shows the temperature variation in the ETC measurement process. The temperatures were a minor fluctuation and approximately maintained constant value, indicating that the measuring temperature was steady. Figure 2-13 (a1-d1) shows the average temperature with respect to the distance between thermocouples and the heating plate. The solid black dot is the average temperature of the thermocouples, the solid black line is the temperature gradient and is calculated using the equations (2-2) to (2-4). The dashed line is the location of the interface between the sample and hot bar or cold bar, and the corresponding temperature gap indicates the temperature difference at the interface.



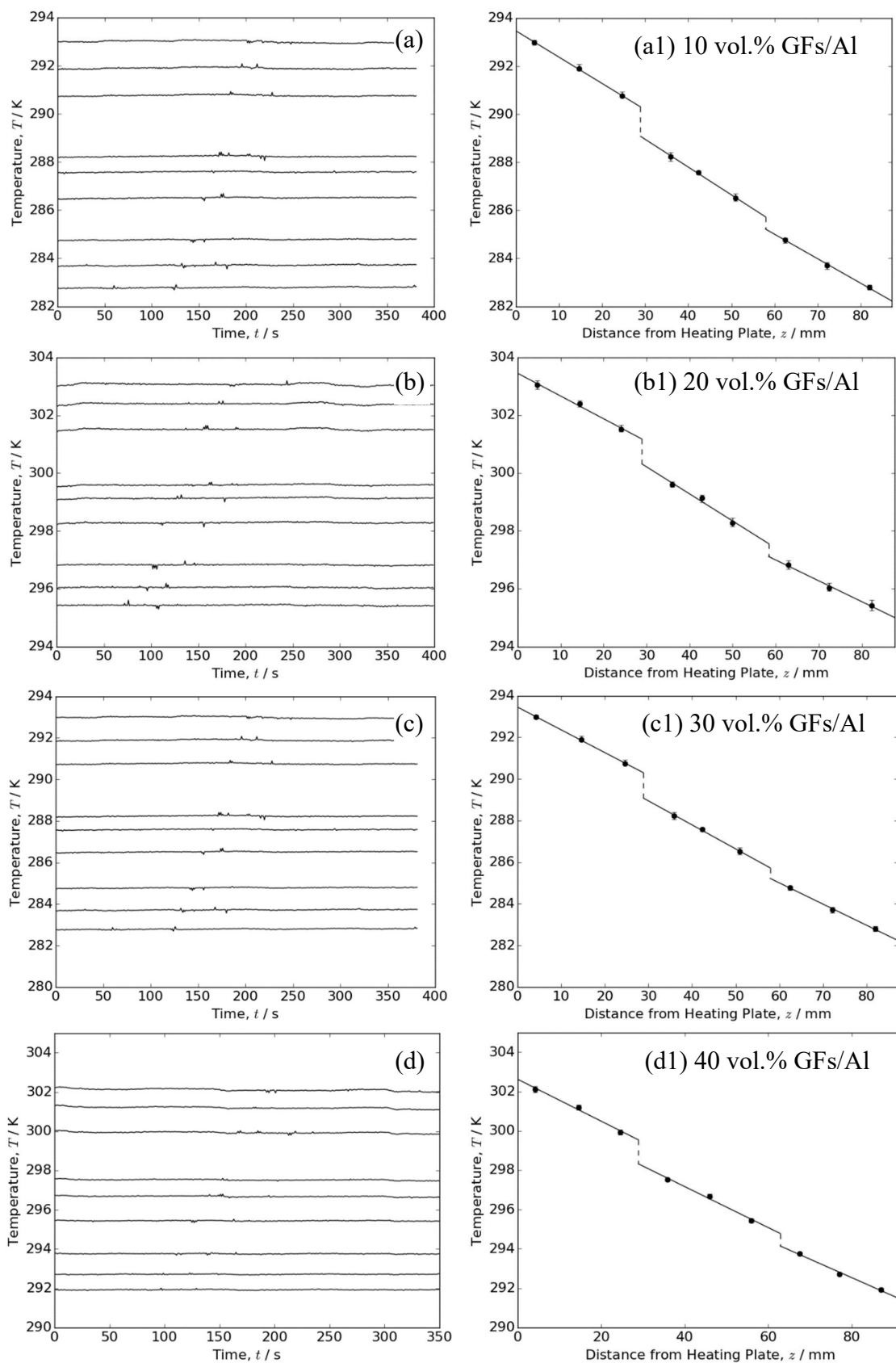


Figure 2-13 Temperature variation in the ETC measurement process of 10-40vol.%GFs/Al composites, (a-d) relationship between temperature and time, and (a1-d1) relationship between temperature and distance of the measuring points.

The ETCs of 10-40 vol.% GFs/Al composites and the pure Al sample were calculated by equations (2-2) to (2-5) and listed in Table 2-5. As a result, the TC of Al was  $217 \text{ Wm}^{-1}\text{K}^{-1}$ , and the ETC of GFs/Al composites  $\lambda_{eff}$  was increased from 248 to  $346 \text{ Wm}^{-1}\text{K}^{-1}$  as GFs volume fraction increasing from 10 to 40 vol.%. Moreover, the theoretical ETC ( $\lambda_{Theo.}$ ) of GFs/Al composites were calculated using the layer-in-parallel model<sup>[29]</sup> and the effective medium approximation (EMA) model.

The layer-in-parallel model considers the basal-plane of GFs parallel to heat flow direction, ETC ( $\lambda_{Theo.}$ ) of GFs/Al composites is calculated as follow:

$$\lambda_{Theo.} = v_{Al\_matrix}\lambda_{Al\_matrix} + v_{GFs}\lambda_{GFs} \quad \dots\dots (2-6)$$

where  $v$  is volume fraction,  $\lambda$  is thermal conductivity, the subscript  $Al\_matrix$  and  $GFs$  are Al-matrix and GFs, respectively. If to consider the effect of interfacial thermal resistance ( $h$ ) on ETC, the TC of GFs ( $\lambda_{GFs}$ ) can be replaced by  $\lambda_{GFs}^{eff}$ , given by<sup>[15]</sup>:

$$\lambda_{GFs}^{eff} = \lambda_{GFs} / \left( \frac{2\lambda_{GFs}}{hd} + 1 \right) \quad \dots\dots (2-7)$$

where a value of  $h$ ,  $5 \times 10^7 \text{ W m}^{-2} \text{ K}^{-1}$  <sup>[30]</sup> was used. Thus, equation (2-6) can be expressed as:

$$\lambda_{Theo.}^h = v_{Al\_matrix}\lambda_{Al\_matrix} + v_{GFs}\lambda_{GFs} / \left( \frac{2\lambda_{GFs}}{hd} + 1 \right) \quad \dots\dots (2-8)$$

The EMA model considers the geometrical factor and  $h$  to evaluate the ETC of composites, and the equation can be expressed as <sup>[15]</sup>:

$$\lambda_{EMA} = \lambda_{Al} \left( 1 + \frac{v_{GFs}}{\frac{\pi}{4p} (1 - v_{GFs}) + \frac{1}{\frac{\lambda_{GFs}^{eff}}{\lambda_{Al}} - 1}} \right) \quad (2-9)$$

where  $p$  is the aspect ratio of GFs, the value was 5 (see Figure 2-10).

Moreover, the measured relative densities of GFs/Al composites and pure Al sample were less than 1, indicating that some pores were embedded in the samples. However, equations (2-6), (2-8), and (2-9) do not include the effect of pores on the ETC. Therefore, the pores in this study were presumed to be embedded within the Al matrix. The TC of the pure Al sample and the Al-matrix,  $\lambda_{Al\_eff}$ , can be calculated by the equation derived by Landauer<sup>[31]</sup>. The equation was as follow:

$$\lambda_{Al\_eff} = \frac{1}{4} [\lambda_p(3v_p - 1) + \lambda_{Al}(3v_{Al} - 1) + ([\lambda_p(3v_p - 1) + \lambda_{Al}(3v_{Al} - 1)]^2 + 8\lambda_p\lambda_{Al})^{\frac{1}{2}}] \dots\dots (2-10)$$

where the  $\lambda_{Al}$  is the thermal conductivity of Al powder,  $\lambda_p$  is the thermal conductivity of pores,  $v_p$  and  $v_{Al}$  are the volume fraction of Al and pores. In addition, the values of  $\lambda_{Al}$ ,  $\lambda_p$ ,  $v_p$ , and  $v_{Al}$  were listed in Table 2-6.

Table 2-6 Calculation parameters in equation (2-7).

Samples	$v_{Al}$	$v_p$	$\lambda_p$	$\lambda_{Al}$	$\lambda_{Al\_eff}$	$\lambda_{Al\_matrix}$
Al	0.97	0.03	0.214	228	217	-
10 Vol.% GFs/Al	0.994	0.06	0.214	228	225.9	225.9
20 Vol.% GFs/Al	0.993	0.07	-	-	225.6	225.6
30 Vol.% GFs/Al	0.986	0.014	-	-	223.2	223.2
40 Vol.% GFs/Al	0.971	0.029	-	-	218.1	218.1

First, by substituting the measured TC ( $217 \text{ Wm}^{-1}\text{K}^{-1}$ ) and the relative density (0.97) of the pure Al sample into the equation (2-10), the TC of Al powder ( $\lambda_{Al}$ ) was calculated for  $228 \text{ Wm}^{-1}\text{K}^{-1}$ . Second, the TCs of the Al-matrix for 10-40 vol.% GFs/Al composites were calculated by substituting the  $\lambda_{Al}$ ,  $\lambda_p$ ,  $v_p$ , and  $v_{Al}$  from Table 2-6 into equation (2-10). The calculated values of  $\lambda_{Al\_eff}$  were 225.9, 225.6, 223.2, and 218.1  $\text{Wm}^{-1}\text{K}^{-1}$  for 10-40 vol.% GFs/Al composites, respectively.  $\lambda_{Al\_eff}$  was equal to  $\lambda_{Al\_matrix}$ . Therefore,

the theoretical ETC of the 10-40 vol.% GFs/Al composites can be calculated by equations (2-6), (2-8) and (2-9), and the values were listed in Table 2-7.

Table 2-7 Theoretical ETC of GFs/Al composites.

ETC	GFs/Al composites			
	10 Vol.% GFs	20 Vol.% GFs	30 Vol.% GFs	40 Vol.% GFs
$\lambda_{eff}$ ( $Wm^{-1}K^{-1}$ )	248	280	313	346
$\lambda_{Theo.}$ ( $Wm^{-1}K^{-1}$ )	291	357	420	483
$\lambda_{Theo.}^h$ ( $Wm^{-1}K^{-1}$ )	289	353	417	480
$\lambda_{EMA}$ ( $Wm^{-1}K^{-1}$ )	271	320	370	422

Figure 2-14 shows a comparison between the measured and theoretical ETCs. The theoretical value was considerably higher than the measured value  $\lambda_{eff}$ , and the deviation between the measured and the theoretical ETCs increased with GFs content. The largest deviation between the measured and theoretical ETCs was up to 28.4%.

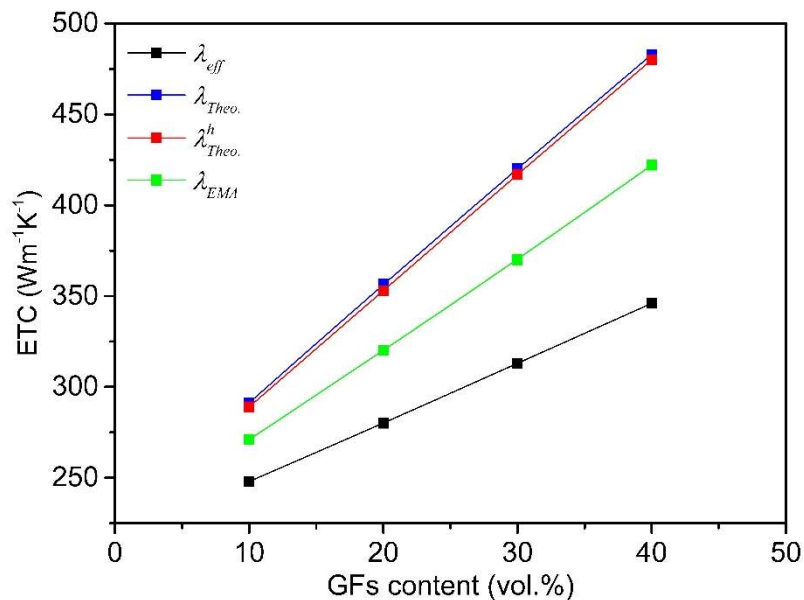


Figure 2-14 Comparison of measured and theoretical ETC of 10-40 vol.% GFs/Al composites.

The value of  $\lambda_{Theo}$  was approximately equal to that of  $\lambda_{Theo}^h$ , indicating that the effect of interfacial thermal resistance on the ETC was small. The deviation between  $\lambda_{Theo}$  and  $\lambda_{eff}$  can be attributed to the effects of the orientation of GFs and the interface thermal resistance. In addition, Zhou et al.<sup>[15]</sup> calculated the ETC of GFs/Al composites using the EMA approach, and the calculated values were agreed with the experimental values. However, the  $\lambda_{EMA}$  were greater than the  $\lambda_{eff}$ . The deviation between them can be attributed to the effect of the orientation of GFs.

The coefficient of thermal expansion of 10-40 vol.% GFs/Al composites were measured along the direction parallel to the basal plane of GFs, and the results were summarized in Table 2-7. Figure 2-15 (a) shows the thermal expansion at temperatures from 25 to 200 °C. The thermal expansion of GFs/Al composites decreased with the GFs content under the temperature from 25 to 200 °C, indicating that the coefficient of thermal expansion of GFs/Al composites decreases with the increase of GFs content.

Table 2-7 Coefficient of thermal expansion.

Samples	Al	GFs	GFs content (vol.%)			
			10	20	30	40
$\alpha_{Exp.}$ (ppm/K)		-1.5 <sup>[34-35]</sup>	26.2	24.4	23.4	22.1
$\alpha_{Com.}$ (ppm/K)	23.8 <sup>[12]</sup>	25 <sup>[37]</sup>	21.27	18.74	16.21	13.68

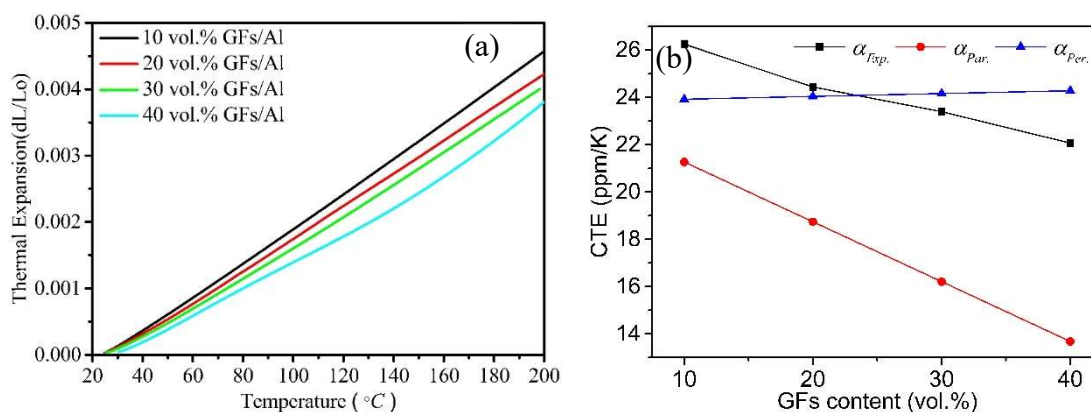


Figure 2-15 Thermal expansion variation versus temperature (a), coefficient of thermal expansion (CTE) versus GFs content (b).

To calculate the coefficient of thermal expansion of composite, the commonly used theoretical models include the rule of mixture [12], Turner model [32], Kerner model [33], and Schapery model [34]. But Chen et al. [12] and Kurita et al. [11] revealed that the Turner and Kerner models were failed to predict the coefficient of thermal expansion of 10-40 vol.% GFs/Al composites. Thus, the mixture rule here was used to evaluate the coefficient of thermal expansion of 10-40 vol.% GFs/Al composites. The rule of the mixture is described as follows:

$$\alpha_{Com.} = v_{Al\_matrix}\alpha_{Al\_matrix} + v_{GFs}\alpha_{GFs} \dots\dots (2-8)$$

where  $\alpha$  is the coefficient of thermal expansion, and the subscript *Com.* is composite. The coefficient of thermal expansion is 23.8 ppm/K [12] for Al and -1.5ppm/K [35, 36] for GFs along the basal plane direction and 25ppm/K [37] for GFs along the direction perpendicular to the basal plane. The calculated values were list in Table 2-7. Figure 2-15 (b) shows the measured and calculated coefficient of thermal expansion. The black line represents the experimental coefficient of thermal expansion, the experimental values of thermal expansion decreased with the increase of GFs content. The blue line presents the coefficient of thermal expansion parallel to GFs basal plane, the red line presents the coefficient of thermal expansion perpendicular to GFs basal plane. Because the orientations of some GFs were disorder in composites, the coefficient of thermal expansion of composite should be located between the red and blue lines. However, the black line was intersected with the blue line. The experimental values for the composites with 10-20 vol.% GFs were higher than the calculated values, and the experimental values exceeded even the coefficient of thermal expansion of Al and GFs (25 ppm/K [37] along the direction perpendicular to the basal plane). Besides, the black line was around the blue line, the reasons can be attributed to pores embedded in the composites and the disorder orientations of GFs.

## 2.5 Summary

GFs/Al composites with 10-40 vol.% GFs were fabricated by using spark plasma sintering in this chapter. Subsequently, the following studies on the GFs/Al composites were completed: 1) The relative density of composites was measured by the Archimedes method; 2) The microstructure images analyzed the distribution and orientation of GFs in the composites; 3) The formation of  $\text{Al}_4\text{C}_3$  was investigated by XRD method; 4), the effective thermal conductivities of the composites were measured using a steady-state thermal conductivity measurement device and calculated by theoretical methods; 5) The coefficient of thermal expansions of the composites were measured by a thermal dilatometer (NETZSCH DIL 402 C) and calculated by the rule of mixture. Finally, the following conclusions were obtained.

- 1) The relative density was up to 99% for the GFs/Al composites; with the GFs content up to 40%, the relative density of composites decreased to 97.1%.
- 2) GFs were homogeneously distributed and stacked by layers in the Al matrix. The orientations of nearly 70% GFs were parallel to each other. Moreover, with the GFs content up to 40 vol.%, GFs exhibits aggregation and tend to form the network system by connecting GFs.
- 3) The formation of  $\text{Al}_4\text{C}_3$  was absent in the 10-40 vol.% GFs/Al composites, but amorphous layers at the Al-GFs interface may be formed at the interface.
- 4) The measured effective thermal conductivities were 248, 280, 313, and 346  $\text{W m}^{-1} \text{K}^{-1}$  for 10-40 vol.% GFs/Al composites, respectively. The effects of orientation of GFs and interface thermal resistance caused the ETC of GFs/Al composites to be lowered by 28.4%.
- 5) The experimental values of the coefficient of thermal expansion were 26.2, 24.4, 23.4, and 22.1 ppm/K for 10-40 vol.% GFs/Al composites, respectively. Moreover, because of the effects of the pores embedded in composites and the disorder orientations of GFs, the experimental values for the composites with 10-20 vol.% GFs were higher than the calculated values perpendicular to GFs basal plane, and the experimental values exceeded even the coefficient of

thermal expansion of Al and GFs (25 ppm/K along the direction perpendicular to the basal plane).



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## *Chapter 3*

### **Relationship between the thermal conductivity of GFs in 3D and 2D microstructure images**

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### **3.1 Introduction**

Graphite flakes (GFs) reinforced aluminum (Al) matrix composites (GFs/Al) are light-weight, low cost, and exhibit superior thermal properties, such as high effective thermal conductivity (ETC) and low coefficient of heat expansion (CTE). Therefore, GFs/Al composites are considered as one of the most promising thermal management materials. However, the ETCs of GFs/Al composites are affected by the orientation of GFs, the interfacial thermal resistance, and the interface reaction between Al and GFs [1-5]. In particular, the interfacial thermal resistance and the orientation of GFs have a great effect on the ETCs of GFs/Al composites, which results in lower experimental ETCs than expected.

The commonly used theoretical methods using to calculate the ETC of GFs/Al composites include mixing rules [1, 2, 6], Fricke equation<sup>[7]</sup>, effective medium approach (EMA)<sup>[8]</sup>, and three-dimensional simulations [9, 10]. However, the orientations of GFs in composites are uncontrollable diverse in the composites. Therefore, obtaining the reliable ETC of GFs/Al composites is challenging using those theoretical models. Three-dimensional (3D) image simulations can be used to calculate the ETC of GFs/Al composites well. But the 3D-images simulation is expensive and has low timeliness because the 3D models are constructed using the ultra-high-resolution X-ray computed tomography images<sup>[11]</sup>. In addition, two-dimensional (2D) image simulations are also often used to evaluate the ETC of composites [12, 13], and it is simple, saves time, and is low cost. However, the orientation of GFs in the 2D microstructure images may not be equivalent to that in the 3D microstructure images due to a 2D microstructure image cannot provide the information in the depth direction with respect to the viewing surface. Therefore, the ETC calculated by the 2D image-based simulation is not reliable.

In order to develop a 2D image simulation method for calculating the ETC of GFs/Al composites considering the orientation of GFs and the interfacial thermal resistance, the relationship between the orientation of GFs in 2D and 3D models was investigated. Moreover, the ETC of the 10 vol % GFs/Al composite was calculated

using 2D image-based simulation, and the interfacial thermal conductance at the GFs-Al interface was calculated using the reversed method.

## 3.2 Calculation Procedure

### 3.2.1 Creating 3D and 2D models

Figure 3-1 shows the flowchart of the calculation procedure for comparing the orientations of GFs in the 2D model and the corresponding 3D model. First, some 3D models were created, as shown in Figure 3-2 (a).

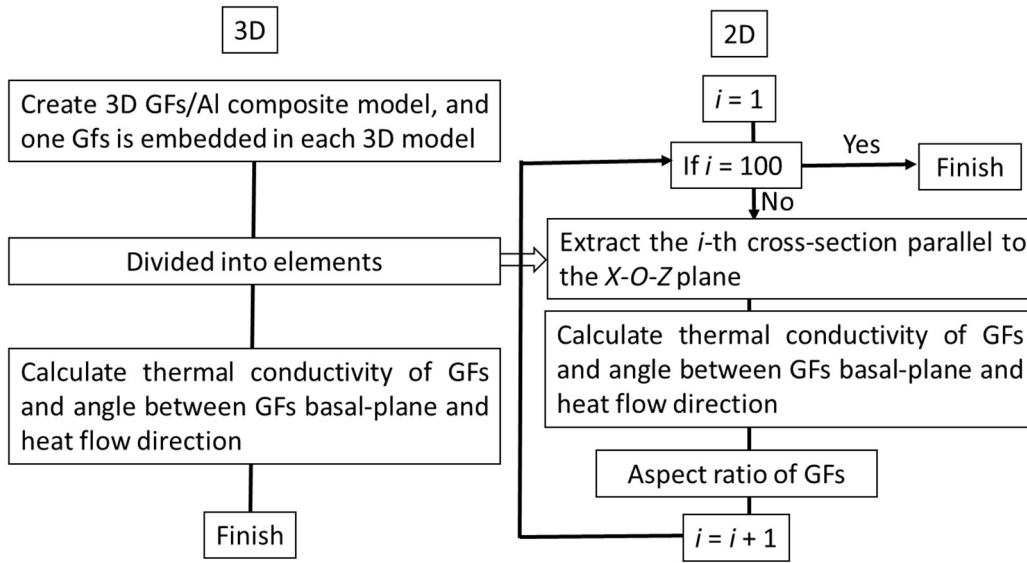


Figure 3-1. Flowchart of the calculation procedure for comparing the orientations of GFs in two- and three-dimensional models.

The dimensions of the 3D model were  $110 \times 100 \times 100$  elements in the  $X$ -,  $Y$ -, and  $Z$ -axis. One GFs was embedded in the center of one 3D model. The shape of GFs was an elliptical cylinder, the semi-major axis( $D$ ) was 40 elements, and the semi-minor axis ( $d$ ) was 20 elements. The thickness of the ellipse ( $T$ ) was 10 elements. As shown in Figure 3-2 (b), the GFs in the 3D model can rotate  $\theta_x$ ,  $\theta_y$ , and  $\theta_z$  around the  $X$ -,  $Y$ -, and  $Z$ -axis, respectively, according to the right-hand rotation rule. The rotated coordinates  $(x, y, z)$  were calculated using the following rotation matrix:

$$\begin{bmatrix} x \\ y \\ z \end{bmatrix} = \begin{bmatrix} \cos\theta_z & -\sin\theta_z & 0 \\ \sin\theta_z & \cos\theta_z & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} \cos\theta_y & 0 & \sin\theta_y \\ 0 & 1 & 0 \\ -\sin\theta_y & 0 & \cos\theta_y \end{bmatrix} \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos\theta_x & -\sin\theta_x \\ 0 & \sin\theta_x & \cos\theta_x \end{bmatrix} \begin{bmatrix} x_0 \\ y_0 \\ z_0 \end{bmatrix} \dots\dots (3-1)$$



where  $(x_0, y_0, z_0)$  is the initial coordinates. The  $x$  and  $y$  are set to coincide with the semi-major axis( $D$ ) and semi-minor axis( $d$ ) of GFs, respectively. In this study, the orientation of GF in 3D models was represented by the angles  $\alpha$ ,  $\beta$ , and  $\gamma$ , which is the difference between the global coordinate  $X$ - $Y$ - $Z$  and the local coordinate  $x$ - $y$ - $z$ . In other words,  $\alpha$ ,  $\beta$ , and  $\gamma$  were the angles between the axes ( $X$ - $x$ ,  $Y$ - $y$ , and  $Z$ - $z$ ), and their values were calculated as follows:

$$\alpha = \arccos \left( \frac{\overline{ox} \cdot \overline{OX}}{|\overline{ox}| |\overline{OX}|} \right), \beta = \arccos \left( \frac{\overline{oy} \cdot \overline{OY}}{|\overline{oy}| |\overline{OY}|} \right), \gamma = \arccos \left( \frac{\overline{oz} \cdot \overline{OZ}}{|\overline{oz}| |\overline{OZ}|} \right) \dots\dots (3-2)$$

where  $\overline{OX}$ ,  $\overline{OY}$  and  $\overline{OZ}$  are the unit coordinate vectors, and their values are (1,0,0), (0,1,0), and (0,0,1), respectively.  $\overline{ox}$ ,  $\overline{oy}$ , and  $\overline{oz}$  were calculated using equation (3-1).

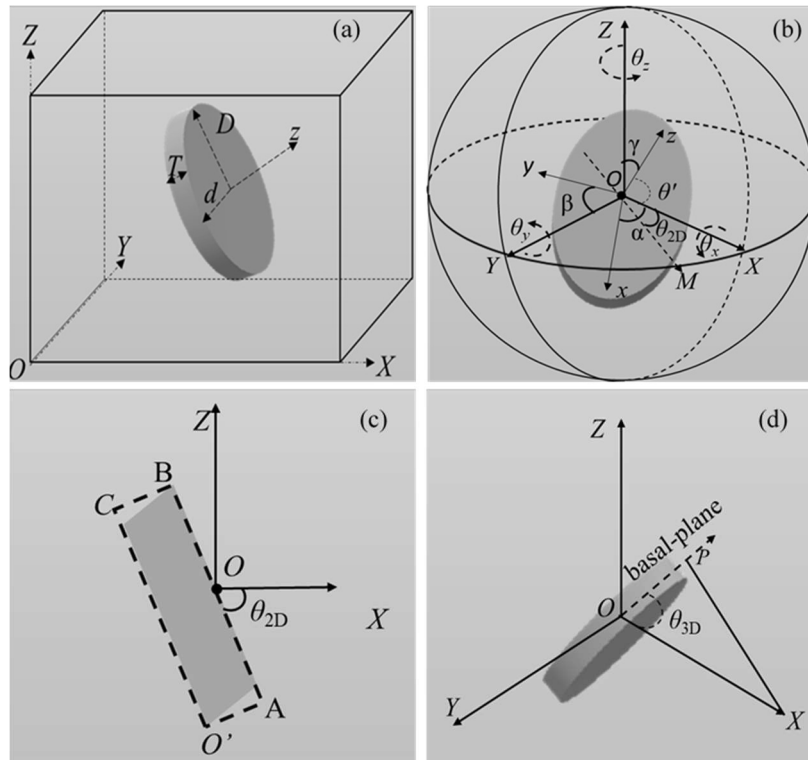


Figure 3-2. Three-dimensional model and orientation of GFs. a) 3D model of GFs/Al composites; b) Orientation of GFs in 3D model; c) Orientation of GFs in the 2D cross-sectional image; (d) Describing the angle between of GFs basal plane and heat flow;  $\alpha$ ,  $\beta$ ,  $\gamma$  and  $\theta'$  are the angles between the axes,  $x$ - $X$ ,  $y$ - $Y$ ,  $z$ - $Z$ , and  $z$ - $X$ , respectively;  $\alpha$ ,  $\beta$  and

$\gamma$  are within  $[0,90^\circ]$ ;  $OM$  is the line perpendicular to  $OY$  and  $Oz$ , and  $\theta_{2D}$  is the angle between  $OX$  and  $OM$ ; Rectangle  $O'ABC$  is the minimum bounding rectangle of GFs,  $\theta_{2D}$  is the angle between  $X$ -axis and  $AB$ ;  $XP$  is perpendicular to the basal plane of GFs.

Second, as many 2D cross-sections were extracted from each 3D model as possible, the cross-sections were extracted parallel to the  $X$ - $O$ - $Z$  plane, as shown in Figure 3-2 (b). The shape of GFs may be irregular in the cross-sectional images; therefore, a minimum bounding rectangle ( $O'ABC$ , see Figure 3-2 (c)) of the GFs was introduced to measure the orientation of the GFs. The orientation of GFs in cross-sectional images was defined by the angle ( $\theta_{2D}$ ) between GFs and the  $X$ -axis, as shown in Figure 3-2 (c).  $\theta_{2D}$  was calculated from the 2D cross-sectional images, excluding GFs smaller than  $5 \times 5$  elements. Moreover, the aspect ratio ( $R$ ) of GFs in each extracted 2D cross-sectional image was also calculated.

### 3.2.2 TC calculation of GFs

Here, the heat flow direction in the 3D model and the 2D cross-sectional image was set to be along the  $X$ -axis. The TC of GFs along the  $X$ -axis thus can be calculated as follows <sup>[14]</sup>:

$$\begin{aligned}\lambda_{//} &= \lambda_a \left[1 - \left(1 - \frac{\lambda_c}{\lambda_a}\right) \sin^2 \theta\right] \\ \lambda_{\perp} &= \lambda_a \left[1 - \left(1 - \frac{\lambda_c}{\lambda_a}\right) \cos^2 \theta\right]\end{aligned}$$

..... (3-3)

where  $\lambda_{//}$  is the TC of GFs parallel to the  $X$ -axis direction,  $\lambda_a$  is the TC of GFs in the basal-plane direction ( $880 \text{ W m}^{-1} \text{ K}^{-1}$ ), and  $\lambda_c$  is the TC of GFs in the out-of-plane direction ( $38 \text{ W m}^{-1} \text{ K}^{-1}$ ).  $\theta$  in the 3D models refers to the orientations of GFs relative to the heat flow direction, which are equal to the angles between the  $X$ -axis and basal plane of the GFs (Figure 3-2 (d)), and it is denoted as  $\theta_{3D}$ . Thus,  $\theta_{3D}$  can be obtained as follows:

$$\theta_{3D} = |90 - \theta'|, \quad \theta' = \arccos \left( \frac{|\vec{Oz} \cdot \vec{OX}|}{|\vec{Oz}| |\vec{OX}|} \right)$$

..... (3-4)

where  $\theta'$  is the angle between  $X$ -axis and  $z$ -axis (Figure 3-2 (b)). In other words,  $\theta'$  is the angle between the heat flow direction and the normal vector of the basal plane, and  $\theta_{3D}$  is the angle between the heat flow direction and basal plane.

In the 2D cross-section images,  $\theta$  refers to the angle between the cross-section of the GFs and the  $X$ -axis, i.e.,  $\theta = \theta_{2D}$ . Therefore,  $\theta_{2D}$  can be calculated mathematically, as follows:

$$\theta_{2D} = \arccos \left( \frac{\overrightarrow{OM} \cdot \overrightarrow{OX}}{|\overrightarrow{OM}| |\overrightarrow{OX}|} \right) \quad \dots\dots (3-5)$$

where  $OM$  is the line perpendicular to  $OY$  and  $Oz$  (Figure 3-2 (b)), and  $\theta_{2D}$  is in the range  $[0, 90)$ .

### 3.2.3 Image-based simulation method

The finite-volume method was employed to calculate the 2D temperature distributions. The temperature of the elements can be calculated using the following equation [13]:

$$T_{x,y}^{n+1} = T_{x,y}^n + \frac{\Delta t}{\rho c} \left( \frac{q_{x+1,y}^n - q_{x-1,y}^n}{\Delta x} + \frac{q_{x,y+1}^n - q_{x,y-1}^n}{\Delta y} \right) \quad \dots\dots (3-6)$$

where  $T_{x,y}^n$  is the temperature of the element at  $(x, y)$  coordinates,  $T_{x,y}^{n+1}$  is the temperature of the element at  $(x, y)$  coordinates after a time  $\Delta t$ ,  $\rho$  is the density,  $c$  is the specific heat, and  $\Delta x$  and  $\Delta y$  are the sizes of the elements. The heat flow  $q^n$  can be calculated as follows:

$$q_{x+1,y}^n = \lambda \left( \frac{T_{x+1,y}^n - T_{x,y}^n}{\Delta x} \right), \quad q_{x-1,y}^n = \lambda \left( \frac{T_{x,y}^n - T_{x-1,y}^n}{\Delta x} \right)$$

$$q_{x,y+1}^n = \lambda \left( \frac{T_{x,y+1}^n - T_{x,y}^n}{\Delta y} \right), \quad q_{x,y-1}^n = \lambda \left( \frac{T_{x,y}^n - T_{x,y-1}^n}{\Delta y} \right)$$

When the heat flow  $q^n$  move at the GFs-Al interface:

$$q_{x+1,y}^n = h(T_{x+1,y}^n - T_{x,y}^n), \quad q_{x-1,y}^n = h(T_{x,y}^n - T_{x-1,y}^n)$$

$$q_{x,y+1}^n = h(T_{x,y+1}^n - T_{x,y}^n), \quad q_{x,y-1}^n = h(T_{x,y}^n - T_{x,y-1}^n) \quad \dots\dots (3-7)$$

where  $h$  is the interfacial thermal conductance between different materials.

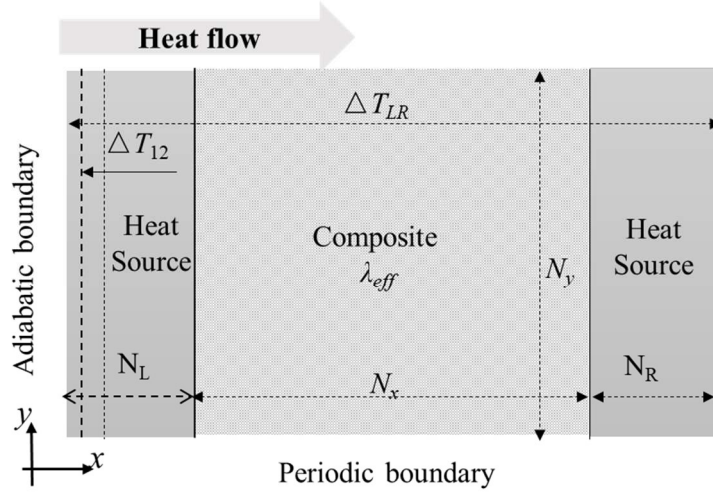


Figure 3-3 Simulation model for the ETC calculation.

Figure 3-3 shows the schematics of the simulation model. The model is a sandwich structure consisting of two heat sources and a composite part. The composite part was obtained from the microstructure image of the GFs/Al composites. The size of the composite part ( $N_x \times N_y$ ) was  $450 \times 600$  elements, and that of the heat source was  $5 \times 600$  elements ( $N_L = N_R = 5$  elements). The size of each element was  $1.18 \times 10^{-6}$  m. Both the upper and lower sides were the periodic boundary, while the left and right sides were the adiabatic boundary. The temperatures of the left and right edge elements were fixed at 301 K and 300 K, respectively. The initial temperature of the other elements was set at 300 K, and the temperature was iteratively updated until the temperature variation was lower than  $10^{-13}$  K. The temperature distribution at this moment was in a steady state, and the ETC of the GFs/Al composite,  $\lambda_{eff}$ , was calculated as follows:

$$\lambda_{eff} = \frac{\lambda_{Al} \Delta T_{12} N_x}{\Delta T_{LR} - N_L \Delta T_{12} - N_R \Delta T_{12}} \quad \dots\dots (3-8)$$

where  $\lambda_{Al}$  is the TC of the Al matrix,  $\Delta T_{12}$  is the average value of the temperature

difference between the first and second columns, and  $N_L$ ,  $N_R$ ,  $N_x$ , and  $N_y$  are the number of elements along the corresponding direction.

### 3.3 Results and Discussions

#### 3.3.1 Orientations of GFs in 2D and 3D models

The 2D orientations of GFs, i.e.,  $\theta_{2D}$ , were obtained from Equation (5) and the 2D-cross-sectional images, respectively. Moreover, the  $\theta_{2D}$  from the 2D-cross-sectional images extracted from the same 3D model showed the same values. Figure 3-4 shows a comparison of  $\theta_{2D}$  obtained from the mathematical calculation and the 2D-cross-sectional images. The solid line is the  $\theta_{2D}$  from Equation (3-5), and the open circles are the  $\theta_{2D}$  from the 2D-cross-sectional images. Thus, the open circles are on the solid line ( $x = y$ ), confirming that the  $\theta_{2D}$  values obtained from the 2D cross-sectional images were reliable.

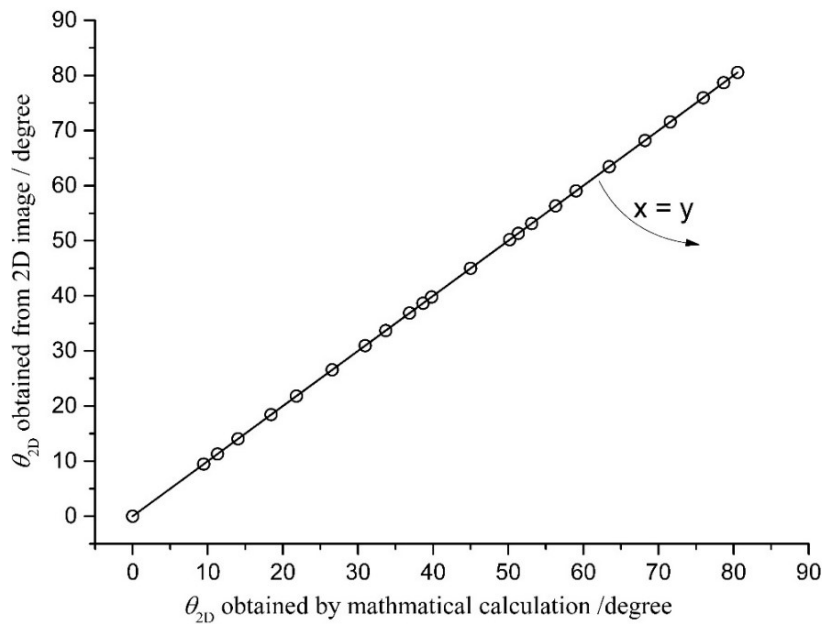


Figure 3-4 Comparison of the  $\theta_{2D}$  obtained from the mathematical calculation and 2D cross-sectional image.

The angle difference between  $\theta_{2D}$  and  $\theta_{3D}$  ( $\theta_{2D-3D}$ ) was calculated. Figure 3-5 (a) shows  $\theta_{2D-3D}$  in a standard regular triangle. The black dots represent the calculated points. The black color indicates that the value of  $\theta_{2D-3D}$  is zero. The other colors indicate that the value of  $\theta_{2D-3D}$  is between  $0^\circ$  and  $90^\circ$ . A few black dots were in the

black region, whereas the rest were distributed in other colored regions. In the triangle, most of the  $\theta_{2D-3D}$  in the triangle was under  $20^\circ$ , but there were some  $\theta_{2D-3D}$  that exceeded  $40^\circ$ . Additionally, the difference between the TCs of GFs in the 2D and 3D models (i.e.,  $TC_{3D-2D}$ ) was shown in Figures 3-5 (b). The maximum value of  $TC_{3D-2D}$  was up to  $840 \text{ W m}^{-1} \text{ K}^{-1}$ . Moreover, comparing Figure 3-5 (a) and Figure 3-5 (b), the value of  $TC_{3D-2D}$  can be large, even if the value of  $\theta_{2D-3D}$  was small in the triangle.

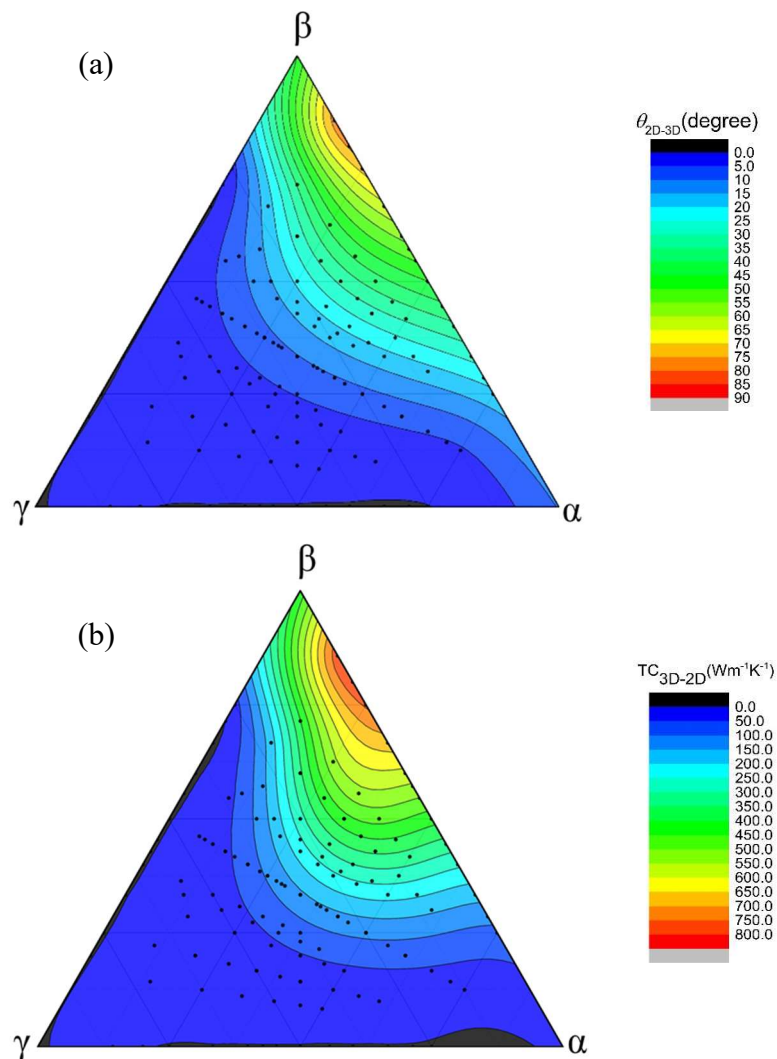


Figure 3-5. Orientation difference of GFs in the 2D cross-sections and 3D models (a); TC difference of GFs in the 2D cross-sections and 3D models (b).

### 3.3.2 Relationship between the aspect ratio of GFs and the orientation of GFs

The aspect ratio of the GFs ( $R$ ) in each 2D cross-section extracted from the same 3D model showed different values, while the corresponding  $\theta_{2D}$  was the same as each other. In order to investigate the relationship between the aspect ratios and the orientations, the average arithmetic value of  $R$  was calculated and marked as  $\bar{R}$ , i.e.,  $\bar{R} = \frac{\sum R}{\sum i}$  ( $i$  was the number of  $R$ ). Figure 3-6 shows the relationships between  $\theta_{2D}$ ,  $\bar{R}$ , and  $\theta_{2D-3D}$ .

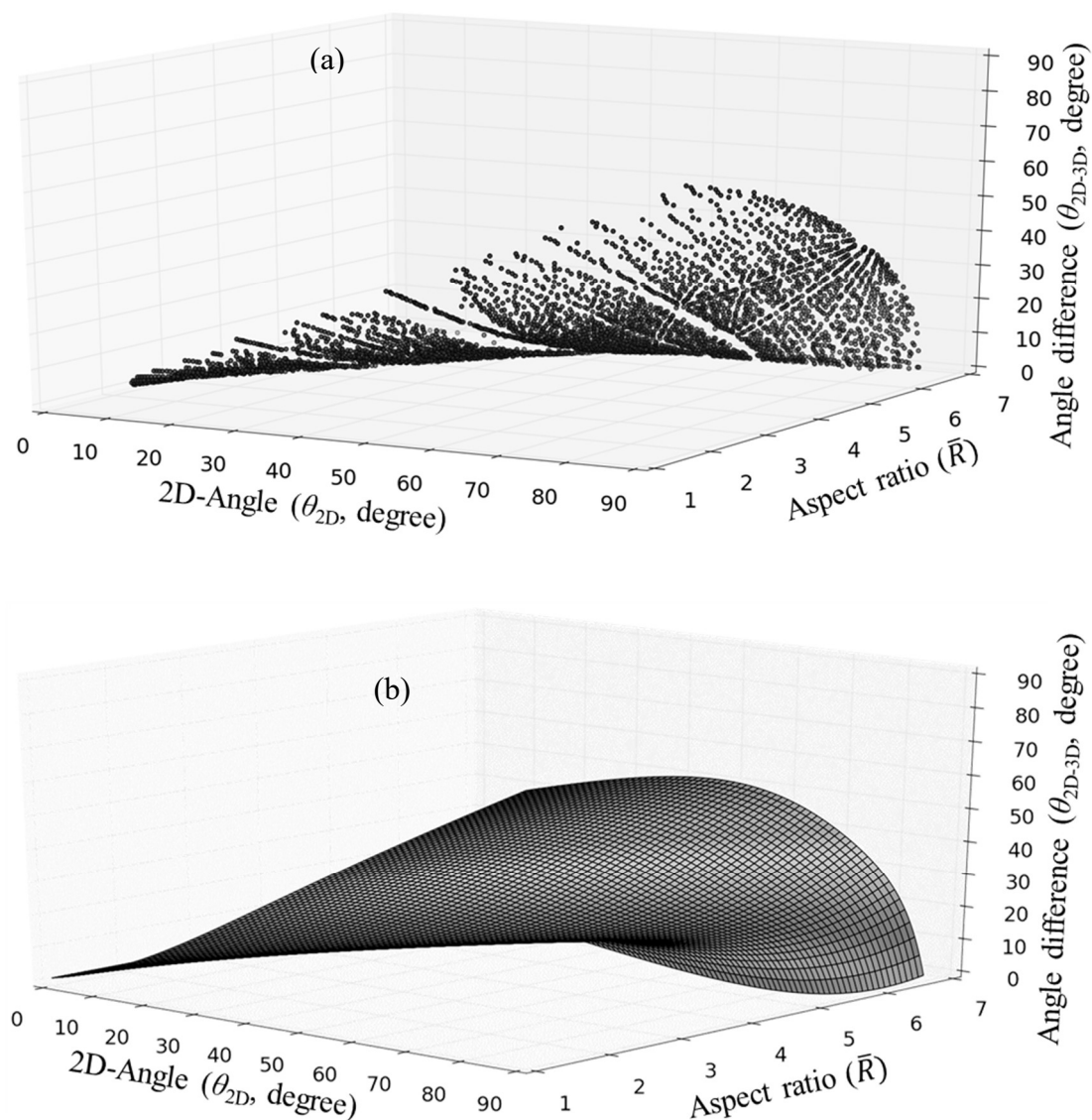


Figure 3-6. Relationship between the  $\theta_{2D}$ ,  $\theta_{2D-3D}$  and  $\bar{R}$ . Scatter image (a); Nonlinear polynomial fitting surface (b).



Figure 3-6 (a) shows that all data was distributed on a curved surface, which means that anyone of the  $\theta_{2D-3D}$ ,  $\bar{R}$  and  $\theta_{2D}$  can be expressed by using the curved surface function. Subsequently, a similar curved surface was obtained by the surface fitting method, as shown in Figures 3-6 (b). The function of the fitted curved surface can be expressed as:

$$\theta_{fitted-3D} = \theta_{2D} - [A_0 + A_1 \times \theta_{2D} \times f - A_4 \times \theta_{2D} \times e^{(A_5 \times \bar{R})}]$$

$$f = \frac{1}{1 + e^{-(A_2 + A_3 \times \theta_{2D} - \bar{R})}}$$

$$(0^\circ < \theta_{2D} < 90^\circ, 1 < \bar{R} < 6.7)$$

..... (3-9)

The coefficients,  $A_0, A_1, A_2, A_3, A_4, A_5$ , were listed in Table 3-1,  $\theta_{fitted-3D}$  represents the angle between GFs basal-plane and heat flow direction in the 3D model. This function can effectively convert  $\theta_{2D}$  into  $\theta_{3D}$  and improve the reliability of the ETC calculated by 2D image simulations. The deviation between the  $\theta_{fitted-3D}$  and  $\theta_{3D}$  was within  $6^\circ$ .

Table 3-1. Calculation parameters in equation (3-9)

Equation	$A_0$	$A_1$	$A_2$	$A_3$	$A_4$	$A_5$
(3-6)	0.64	0.86	0.62	0.059	1.80	3.15

To determine the reliability of function Equation (3-9), a comparison between the TCs of GFs calculated with  $\theta_{3D}$  and  $\theta_{fitted-3D}$  was conducted. Figure 3-7 shows the GFs' TCs calculated with  $\theta_{3D}$  and  $\theta_{fitted-3D}$  by Equation (3-3). The solid black line represents TCs calculated with  $\theta_{3D}$ , and grey circles were for  $\theta_{fitted-3D}$ . The circles were tightly bound around the solid black line, and the errors were so minor that they can be ignored. Thus, Figure 3-7 indicated that equation (3-6) is successful in converting  $\theta_{2D}$  to  $\theta_{3D}$ .

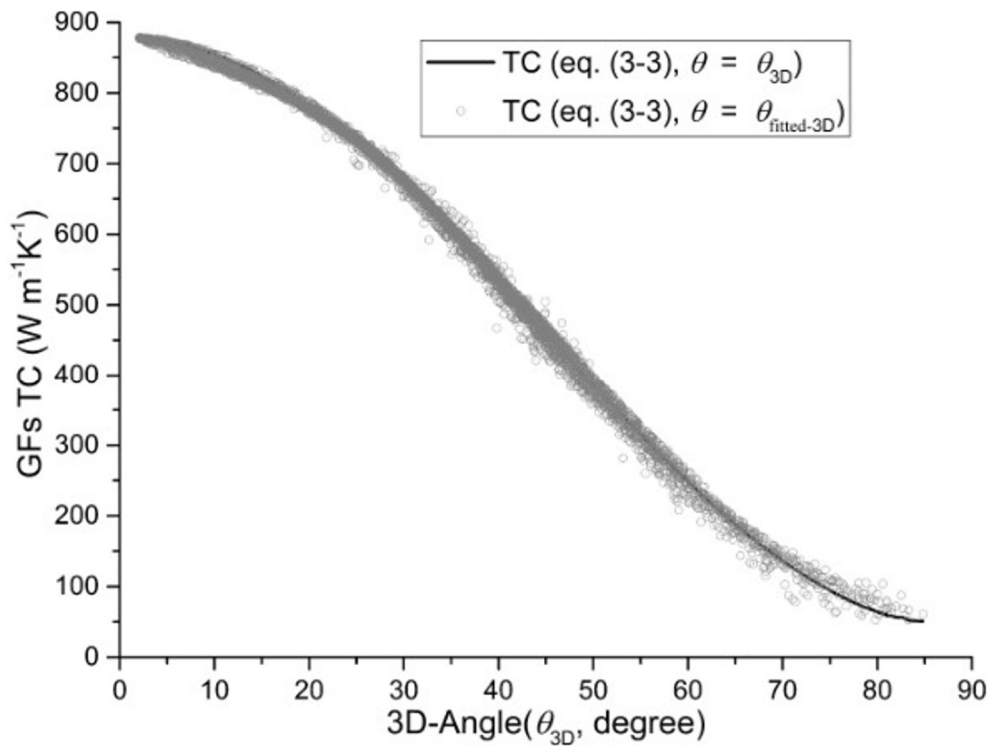


Figure 3-7. TC of GFs calculated by equation (3-3),  $\theta$  was equal to the  $\theta_{3D}$  for the black solid line and  $\theta_{fitted-3D}$  for the gray circle.

### 3.3.3 ETC calculation of GFs/Al composites

10 vol.% GFs/Al composites were fabricated via spark plasma sintering. The diameter of the GFs was  $137.02 \mu\text{m}$  (98% purity), and the average particle of Al powder was  $30 \mu\text{m}$  (99.9% purity). The relative density of the samples was 99.4%. The ETC of the sample, measured using a steady-state thermal-conductivity-measuring device, was  $238 \text{ W m}^{-1} \text{ K}^{-1}$ . The microstructures of the samples were observed using an optical microscope. Figure 3-8 (a) shows the microstructures of the 10 vol.% GFs/Al composites. To evaluate the TC distributions of GFs in the microstructures, different colors were used to mark the GFs according to the orientations of the GFs, as shown in Figure 3-8(b). For example, the red color indicates that the orientation of the GFs was parallel to the heat flow direction, and the GFs had high TC; the other colors indicate that the GFs tilt from the heat flow direction, decreasing TC.

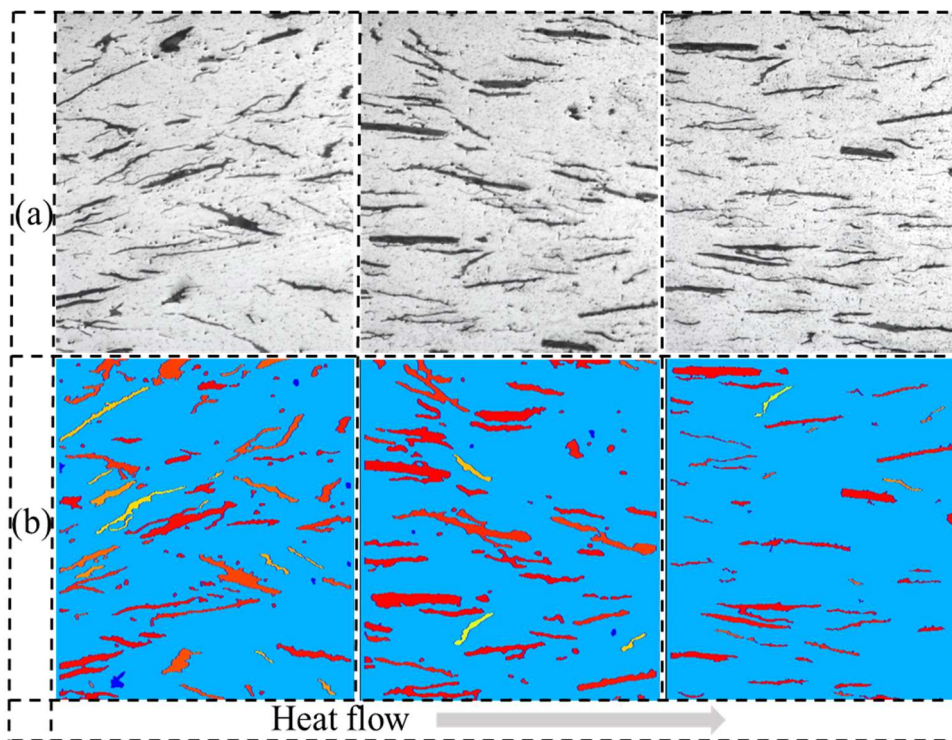


Figure 8 Microstructure images of the 10 vol.% GFs/Al composite; (a) The OM images; (b) Distribution of the GFs TC.

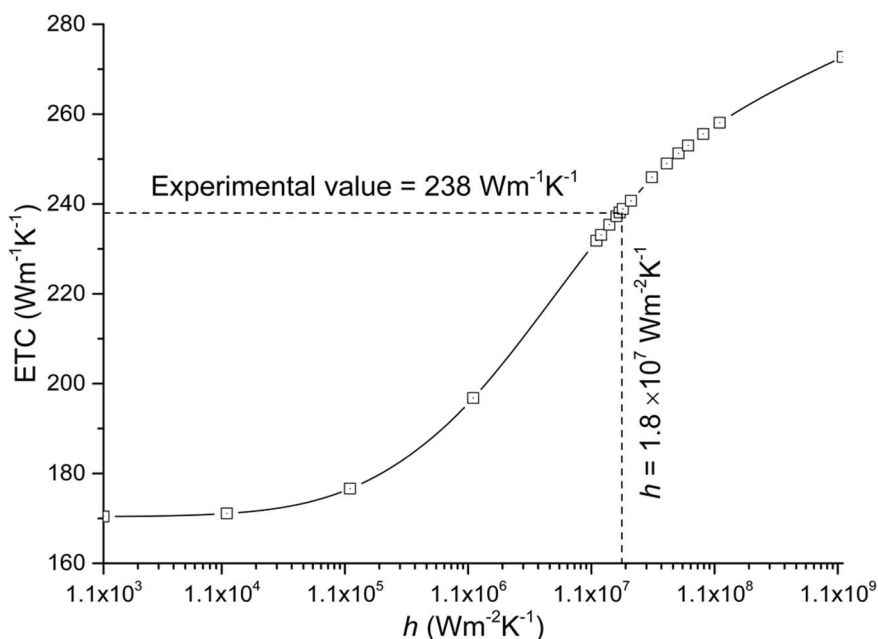


Figure 9. ETC of the 10 vol.% GFs/Al composite considering the interfacial thermal conductance ( $h$ ) from  $1.1 \times 10^3$  to  $1.1 \times 10^9$  W m<sup>-2</sup> K<sup>-1</sup>.

To calculate the ETC of 10 vol.% GFs/Al composites using 2D image-based simulation, the orientations of GFs in the microstructures were corrected using Equation (9), and the TC of GFs was calculated using Equation (3-3). The TC of the Al

matrix ( $\lambda_{\text{Al-matrix}}$ ) was calculated using the following equation <sup>[15]</sup>:

$$\lambda_{\text{Al-matrix}} = \frac{1}{4} [\lambda_p(3v_p - 1) + \lambda_{\text{Al}}(3v_{\text{Al}} - 1) + \left( [\lambda_p(3v_p - 1) + \lambda_{\text{Al}}(3v_{\text{Al}} - 1)]^2 + 8\lambda_p\lambda_{\text{Al}} \right)^{\frac{1}{2}}] \dots\dots (3-10)$$

where the TCs of the raw material Al powder ( $\lambda_{\text{Al}}$ ) was taken as 228 W m<sup>-1</sup> K<sup>-1</sup> as in our previous works, TC of the pores ( $\lambda_p$ ) was 0.214 W m<sup>-1</sup> K<sup>-1</sup>, volume fraction of Al ( $v_{\text{Al}}$ ) was 0.994, and volume fraction of pores ( $v_p$ ) was 0.006. The pores in this study were assumed to be embedded in the Al matrix. The  $\lambda_{\text{Al-matrix}}$  was 226 W m<sup>-1</sup> K<sup>-1</sup>. Figure 3-9 shows the calculated ETC considering the interfacial heat conduction ( $h$ ) from 1.1 × 10<sup>3</sup> to 1.1 × 10<sup>9</sup> W m<sup>-2</sup> K<sup>-1</sup>. When the calculated ETC was equal to the experimental ETC,  $h$  was 1.8 × 10<sup>7</sup> W m<sup>-2</sup> K<sup>-1</sup>. This value (1.8 × 10<sup>7</sup> W m<sup>-2</sup> K<sup>-1</sup>) was the same order of magnitude as the  $h$  given by Xue, Zhou, and Schmidt (4.5 ~ 5 × 10<sup>7</sup> W m<sup>-2</sup> K<sup>-1</sup>) <sup>[2, 8, 16]</sup>. Moreover, the calculated ETC was 250 W m<sup>-2</sup> K<sup>-1</sup> when  $h$  was in 4.5 ~ 5 × 10<sup>7</sup> W m<sup>-2</sup> K<sup>-1</sup>.

To view the Al-GFs interface (see Figure 3-10), pores can be observed. Additionally, Li <sup>[4]</sup>, Zhou <sup>[5]</sup>, and Huang <sup>[17]</sup> reported that an amorphous layer interface could be readily formed at the Al-GFs interface. Therefore, the small difference between  $h$  calculated in this work and  $h$  reported by Xue, Zhou, and Schmidt can be attributed to the pores and amorphous layer at the interface.

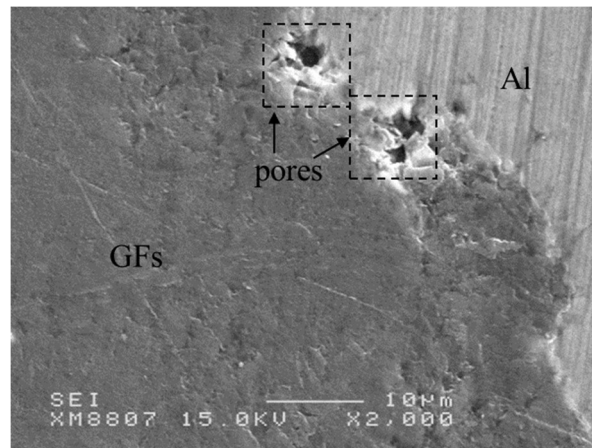


Figure 10. A SEM image of the microstructure at the Al-GFs interface.

### 3.4 Summary

This study aimed to calculate the ETC of GFs/Al composites using 2D image-based simulations. However, the ETC calculated from the microstructure images might not be equivalent to that measured using experimental methods. Because the 2D microstructure image cannot reveal the depth information related to the observation surface, the orientation of GFs in the 2D microstructure image differs from that of the experimental sample. Thus, we studied the relationships between the orientation of GFs in 2D cross-sectional images, the orientation of GFs in 3D models, and the aspect ratio of GFs in the 2D cross-sectional images.

The GFs in the modeled 3D models are elliptical cylinders. In light of the morphological characteristics of GFs, the conclusions of this study are limited to the study of the thermal conductivity of flake graphite reinforced aluminum matrix composites. GFs orientation was defined by the angle between GFs basal plane and heat flow direction in this study. Angles were marked as  $\theta_{2D}$  in the 2D cross-section and  $\theta_{3D}$  in the 3D model. Based on the research results, the following conclusions were drawn:

- 1)  $\theta_{2D}$  was generally larger than  $\theta_{3D}$ , and the difference between them could cause the TC error to be as high as  $840 \text{ W m}^{-1} \text{ K}^{-1}$ . No regularity was observed between the  $\theta_{2D}$  and  $\theta_{3D}$  differences.
- 2) The data for  $\theta_{2D}$ ,  $\theta_{3D}$ , and aspect ratio were on a curved surface, and the curved surface function can be used to convert  $\theta_{2D}$  into  $\theta_{3D}$ . The difference between the actual  $\theta_{3D}$  and  $\theta_{3D}$  values calculated using the curved surface function was less than  $6^\circ$ .
- 3) The ETC of the 10 vol.% GFs/Al composite was evaluated using 2D image-based simulation, where the suggested converting  $\theta_{2D}$  into  $\theta_{3D}$  was employed.
- 4) The calculated interfacial thermal conductance was  $1.8 \times 10^7 \text{ W m}^{-2} \text{ K}^{-1}$ , and it is the same order of magnitude with the reference value ( $4.5 \sim 5 \times 10^7 \text{ W m}^{-2} \text{ K}^{-1}$ ).

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## *Chapter 4*

# **Effect of the orientation of GFs and interfacial thermal resistance on the effective thermal conductivity of GFs/Al composites**

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## 4.1 Introduction

Graphite flakes (GFs) have excellent thermal properties, including high thermal conductivity (TC) and low coefficient of thermal expansion (CTE). In the basal plane, GFs TC and CTE are  $2800 \text{ W m}^{-1} \text{ K}^{-1}$  [1-4] and  $-1.5 \text{ ppm K}^{-1}$  [5, 6], respectively. Therefore, GFs/Al composites have attracted much attention from researchers of thermal management materials. Many experimental and theoretical studies have been conducted to study the effective thermal conductivities of GFs/Al composites. GFs/Al composites exhibit high effective thermal conductivities [7-20], e.g., Chen et al. [7] reported that the ETC of 80 vol.% GFs/Al composites ( $783 \text{ W m}^{-1} \text{ K}^{-1}$ ) was higher than that of 80 vol.% diamond/copper composites ( $724 \text{ W m}^{-1} \text{ K}^{-1}$ ). Li et al. [14] reported that the ETC of 70 vol.% GFs /Al composite was  $714 \text{ W m}^{-1} \text{ K}^{-1}$  in the plane parallel to the GFs layers.

However, the past studies [7, 11, 19, 21-23] also revealed that the ETCs of GFs/Al composites were affected by the orientation of GFs, interfacial thermal resistance, and interface reaction. The ETCs of GFs/Al composites are not as high as the expected theoretical values. The TCs of GFs are significant anisotropy, and its TC in the out-plane is  $38 \text{ W m}^{-1} \text{ K}^{-1}$  [3, 4, 12]. The anisotropic TC of GFs causes its orientation to influence the ETC of GFs/Al composites greatly. Moreover, it is almost impossible to control perfectly the orientation of GFs; thus, the orientations of GFs may be diverse in composites. On the other hand, the interfacial thermal resistance at the Al-GFs interface may be anisotropic due to the anisotropic crystal structure of GFs. Therefore, if the orientations of GFs in composites cannot be obtained, it would be difficult to calculate the effects of the orientation of GFs and the interfacial thermal resistance on the ETC using theoretical models, such as the rule of mixture [24, 25], Fricke model [26], Nan model and effective medium approximation (EMA) approach [27].

The interface reaction is a chemical reaction.  $\text{Al}_4\text{C}_3$  [10, 21, 28, 29] and amorphous layers [9, 14, 17] are easily formed at the Al-GFs interface.  $\text{Al}_4\text{C}_3$  phase is brittle and harmful for the bonding between Al and GFs. Meanwhile,  $\text{Al}_4\text{C}_3$  has low TC and can

hinder the heat transfer in the Al-GFs interface <sup>[7, 11, 21]</sup>. Fortunately, some approaches can suppress the interface reaction while fabricating the GFs/Al composites. The approaches include lowering the sintering temperature <sup>[7, 13, 17, 19]</sup> or reducing the reaction time <sup>[9, 21, 30, 31]</sup>, or adding Si or SiC particles to the starting materials <sup>[8, 9, 12, 13, 15, 19]</sup>, etc. Besides, the amorphous layer at the Al-GFs interface may act as the nucleation site for Al<sub>4</sub>C<sub>3</sub> <sup>[32]</sup>. The amorphous layer thus can also hinder the heat transfer itself <sup>[9, 33]</sup>.

Therefore, in this chapter, 10~20 vol.% GFs/Al composites were prepared. two-dimensional(2D) image-based simulation method <sup>[34]</sup> is used to calculate the ETC of GFs/Al composites considering the effect of the orientations of GFs and the interfacial thermal resistance. In addition, the interfacial thermal resistances in the in-plane and the out-of-plane of GFs were calculated.

## 4.2 Experimental and calculation procedures

### 4.2.1 Experiments

Al powder (99.9%, 30  $\mu\text{m}$  in diameter) and GFs (98%, average particle size: 137.02  $\mu\text{m}$ ) were used. The nominal TC of the Al powder was  $236 \text{ W m}^{-1} \text{ K}^{-1}$ , while that of GFs was 880 and  $38 \text{ W m}^{-1} \text{ K}^{-1}$  at the basal-plane and out-of-plane directions, respectively. Al powders were mixed with 10 vol.% and 20 vol.% of GFs, respectively. Subsequently, appropriate amounts of alumina balls were added through a wet process for 2 h using a V-type mixer at 50 rpm. Then, spark plasma sintering (SPS) was conducted at 873 K and 60 MPa pressure for 0.5 h under a vacuum of  $1.3 \times 10^{-2} \text{ Pa}$ . In order to avoid the interface reaction between Al and GF, the sintering temperature was maintained lower than the melting point of Al. The temperature was raised by 200 K/min: one pure Al sample, 10 vol.% GF/Al (Samples 1-2), and 20 vol.% GF/Al (Samples 3-4) composites were fabricated. Figure 4-1 depicts the shape of the samples sintered using SPS.

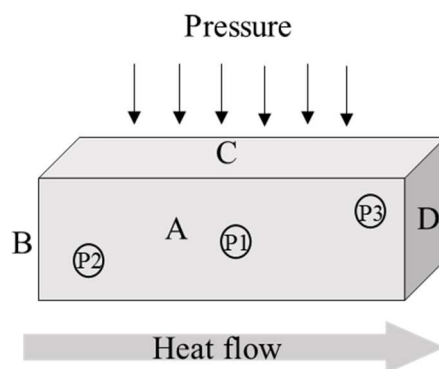


Figure 4-1. Shape of the sample sintered by SPS; P1, P2, and P3 at the A plane denote the regions that correspond to the optical micrographs. The grey arrow illustrates that the heat flow was from B to D during ETC measurement, C is the plane bearing the pressure during sintering.

The microstructure images of the sintered samples were obtained via optical

microscopy from the P1, P2, and P3 regions on A plane in Figure 4-1. The relative densities of the sintered samples were measured using the Archimedes method. ETCs of the sintered samples were measured at 25 °C using a steady-state thermal-conductivity device.

## 4.2.2 Numerical methods

2D image-based simulations were performed in two steps. First, the ETCs of composites were calculated considering the effect of the orientations of GFs. Then, ETCs of the composites were calculated considering the orientations of the GFs and the interfacial thermal resistance at the Al-GFs interface. The obtained microstructure images of the experimental samples were used to perform 2D image-based simulations.

### 4.2.2.1 Temperature distribution

We calculated 2D-temperature distributions using the finite volume method as follows:

$$T_{x,y}^{n+1} = T_{x,y}^n + \frac{\Delta t}{\rho c} \left( \frac{q_{x+1,y}^n - q_{x-1,y}^n}{\Delta x} + \frac{q_{x,y+1}^n - q_{x,y-1}^n}{\Delta y} \right) \dots\dots (4-1)$$

where  $x$  and  $y$  denote the coordinate positions of the element,  $T_{x,y}^n$  is the temperature of element at  $(x, y)$  coordinates, and  $T_{x,y}^{n+1}$  is the element's temperature at  $(x, y)$  coordinates after a time  $\Delta t$ ,  $\rho$  is density,  $c$  is specific heat,  $\Delta x$  and  $\Delta y$  denote the sizes of elements along the  $x$ -axis and the  $y$ -axis, respectively,  $q^n$  denotes heat flow. When  $q^n$  is the heat conduction between adjacent elements (e.g., as depicted in Figure 4-2,  $q^n$  moves from the E to adjacent elements M or N),  $q^n$  can be calculated as follows:

$$q_{x+1,y}^n = \lambda_{x+\frac{1}{2},y} \left( \frac{T_{x+1,y}^n - T_{x,y}^n}{\Delta x} \right), \quad q_{x-1,y}^n = \lambda_{x-\frac{1}{2},y} \left( \frac{T_{x,y}^n - T_{x-1,y}^n}{\Delta x} \right)$$

$$q_{x,y+1}^n = \lambda_{x,y+\frac{1}{2}} \left( \frac{T_{x,y+1}^n - T_{x,y}^n}{\Delta y} \right), \quad q_{x,y-1}^n = \lambda_{x,y-\frac{1}{2}} \left( \frac{T_{x,y}^n - T_{x,y-1}^n}{\Delta y} \right)$$

$$\lambda_{x+\frac{1}{2},y} = \frac{2\lambda_{x,y}\lambda_{x+1,y}}{\lambda_{x,y} + \lambda_{x+1,y}}, \quad \lambda_{x-\frac{1}{2},y} = \frac{2\lambda_{x,y}\lambda_{x-1,y}}{\lambda_{x,y} + \lambda_{x-1,y}}$$

$$\lambda_{x,y+\frac{1}{2}} = \frac{2\lambda_{x,y}\lambda_{x,y+1}}{\lambda_{x,y} + \lambda_{x,y+1}}, \quad \lambda_{x,y-\frac{1}{2}} = \frac{2\lambda_{x,y}\lambda_{x,y-1}}{\lambda_{x,y} + \lambda_{x,y-1}}$$

..... (4-2)

where  $\lambda_{(x,y)}$  denotes the TC of element at  $(x,y)$  coordinates,  $\lambda_{(x+1,y)}$ ,  $\lambda_{(x-1,y)}$ ,  $\lambda_{(x,y+1)}$ , and  $\lambda_{(x,y-1)}$  denote the TC of elements adjacent to  $(x,y)$ , and  $\lambda_{(x+1/2,y)}$ ,  $\lambda_{(x-1/2,y)}$ ,  $\lambda_{(x,y+1/2)}$ , and  $\lambda_{(x,y-1/2)}$  denote the harmonic mean of TC.

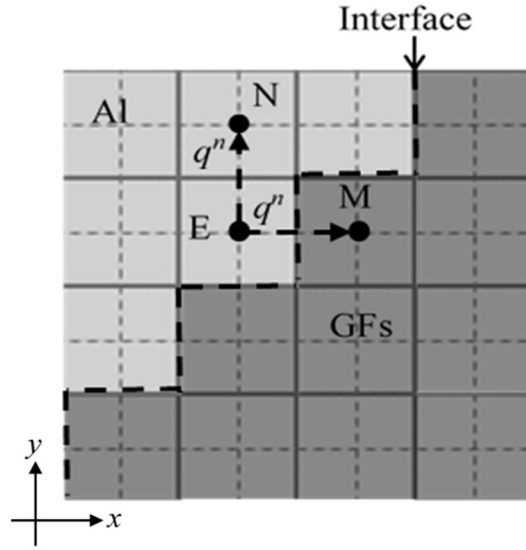


Figure 4-2. Schematic diagram of heat flow at the interface. The dashed line denotes the Al–GF interface,  $q^n$  denotes the heat flow in the direction of the dashed arrows.

Considering the effect of the interfacial thermal resistance in the composite, if  $q^n$  moves by heat transfer between different types of elements, the interfacial heat transfer coefficient ( $h$ ) must be considered. For example, as depicted in Figure 4-2, considering that  $q^n$  moves from E to M and through the interface between Al and GFs (Al-GF),  $q^n$  can be calculated as follows:

$$\begin{aligned}
 q_{x+1,y}^n &= h(T_{x+1,y}^n - T_{x,y}^n), & q_{x-1,y}^n &= h(T_{x,y}^n - T_{x-1,y}^n) \\
 q_{x,y+1}^n &= h(T_{x,y+1}^n - T_{x,y}^n), & q_{x,y-1}^n &= h(T_{x,y}^n - T_{x,y-1}^n) \\
 & & & \dots\dots (4-3)
 \end{aligned}$$

#### 4.2.2.2 Calculation of effective thermal conductivity

As depicted in Figure 4-3, the simulation model is comprised of two heat sources and composite components. The top and bottom surfaces correspond to the periodic boundary, while the left and right sides correspond to the adiabatic boundary. The composite part is based on the microstructure images of Samples 1–4. Therefore, the size of the composite part ( $N_x \times N_y$ ) was  $450 \times 600$  elements for Samples 1, 2, and 4, and it was  $570 \times 450$  elements for Sample 3. The heat source size was  $5 \times 600$  elements for Samples 1, 2, and 4, and it was  $5 \times 450$  elements for sample 3 ( $N_L = N_R = 5$  elements). The size of each element was  $1.18 \times 10^{-6}$  m. The initial TC of the GFs was set to be  $(880, 38)$  W m<sup>-1</sup> K<sup>-1</sup> in the  $(x, y)$  directions.

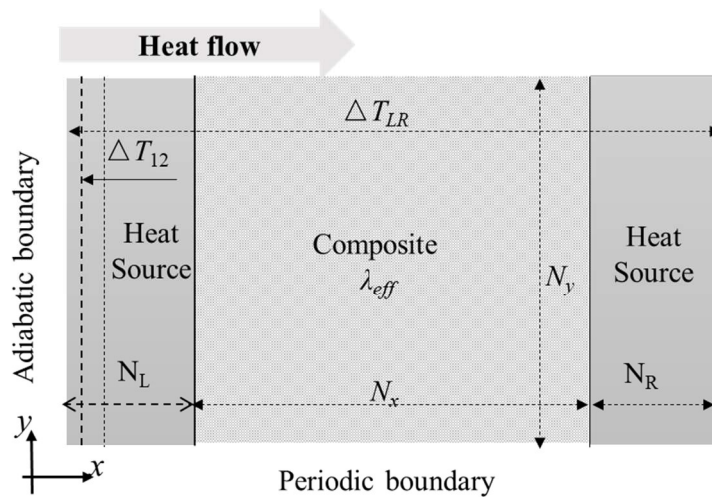


Figure 4-3 Simulation model for the ETC calculation.

The temperature of the left edge elements was set to 301 K, and the initial

temperature of the other elements was set to 300 K. The temperatures at the left and right sides were fixed. The temperature of the remaining elements was iterated until the temperature variation was lower than  $10^{-13}$  K, and the temperature distribution in the steady-state was obtained. The ETC values of the GF/Al composites,  $\lambda_{s-eff}$ , were calculated when the temperature distribution was in the steady-state, as follows:

$$\lambda_{s-eff} = \frac{\lambda_{Al-e} \Delta T_{12} N_x}{\Delta T_{LR} - N_L \Delta T_{12} - N_R \Delta T_{12}} \dots\dots (4-4)$$

where  $\lambda_{Al-eff}$  denotes the TC of the Al matrix considering the effect of pores, and  $\Delta T_{12}$  is the average temperature difference between the first and the second columns;  $N_L$  and  $N_R$  denote the numbers of elements in the heat sources,  $N_x$  and  $N_y$  are the numbers of elements in the composite region.

#### 4.2.2.3 Thermal conductivity of GFs

This study used an angle ( $\theta$ ) between the basal plane of GFs and the direction of heat flow to characterize the GFs' orientation in composites. Thus, by using the following equation (4-5)<sup>[18]</sup>, the TC of GFs can be calculated along the direction of the basal plane of GFs parallel ( $\lambda_{//}$ ) and perpendicular ( $\lambda_{\perp}$ ) to the heat flow direction.

$$\begin{aligned} \lambda_{//} &= \lambda_a \left[ 1 - \left( 1 - \frac{\lambda_c}{\lambda_a} \right) \sin^2 \theta \right] \\ \lambda_{\perp} &= \lambda_a \left[ 1 - \left( 1 - \frac{\lambda_c}{\lambda_a} \right) \cos^2 \theta \right] \end{aligned} \dots\dots (4-5)$$

where  $\lambda_a$  and  $\lambda_c$  are the TC in the basal-plane and the out-of-plane directions of GFs, respectively.

Figure 4-4 illustrates the relationship between the TC of GFs and the angle  $\theta$ . It can be observed that TC decreased significantly from 880 to 38 W m<sup>-1</sup> K<sup>-1</sup> when the angle increased from 0° to 90°, revealing that the orientation of the GFs significantly

affects the ETC of composites. In addition, Figure 4-4 shows that the TC of GFs was lower than that of the Al matrix ( $236 \text{ W m}^{-1} \text{ K}^{-1}$ ) if the angles were greater than  $61^\circ$ , indicating that the ETC of composites could be enhanced at critical angles smaller than or equal to  $61^\circ$ .

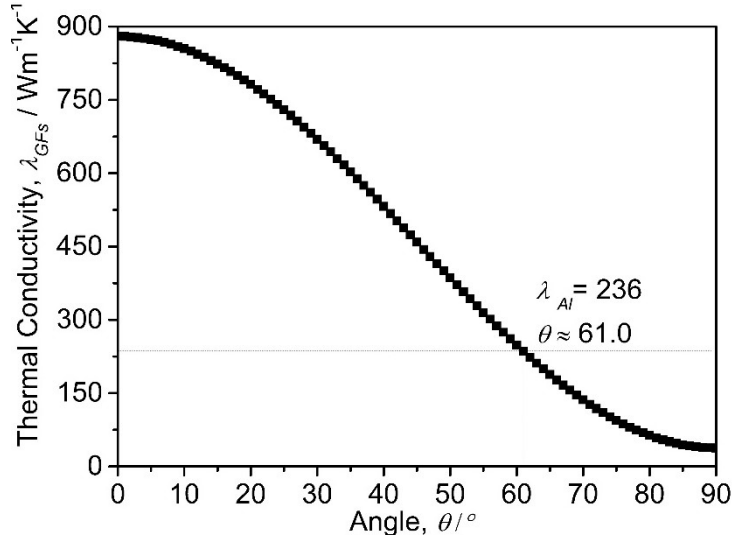


Figure 4-4 TC of GFs as a function of angle ( $\theta$ ). The nominal TC of Al matrix,  $\lambda_{Al}$ , was  $236 \text{ W m}^{-1} \text{ K}^{-1}$  (gray horizontal dotted line).

#### 4.2.2.4 Interfacial thermal resistance

The interfacial thermal conductance in the in-plane and the out-of-plane may be different due to the anisotropic crystal structure of GFs. Thus, the interfacial thermal conductance,  $h$ , at the Al-GFs interface was calculated by the DMM method.  $h$  can be calculated as follows:

$$h = \frac{\hbar^2}{8\pi^2 K T^2} \sum_j \left[ \frac{1}{v_{1,j}^2 v_{2,j}^2 (v_{1,j}^{-2} + v_{2,j}^{-2})} v_{1,j} \int_0^{\omega_{1,j}^c} \frac{\omega^4 e^{\frac{\hbar\omega}{KT}}}{\left(e^{\frac{\hbar\omega}{KT}} - 1\right)^2} d\omega \right] \dots (4-6)$$

where  $\hbar$  is Planck's constant,  $K$  is Boltzmann constant,  $T$  is temperature,  $\omega$  is phonon angular frequency, and  $\omega_{1,j}^c$  is the maximum phone angular frequency. The phono sound velocity,  $v$ , can be written as:



$$\sum_j v_j = v_L + 2v_T$$

$$v_L = \sqrt{\frac{E}{\rho} \cdot \frac{(1 - \mu)}{(1 + \mu)(1 - 2\mu)}}$$

$$v_T = \sqrt{\frac{E}{\rho} \cdot \frac{1}{2(1 + \mu)}}$$

..... (4-7)

where  $v_L$  is longitudinal sound velocity,  $v_T$  is transverse sound velocity,  $E$  is Young's modulus,  $\rho$  is density,  $\mu$  Poisson's ratio. The maximum phone angular frequency is calculated:

$$\omega_{1,j}^c = v_{1,j} \left( 6\pi^2 \frac{\rho_1 N_A}{M_1} \right)^{\frac{1}{3}}$$

..... (4-8)

where  $N_A$  is Avogadro constant,  $M_1$  is relative molecular mass. The material parameters are given in Table 4-3.

## **4.3 Results and discussion**

### **4.3.1 Microstructures**

Figure 4-5 illustrates the microstructures of the samples. The GFs were homogeneously distributed in the Al matrix. However, GFs' orientations were not uniform in the samples. As shown in Figures 4-5, the orientations of GFs were almost parallel to each other in Samples 1 and 4. However, GFs' orientations were partially disordered in Samples 2 and 3, forming different angles between GFs' basal-plane and heat flow direction. The orientations of GFs in Figures 4-5 were determined and shown in Figures 4-6. As depicted in Figures 4-6 (a), GFs were bounded by the smallest circumscribed rectangle, and the angle  $\theta$  represents the orientation of GFs. The  $\theta$  values in the microstructures were measured as depicted in Figure 4-6 (b), excluding GFs smaller than  $5 \times 5$  pixels. The measured  $|\theta|$  values were shown in Figure 4-6 (c)–(f). Most of the angles ranged from  $0^\circ$  to  $30^\circ$  in all the samples. Nevertheless, for the angles ranging between  $50^\circ$  and  $90^\circ$ , the relative frequency in Samples 1 and 4 was  $< 10\%$ , while it was  $> 10\%$  in Samples 2 and 3. Additionally, as depicted in Figures 4-4, the TC of GFs along the heat flow direction was smaller than that of the Al matrix when the angle was  $> 61.0^\circ$ . Thus, it can be derived from Figure 4-6 that the average TCs of GFs in Samples 2 and 3 along the heat flow direction were smaller than the average TCs of GFs in Samples 1 and 4. Moreover, the average angle was calculated for each sample, as shown in Figure 4-6 (c)–(f).

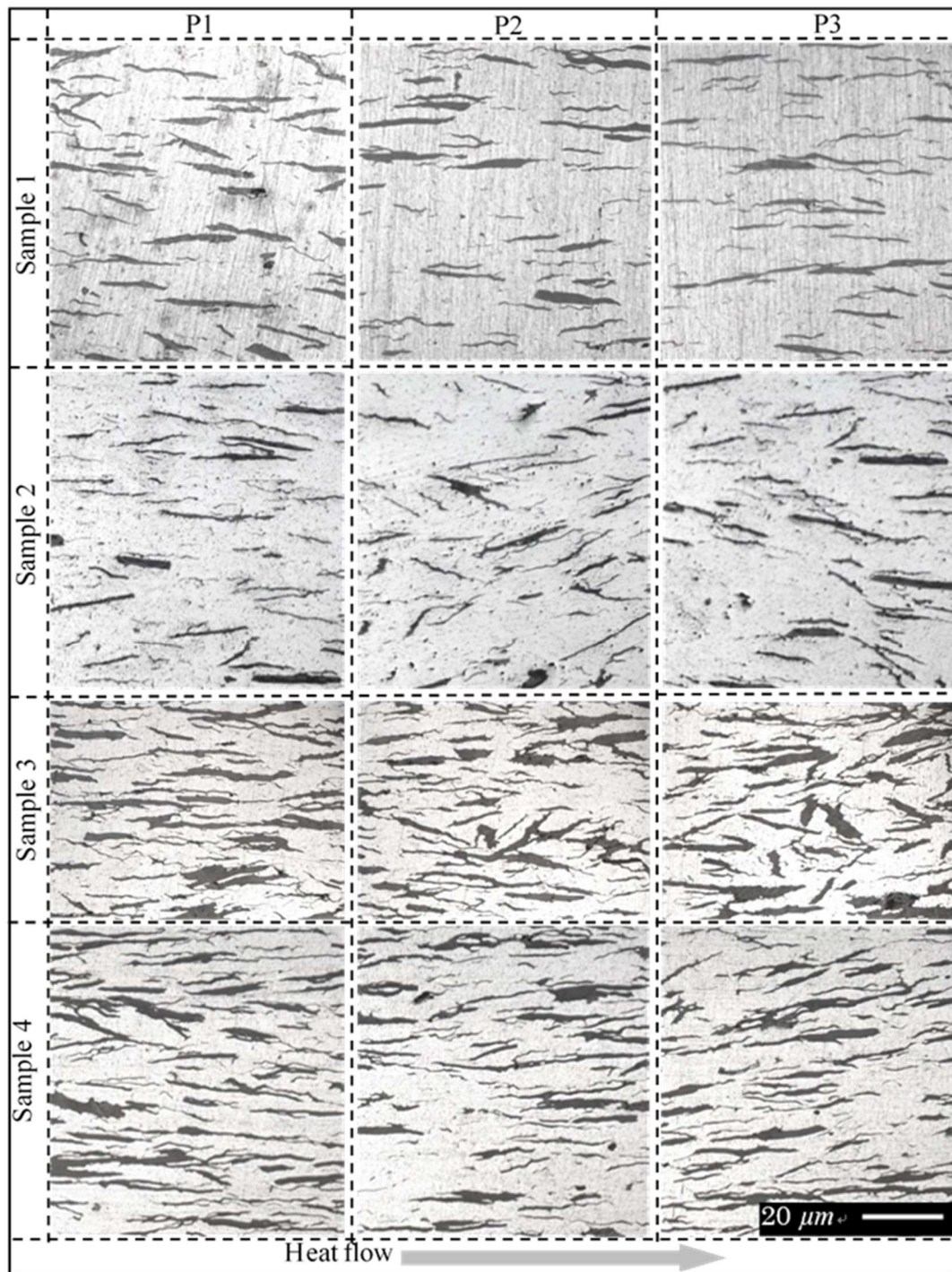


Figure 4-5 Optical micrographs of Samples 1–4. Regions P1, P2, and P3 correspond to the black circles at the A-plane in Figure 1. The gray arrow represents the heat flow direction during the ETC measurements. The brighter phases correspond to the Al matrix and the elongated dark phases to the GFs.

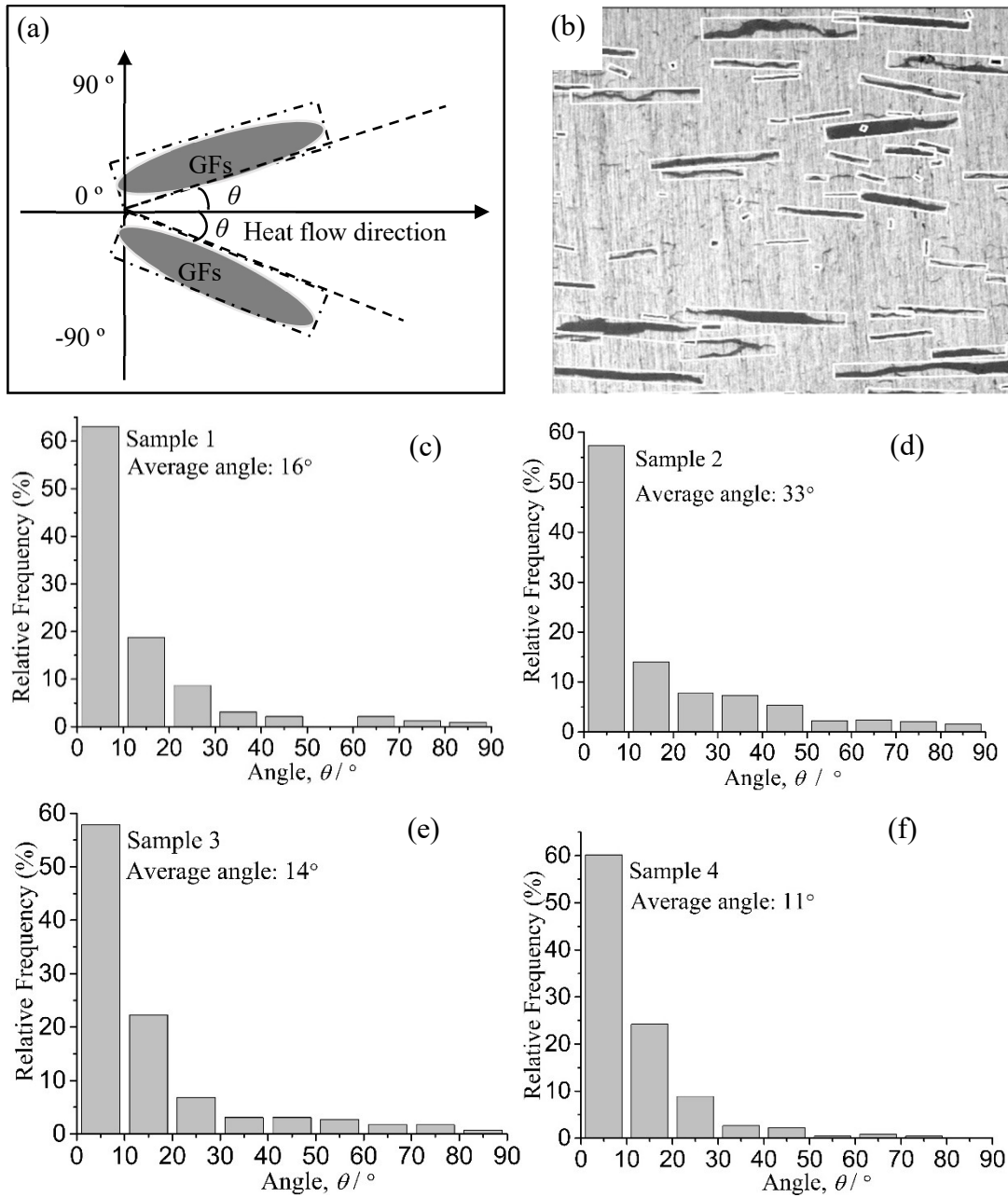


Figure 4-6 Distribution of the GF orientations in the Al matrix. Gray and white rectangles in (a) and (b) represent the smallest circumscribed rectangles and were used to measure the angle ( $\theta$ ) between the basal-plane of GFs and the heat flow direction; in case  $-90^\circ < \theta < 0^\circ$ ,  $\theta$  was taken as the absolute value.

### 4.3.2 Effect of porosity on TC of Al

The measured ETCs and the relative densities of the experimental samples were listed in Table 4-1. The ETC of Sample 1 was higher than that of Sample 2, and the ETC of Sample 3 was lower than that of Sample 4. These measurement results agree with the results shown in Figures 4-6.

Table 4-1 Volume fraction of GFs ( $v_{GFs}$ ), measured ETCs ( $\lambda_{eff}$ ), and relative densities of Samples 1–4 and pure Al sample.

Sample	$v_{GFs}$	$\lambda_{eff}$	Relative density
	%	$Wm^{-1}K^{-1}$	
1	10	248	99.2%
2	10	238	99.4%
3	20	273	98.0%
4	20	280	99.3%
Al	0	217	97.0%

The measured relative density (Table 4-1) indicated that some pores existed in the composites. Therefore, in this study, the pores for all the samples were assumed to be in the Al matrix. The following equation derived by Landauer<sup>[35]</sup> was used to eliminate the pores' effect on the TC of Al-matrix:

$$\lambda_{Al-eff} = \frac{1}{4} [\lambda_p(3v_p - 1) + \lambda_{Al}(3v_{Al} - 1) + \left( [\lambda_p(3v_p - 1) + \lambda_{Al}(3v_{Al} - 1)]^2 + 8\lambda_p\lambda_{Al} \right)^{\frac{1}{2}}] \dots\dots (4-9)$$

where  $\lambda_{Al-eff}$  denotes the TC of Al matrix,  $\lambda_{Al}$  and  $\lambda_p$  are the TCs of the raw material Al powder and pores, respectively,  $v_{Al}$  is the volume fraction of Al, and  $v_p$  is the volume fraction of the pores. In addition,  $v_{Al}$ ,  $v_p$ , and  $\lambda_p$  are listed in Table 4-2. First, the

measured TC of the pure Al sample ( $\lambda_{Al-eff}$ , 217 W m<sup>-1</sup> K<sup>-1</sup>) and its relative density (97%) were substituted in equation (4-9), and  $\lambda_{Al}$  was determined to be 228 W m<sup>-1</sup> K<sup>-1</sup>, as shown in Table 4-2. Second, the TCs of the Al-matrix for Samples 1–4,  $\lambda_{Al-eff}$ , were calculated by substituting  $v_{Al}$ ,  $v_p$ ,  $\lambda_p$ , and  $\lambda_{Al}$  from Table 4-2 in equation (4-9). The calculated values of  $\lambda_{Al-eff}$  are 225.3, 225.9, 221.2, and 225.6 W m<sup>-1</sup> K<sup>-1</sup> for Samples 1–4, respectively. The ETCs ( $\lambda_{ROM}$ ) of Samples 1–4 were calculated using the rule of mixture (ROM), i.e.,  $\lambda_{ROM} = (v_{GFs})(\lambda_{GFs}) + (1-v_{GFs})(\lambda_{Al-eff})$ , where the orientations of all the GFs were assumed to be aligned with the heat flow direction in the composites. The TC of the GFs ( $\lambda_{GFs}$ ) was set to 880 W m<sup>-1</sup> K<sup>-1</sup> along the heat flow direction. The values of  $\lambda_{ROM}$  were 291, 291, 353, and 357 Wm<sup>-1</sup>K<sup>-1</sup> for Samples 1–4, respectively. The calculated  $\lambda_{ROM}$  was considerably higher than the measured ETC. Furthermore, We performed 2D image-based simulations to confirm the effect of the orientation of GFs on the ETC of composites.

Table 4-2 Calculation parameters in equation (10) and calculated effective thermal conductivities of composites.

Sample	$v_{Al}$	$v_p$	$\lambda_p$	$\lambda_{Al}$	$\lambda_{Al-eff}$	$\lambda_{ROM}$	$\lambda_{s-eff}$
			Wm <sup>-1</sup> K <sup>-1</sup>				
Al	0.97	0.03	0.214	228	217		
1	0.992	0.008			225.3	291	270
2	0.994	0.006	0.214		225.9	291	263
3	0.98	0.02			221.2	353	318
4	0.993	0.007			225.6	357	323

### 4.3.3 Effect of the orientation of GFs on the effective thermal conductivity

The TC values of GFs in the experimental samples were calculated using equation

(4-5), and Figure 4-7 illustrated the results. GFs were denoted in different colors according to the orientations of GFs in the Al matrix. When the orientation of GFs was parallel to the heat flow direction, GFs exhibited high TC, and they were marked in red.

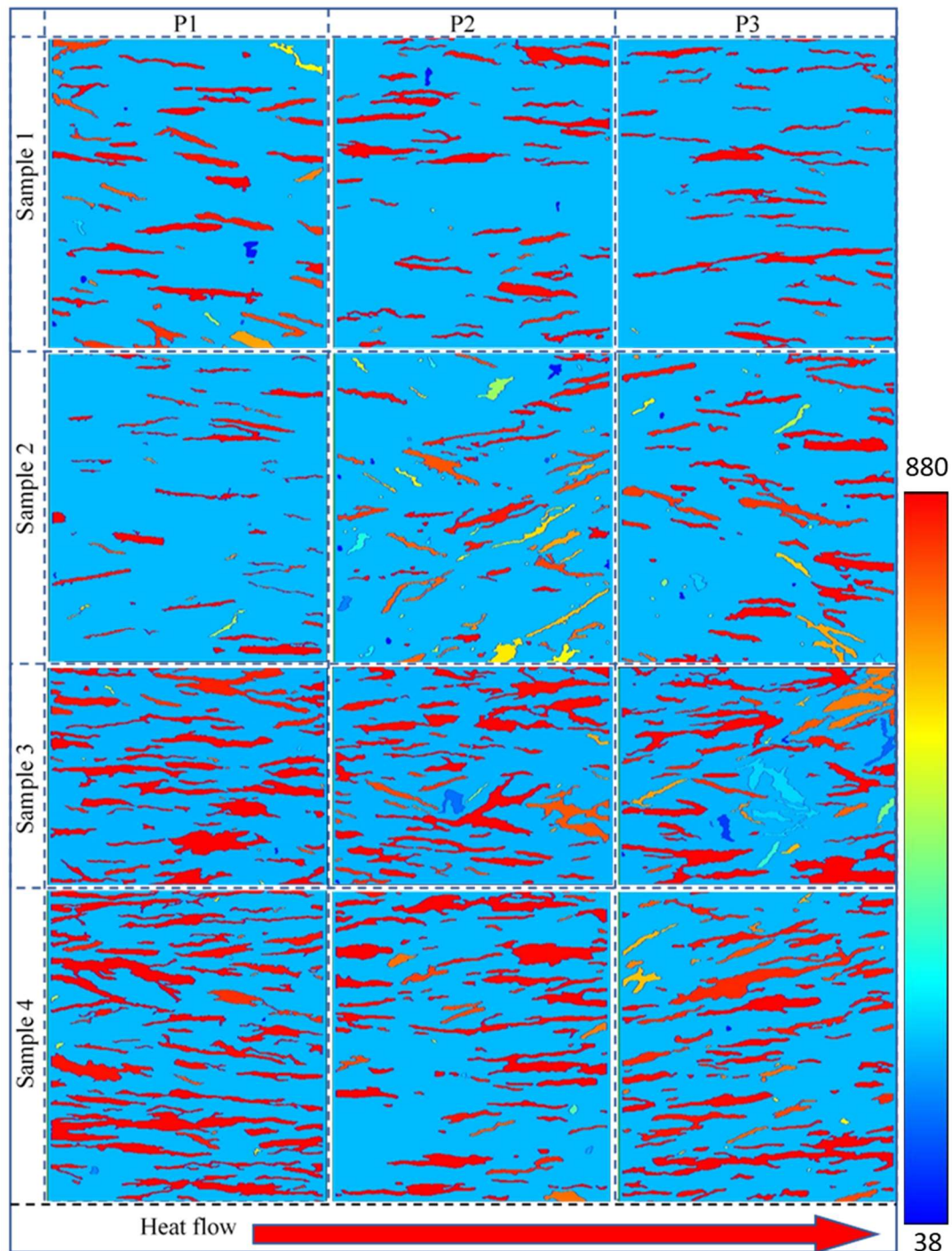


Figure 4-7 TC values of GFs in Samples1-4. Regions denoted as P1, P2, and P3 correspond to the microstructures in Figure. 5; The color bar shows the TCs of GFs in different orientations; The red arrow indicates the heat flow direction.

In contrast, if the GFs tilted from the heat flow direction, color changes indicated a decrease in TC. Figure 4-7 showed that most GFs were colored in red for Samples 1 and 4, which implied that GFs had a high TC, while the yellow, blue, and other colors appeared in the P2 and P3 regions for Samples 2 and 3. Table 4-2 listed the simulated ETCs ( $\lambda_{s-eff}$ ) based on 2D image-based simulations. Figure 4-8 further compared the measured ETC,  $\lambda_{ROM}$ , and  $\lambda_{s-eff}$  of Samples 1–4. The values of  $\lambda_{s-eff}$  were smaller than those of  $\lambda_{ROM}$ , and the deviations between  $\lambda_{ROM}$  and  $\lambda_{s-eff}$  for Samples 1–4 were 7.2%, 9.6%, 9.9%, and 9.5%, respectively. Comparing the ETC loss in Samples 1–4, the difference between Samples 1 and 2 in ETC loss was 2.4%, while it was only 0.4% between Samples 3 and 4. Those deviations may be explained by the angle deviations depicted in Figure 4-6 (c)-(f), where the angle difference between samples 1 and 2 was 17°, and between samples 3 and 4 was 3°. Furthermore, compared the  $\lambda_{s-eff}$  values and the measured ETC values, the deviations between  $\lambda_{s-eff}$  and the measured ETCs were 8.1%, 9.5%, 14.2%, and 13.3% for Samples 1–4, respectively. These deviations can be attributed to the effect of the interfacial thermal resistance between Al and GFs.

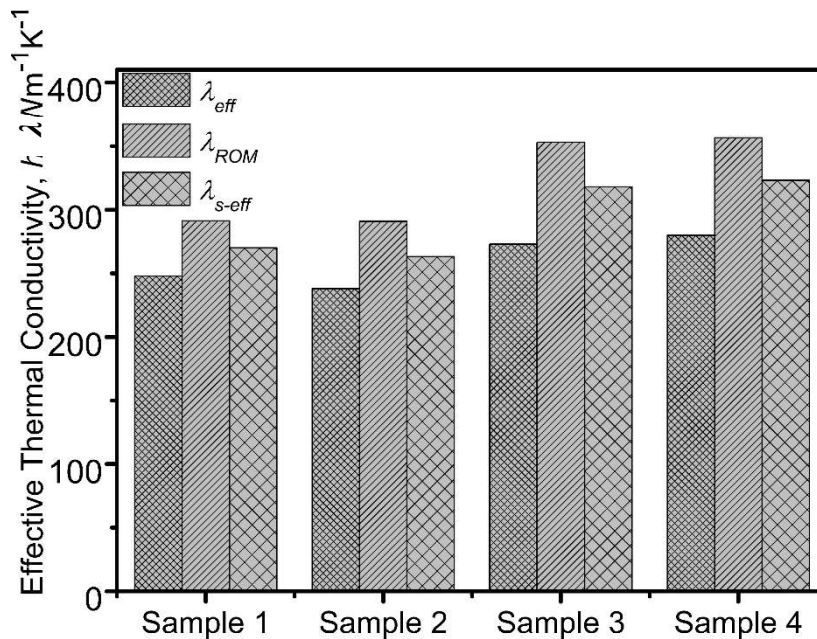


Figure 4-8 ETCs of the composite samples.  $\lambda_{eff}$  denotes the measured ETC,  $\lambda_{ROM}$  is the ETC calculated using the rule of mixture, and  $\lambda_{s-eff}$  is the ETC calculated using the 2D image-based simulation considering GF orientations.



### 4.3.4 Effect of interfacial thermal resistance on the effective thermal conductivity

A series of 2D image-based simulations were performed to study further the effect of the interfacial thermal resistance on the ETC of the composites. The heat transfer coefficient,  $h$ , ranged from  $10^3$  to  $10^9$   $\text{W m}^{-2} \text{K}^{-1}$  at the Al–GF interface. Figure 4-9 illustrated ETC values as a function of  $h$  for composites.

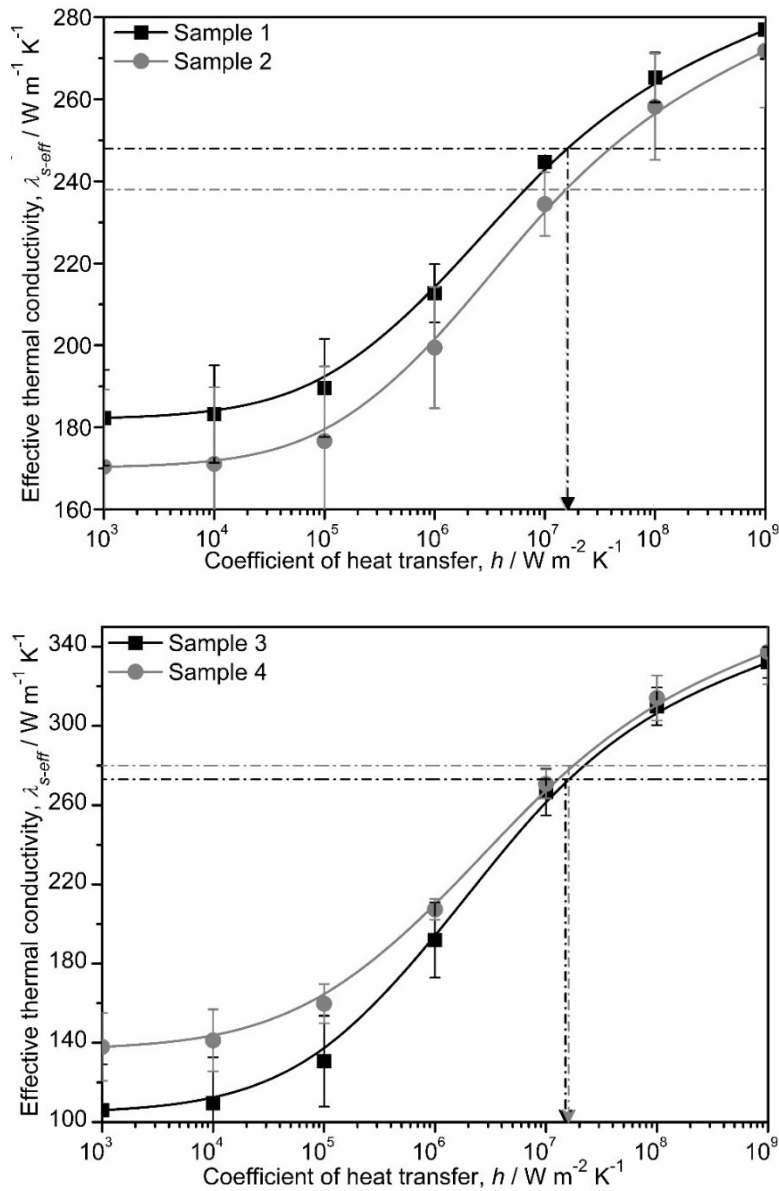


Figure 4-9 ETC as a function of the heat transfer coefficient,  $h$ . The dashed lines represent the experimental ETCs and the arrows denote the  $h$  value at the Al-GF interface.

The ETC curve was similar to KJMA <sup>[36]</sup> equation, and ETC values rapidly increased as  $h$  increased from  $10^5$  to  $10^9$ . Moreover, the ETC curves for Sample 1 and Sample 2 did not overlap in the  $h$  value range between  $10^3$  and  $10^9$ , further proving that the orientations of GFs affected the composite ETC significantly. In Figure 4-9, the dashed lines represented the experimental ETC, and the arrows denoted the heat transfer coefficients ( $h$ ) at the Al-GF interface. Table 4-3 showed the  $h$  values evaluated via inverse analysis, and the values showed the same order of magnitude for all samples, i.e.,  $10^7$ . Moreover, a reference value of  $h$  calculated using the AMM model was listed in Table 4-3. The order of magnitude of the reference value is also  $10^7$ , and the minor differences between the  $h$  reference value and the simulated  $h$  for Samples 1–4 can be attributed to the misfits at the interface. However, the values of  $h$  calculated by DMM were  $1.1 \times 10^8 \text{ Wm}^{-2}\text{K}^{-1}$  in the in-plane and  $1.32 \times 10^9 \text{ Wm}^{-2}\text{K}^{-1}$  in the out-of-plane. Thus, the values were greater than the simulated  $h$ .

Table 4-3 Material parameters of the DMM calculation<sup>[37]</sup>, interfacial thermal conductance ( $h$ ), and interfacial thermal resistance ( $R$ ); Sim. is the simulated  $h$  using 2D-image, Cal. is calculated  $h$  using DMM method, Ref. is the reference value of  $h$ .

Sample	$T$ (K)	$E$ (GPa)	$P$ (g/cm <sup>3</sup> )	$\mu$	$M_l$	$h / 10^7 \text{ Wm}^{-2}\text{K}^{-1}$			$R / 10^{-8}$ m <sup>2</sup> KW <sup>-1</sup>
						Sim.	Cal.	Ref.	
Al	/	70.0	2.7	0.3	27				
GFs	In-plane	300	1153.0	2.2	0.195	12	11	4.8 <sup>[19]</sup>	
	Out-of-plane		39.511	2.2	0.0002	12	130		
1						1.62			6.17
2						1.56			6.41
3						1.64			6.10
4						1.78			5.62

The interfacial thermal resistance,  $R$ , is the reciprocal of heat transfer coefficient, i.e.,  $R = 1 / h$ ,  $R$  values for Samples 1–4 were listed in Table 4-3. It should be noted that  $R$  values were sufficiently small and even negligible. However, as shown in Figure 4-8,  $R$  affected the ETC of the composites significantly, leading to a significant decrease in ETC. The reason can be attributed to the considerable number of interfaces formed in the GF/Al composites.

#### **4.4 Summary**

10 vol.% GF/Al and 20 vol.% GF/Al composites were fabricated via SPS. The microstructures images of the composites were observed, and their relative density and ETCs were measured. The measured ETC values were 248, 238, 273, and 280 W m<sup>-1</sup> K<sup>-1</sup> for Samples 1–4, which were smaller than the ETCs calculated through the *ROM* (291, 291, 353, and 357 W m<sup>-1</sup> K<sup>-1</sup>, for Samples 1–4, respectively).

The average angles of the GFs with respect to the heat flow direction were calculated to be 16°, 33°, 14°, and 11° for Samples 1–4, respectively. The ETCs of composites considering the orientations of GFs were evaluated using 2D image-based simulations. Compared with the ETC calculated using *ROM*, the results showed that ETCs decreased due to the anisotropic TC of GFs by 7.2%, 9.6%, 9.9%, and 9.5% for Samples 1–4, respectively. Due to the interfacial thermal resistance effect, the ETCs decreased by 8.1, 9.5, 14.2, and 13.3% for Samples 1–4, respectively. The *R* values at the Al–GF interface were evaluated to be  $6.17 \times 10^{-8}$ ,  $6.41 \times 10^{-8}$ ,  $6.10 \times 10^{-8}$ , and  $5.62 \times 10^{-8}$  W m<sup>-2</sup> K<sup>-1</sup> for Samples 1–4, respectively.

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## *Chapter 5*

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### **Conclusions**

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Graphite flake reinforced aluminum matrix composites (GFs/Al composite) are lightweight, low cost, and exhibit outstanding thermal conductivity and low coefficient of thermal expansion. Especially for the excellent thermal conductivity, GFs/Al composites have been proposed as one of the most promising electronic packaging materials.

Lots of studies have been carried out to explore the thermal properties of GFs/Al composites. However, past studies reported that the high thermal conductivity of GFs could not be fully displayed in the GFs/Al composites. The effective thermal conductivity of GFs/Al composites was influenced by the volume fraction, dispersion, size, orientation of GFs, interface reaction, and interface thermal resistance, in particular the orientation of GFs and the interfacial thermal resistance. Unfortunately, until today, no method has been proposed to control the orientation of GFs well, and the theoretical models such as mixture rules, Maxwell, Fricke, Every, Nielsen, and the effective medium approximation cannot be considered the effect of the orientation of GFs. Moreover, interface thermal resistance seems to be an inherent property of carbon materials reinforced metal-matrix composites, and the value of interface thermal resistance is affected by the structure, composition, and state of the Al-GFs interface. AMM and DMM models are commonly used to calculate interface thermal resistance. However, the AMM and DMM models calculate the interface thermal resistance by assuming an ideal interface, but in reality, there may be gaps and lattice defects at the interface. Therefore, the interface thermal resistance evaluated by the AMM and DMM methods is not completely reliable.

This study aims to study the effective thermal conductivity of GFs/Al composites

by two-dimensional image simulations, considering the orientation of GFs and the interface thermal resistance. The works were performed in three steps: first, 10-40 vol.% GFs/Al composites were prepared by spark plasma sintering. Moreover, we investigated the relative density of composites, orientation and distribution of GFs, interface reaction between Al and GFs, effective thermal conductivity, and thermal expansion coefficient of GFs/Al composites. Second, since the orientation of GFs in two-dimensional images may differ from that of GFs in three-dimensional images, some three- and two-dimensional models of GFs/Al composites were created to investigate the relationships between the orientations of GFs in three- and two-dimensional models. As a result, a fitted function was obtained to convert the two-dimensional orientation to three-dimensional orientation. Third, the effects of GFs' orientation and the interfacial thermal resistance on the effective thermal conductivities of GFs/Al composites were calculated by two-dimensional image simulations, and the interface thermal resistances at the Al-GFs interface were calculated by inverse analysis, AMM, and DMM methods. Through the above research, this thesis draws the following conclusions:

- 1) 10-40 vol.% GFs/Al composites were fabricated by spark plasma sintering; Most of the GFs were stacked by layers and parallel to each other in the composites; With GFs content up to 40%, GFs exhibit a state of aggregation and tends to form the network by connecting GFs; The formation of  $Al_4C_3$  was absent in the 10-40 vol.% GFs/Al composites; The relative density was up to 99% for the GFs/Al composites; With GFs content up to 40%, the relative density of composites decreased to 97.1%. The measured effective thermal conductivities were 248, 280, 313, and 346  $W\ m^{-1}\ K^{-1}$  for 10-40 vol.% GFs/Al composites, respectively; Due to the orientation of GFs and the interface thermal resistance, the experimental values of effective thermal conductivity of the GFs/Al composites were lower than the theoretical values. The experimental values of the coefficient of thermal expansion were 26.2, 24.4, 23.4, and 22.1 ppm/K for 10-40 vol.% GFs/Al composites, respectively. Moreover, because of the effects of the pores embedded in composites

and the disorder orientations of GFs, the experimental values for the composites with 10-20 vol.% GFs were higher than the calculated values perpendicular to GFs basal plane, and the experimental values exceeded even the coefficient of thermal expansion of Al and GFs (25 ppm/K along the direction perpendicular to the basal plane).

- 2) A small difference between the orientations can cause the difference in the thermal conductivity of GFs to be as high as  $840 \text{ W m}^{-1} \text{ K}^{-1}$ ; The orientations of GFs in two-dimensional and three-dimensional images and two-dimensional aspect ratio of GFs have a strong correlation, and the correlation function can be obtained and use to convert the orientation of GFs from two-dimensional to three-dimensional. Moreover, the ETC of 10 vol% GFs/Al composite was evaluated using 2D image-based simulation and the correlation function. The results showed that the calculated ETC was equal to the experimental ETC  $238 \text{ W m}^{-1} \text{ K}^{-1}$  when the interfacial thermal conductance at the GFs-Al interface was  $1.8 \times 10^7 \text{ W m}^{-2} \text{ K}^{-1}$ , and the calculated ETC was  $250 \text{ W m}^{-1} \text{ K}^{-1}$  when the interfacial thermal conductance value was the same as that in a reference ( $4.5 \sim 5 \times 10^7 \text{ W m}^{-2} \text{ K}^{-1}$ ).
- 3) A two-dimensional image simulation code was developed to calculate the effective thermal conductivity of GFs/Al composites considering the effects of the orientation of GFs and interfacial thermal resistance; Orientations of GFs with respect to the heat flow direction were from  $11^\circ$  to  $33^\circ$  in 10-20 vol.% GFs/Al composites, and the orientation of GFs resulted in a loss of the effective thermal conductivity by 7.2% to 9.9%. According to the DMM method, the interfacial thermal conductance at the Al-GFs interface was  $1.1 \times 10^8 \text{ W m}^{-2} \text{ K}^{-1}$  in the in-plane and  $1.3 \times 10^9 \text{ W m}^{-2} \text{ K}^{-1}$  in the out-plane. In contrast, the result ( $1.1 \times 10^8 \text{ W m}^{-2} \text{ K}^{-1}$  in the in-plane) was much larger than the result of two-dimensional image simulation ( $1.56 \sim 1.78 \times 10^7 \text{ W m}^{-2} \text{ K}^{-1}$ ) and reference value ( $4.8 \times 10^7 \text{ W m}^{-2} \text{ K}^{-1}$ ).



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## **Published Papers in Regards to This Thesis**

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1. Yan Zhao, Kenjiro Sugio, Sasaki Gen, Zhefeng Xu, Jinku Yu. Evaluation of Effective Thermal Conductivity of Graphite Flake/Aluminum Composites by Two-Dimensional Image Simulation under a Correction Function. *Material transactions*, 62 (2021) 1625-1631. (Chapter 3)
2. Yan Zhao, Kenjiro Sugio, Yongbum Choi, Sasaki Gen, Zhefeng Xu, Jinku Yu. Effect of Anisotropic Thermal Conductivity of Graphite Flakes and Interfacial Thermal Resistance on the Effective Thermal Conductivity of Graphite Flakes/Aluminum Composites. *Material transactions*, 62 (2021) 98-104. (Chapter 4)
3. Yan Zhao, Kenjiro Sugio, Yong Bum Choi, Gen Sasaki and Jinku Yu. The relationship between the effective thermal conductivity and the orientation of graphite flakes in Al matrix composites. *Processing of the 10<sup>th</sup> Pacific Rim International Conference on Advanced Materials and Processing (PRICM10)*. (2019) 815-820. (Chapter 4)

## **Presentations**

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1. Yan Zhao, Kenjiro Sugio, Yong Bum Choi, Gen Sasaki and Jinku Yu. Effect of Anisotropic Thermal Conductivity of Carbon Fiber in Aluminum Matrix Composite. 日本金属学会 2018 年春期講演大会. March 19-21, 2018, Tokyo, Japan; Chiba Institute of Technology University.
2. Yan Zhao, Kenjiro Sugio, Yong Bum Choi, Gen Sasaki and Jinku Yu. The relationship between the effective thermal conductivity and the orientation of graphite flakes in Al matrix composites. The Tenth Pacific Rim International Conference on Advanced Materials and Processing (PRICM10). August 18-22, 2019, Xian, China, Qujiang International Conference Center.
3. Yan Zhao, Kenjiro Sugio, Sasaki Gen. Effect of anisotropic thermal conductivity and interfacial thermal resistance on the effective thermal conductivity of graphite flake/Al composites. 日本金属学会 2019 年秋期講演大会. September 11-13, 2019, Okayama, Japan; Okayama University Tsushima Campus.
4. Yan Zhao, Kenjiro Sugio, Sasaki Gen. Kenjiro Sugio, Sasaki Gen. Effect of anisotropic thermal conductivity and interfacial thermal resistance on the effective thermal conductivity of graphite flake/Al composites. 2018 年 12 月 日本金属学会 日本鉄鋼協会 中国四国支部 第 43 回若手フォーラム 広島, 日本.